

# § 3.1 热动平衡判据

## 1. 熵判据

孤立系  $dS \geq 0$

平衡态S极大。

对系统的状态虚变动，熵的虚变动  $\Delta S$

$$\Delta S < 0$$

孤立系统处在稳定平衡状态的必要和充分条件为:

$$\Delta S < 0$$

将**S**泰勒展开，准确到二级。有

$$S = S_0 + \delta S + \frac{1}{2} \delta^2 S$$

$$\Delta S = \delta S + \frac{1}{2} \delta^2 S$$

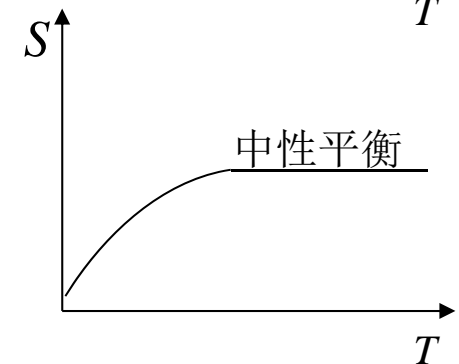
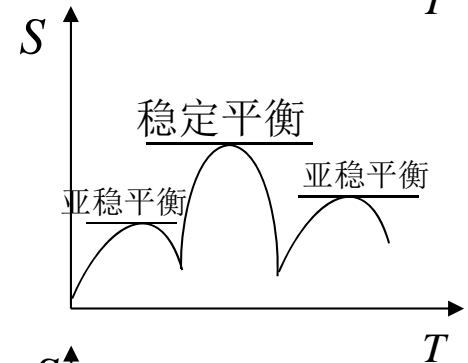
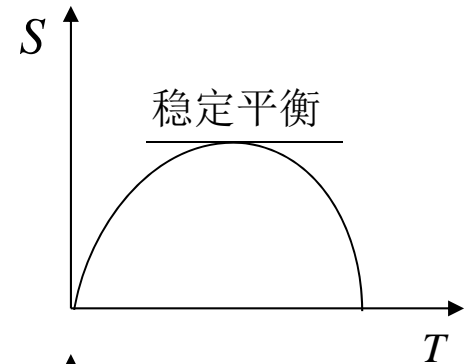
平衡态的必要条件  $\delta S = 0$

$\delta^2 S < 0$     $\Delta S < 0$    极大值   稳定平衡

最大极值   稳定平衡

较小极值   亚稳平衡

$\Delta S = 0$    常数值   中性平衡



## 2. 等温等容系统 ——自由能判据

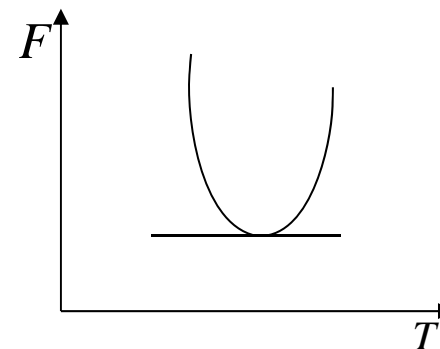
平衡态自由能最小

$$\Delta F > 0$$

$$\Delta F = \delta F + \delta_2 F$$

平衡条件:  $\delta F = 0$

稳定平衡:  $\delta^2 F > 0$



等温等容系统处在稳定平衡状态的必要和充分条件为:

$$\Delta F > 0$$

### 3. 等温等压系统 ——吉布斯函数判据

平衡态吉布斯函数最小  $\Delta G > 0$

$$\Delta G = \delta G + \delta^2 G$$

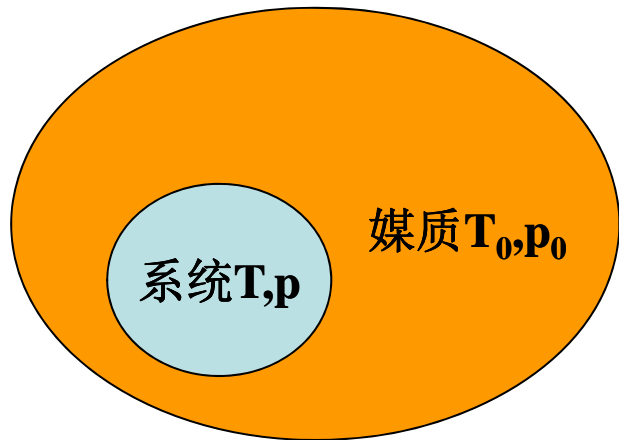
平衡条件:  $\delta G = 0$

稳定平衡:  $\delta^2 G > 0$

等温等压系统处在稳定平衡状态的必要和充分条件为:

$$\Delta G > 0$$

## 4. 孤立系统热动平衡条件和平衡稳定性条件



孤立大系统

媒质很大，有恒定的温度和压强。

整个系统的  $V$  不变， $U$  不变。

子系统虚变动  $\delta U$ ， $\delta V$

$$\delta U_0 + \delta U = 0,$$

$$\delta V_0 + \delta V = 0$$

熵的广延性  $\Delta \tilde{S} = \Delta S + \Delta S_0$

$$\Delta S = \delta S + \frac{1}{2} \delta^2 S$$

$$\Delta S_0 = \delta S_0 + \frac{1}{2} \delta^2 S_0$$

孤立系统极值要求  $\delta \tilde{S} = \delta S + \delta S_0 = 0$

$$dU = TdS - pdV \Rightarrow \delta S = \frac{1}{T} \delta U + \frac{p}{T} \delta V$$

$$\delta \tilde{S} = \delta S + \delta S_0 = \delta U \left( \frac{1}{T} - \frac{1}{T_0} \right) + \delta V \left( \frac{p}{T} - \frac{p_0}{T_0} \right) = 0$$

系统处于平衡状态  $T = T_0, p = p_0$  ——热动平衡条件

稳定平衡  $\delta^2 \tilde{S} = \delta^2 S_0 + \delta^2 S < 0$

$$\delta^2 S_0 \ll \delta^2 S$$

近似有  $\delta^2 \tilde{S} \approx \delta^2 S < 0$

$$S = S(U, V) \quad \delta S = \left( \frac{\partial S}{\partial U} \right)_V \delta U + \left( \frac{\partial S}{\partial V} \right)_U \delta V$$

$$\delta^2 S = \frac{\partial^2 S}{\partial U^2} (\delta U)^2 + \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial U \partial V} \delta U \delta V + \frac{\partial^2 S}{\partial V^2} (\delta V)^2$$

$$\begin{aligned}
\delta^2 S &= \delta U \left( \frac{\partial^2 S}{\partial U^2} \delta U + \frac{\partial^2 S}{\partial U \partial V} \delta V \right) + \delta V \left( \frac{\partial^2 S}{\partial U \partial V} \delta U + \frac{\partial^2 S}{\partial V^2} \delta V \right) \\
&= \delta U \left( \frac{\partial}{\partial U} \left( \frac{\partial S}{\partial U} \right)_V \delta U + \frac{\partial}{\partial V} \left( \frac{\partial S}{\partial U} \right)_V \delta V \right) + \delta V \left( \frac{\partial}{\partial U} \left( \frac{\partial S}{\partial V} \right)_U \delta U + \frac{\partial}{\partial V} \left( \frac{\partial S}{\partial V} \right)_U \delta V \right) \\
&= \delta U \delta \left( \frac{\partial S}{\partial U} \right)_V + \delta V \delta \left( \frac{\partial S}{\partial V} \right)_U
\end{aligned}$$

$$dU = TdS - pdV \Rightarrow dS = \frac{1}{T}dU + \frac{p}{T}dV$$

$$\left( \frac{\partial S}{\partial U} \right)_V = \frac{1}{T}, \quad \left( \frac{\partial S}{\partial V} \right)_U = \frac{p}{T}$$

$$\begin{aligned}
\delta^2 S &= \delta U \delta \left( \frac{1}{T} \right) + \delta V \delta \left( \frac{p}{T} \right) = \delta U \frac{-\delta T}{T^2} + \delta V \frac{T\delta p - p\delta T}{T^2} \\
&= \frac{1}{T} \left[ -\frac{\delta U + p\delta V}{T} \delta T + \delta p \delta V \right] = \frac{1}{T} [-\delta S \delta T + \delta p \delta V] < 0
\end{aligned}$$

以 (T, V) 为变量

$$\delta^2 S = \frac{1}{T} [-\delta S \delta T + \delta p \delta V]$$

$$= \frac{1}{T} \left[ -\left(\frac{\partial S}{\partial T}\right)_V (\delta T)^2 - \left(\frac{\partial S}{\partial V}\right)_T \delta V \delta T + \left(\frac{\partial p}{\partial T}\right)_V \delta T \delta V + \left(\frac{\partial p}{\partial V}\right)_T (\delta V)^2 \right]$$

$$= \frac{1}{T} \left[ -\frac{C_V}{T} (\delta T)^2 - \left(\frac{\partial p}{\partial T}\right)_V \delta V \delta T + \left(\frac{\partial p}{\partial T}\right)_V \delta T \delta V + \left(\frac{\partial p}{\partial V}\right)_T (\delta V)^2 \right]$$

$$= -\frac{C_V}{T^2} (\delta T)^2 + \frac{1}{T} \left(\frac{\partial p}{\partial V}\right)_T (\delta V)^2 < 0$$

平衡的稳定性条件:

$$C_V > 0, \quad \left(\frac{\partial p}{\partial V}\right)_T < 0$$



# § 3.2 开系热力学基本方程

## 1. 单元复相系的特点

- 复相系中的任一相都是均匀的开系，由于有相变发生，因而一个相的质量或摩尔数不守恒；整个系统是闭系。
- 复相系中每一相的平衡态热力学性质都可用状态参量来描述。各相的状态参量不完全独立，互相联系。

## 2. 开系热力学基本方程

$$dG = -SdT + Vdp + \mu dn \quad G(T, p, n)$$

$$\mu = \left( \frac{\partial G}{\partial n} \right)_{T, p}$$

叫系统的化学势。

单元单相  $\mu = G_m = \frac{G}{n} \quad G(T, p, n) = nG_m(T, p) = n\mu$

$$dG = -SdT + Vdp + \mu dn$$

$$S = -\left(\frac{\partial G}{\partial T}\right)_{p,n} \quad V = \left(\frac{\partial G}{\partial p}\right)_{T,n}$$

$$G = U - TS + pV$$

$$U = G + TS - pV$$

$$\begin{aligned} dU &= dG + TdS + SdT - pdV - Vdp \\ &= -SdT + Vdp + \mu dn + TdS + SdT - pdV - Vdp \end{aligned}$$

$$dU = TdS - pdV + \mu dn$$

$$U = U(S, V, n)$$

$$H = G + TS = U + pV$$

$$dH = TdS + Vdp + \mu dn$$

$$H = H(S, p, n)$$

$$F = G - pV = U - TS$$

$$dF = -SdT - pdV + \mu dn$$

$$F = F(T, V, n)$$

### 3. 巨热力势

$$J = F - \mu n = F - G = -pV$$

$$dJ = dF - \mu dn - nd\mu = -SdT - pdV + \mu dn - \mu dn - nd\mu$$

$$dJ = -SdT - pdV - nd\mu$$

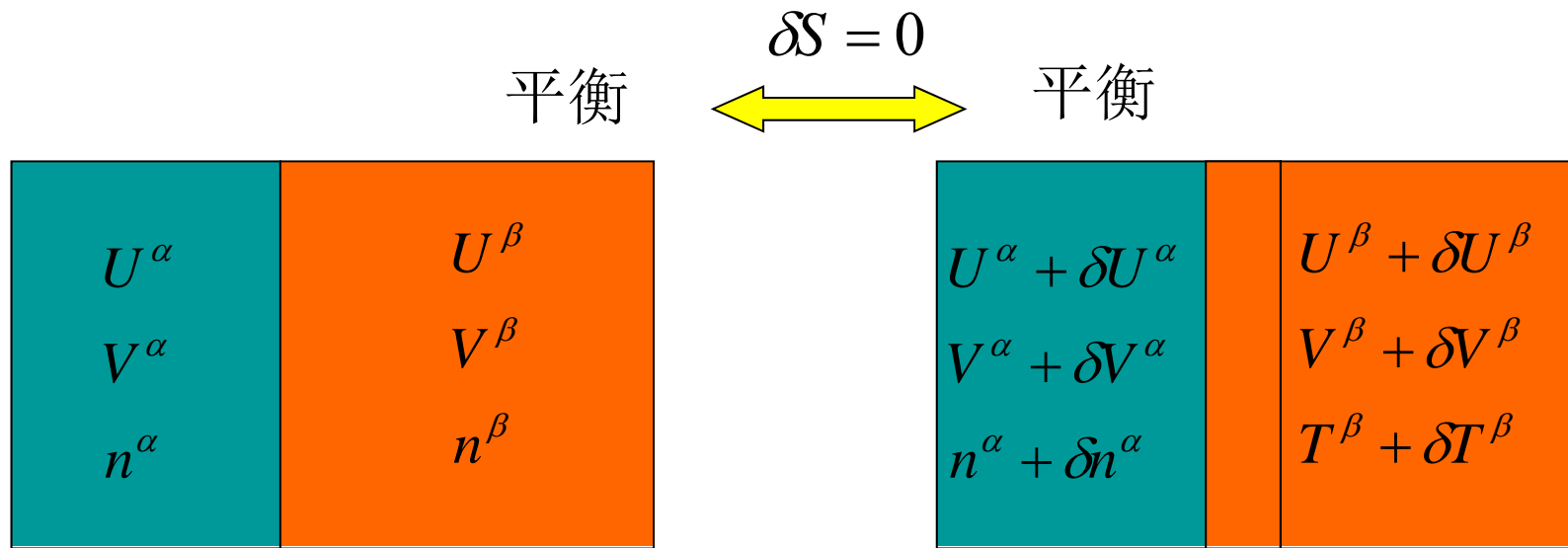
$$J = J(T, V, \mu)$$

$$S = -\left(\frac{\partial J}{\partial T}\right)_{V, \mu} \quad p = -\left(\frac{\partial J}{\partial V}\right)_{T, \mu} \quad n = -\left(\frac{\partial J}{\partial \mu}\right)_{T, V}$$

# § 3.3 单元系的复相平衡

## 1. 单元两相系

$\alpha$ 、 $\beta$ 两相 (或两子系)  $\longrightarrow$  孤立系



$$U^\alpha + U^\beta = \text{常量}$$

$$V^\alpha + V^\beta = \text{常量}$$

$$n^\alpha + n^\beta = \text{常量}$$

$$\delta U^\alpha + \delta U^\beta = 0$$

$$\delta V^\alpha + \delta V^\beta = 0$$

$$\delta n^\alpha + \delta n^\beta = 0$$

$$dU = TdS - pdV + \mu dn \quad \Rightarrow$$

$$\delta S^\alpha = \frac{\delta U^\alpha + p^\alpha \delta V^\alpha - \mu^\alpha \delta n^\alpha}{T^\alpha}$$

$$\delta S^\beta = \frac{\delta U^\beta + p^\beta \delta V^\beta - \mu^\beta \delta n^\beta}{T^\beta}$$

$$\delta S = \delta S^\alpha + \delta S^\beta$$

$$= \delta U^\alpha \left( \frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left( \frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left( \frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right)$$

## 2. 平衡条件

总熵极大值,  $\delta S = 0$

$$\frac{1}{T^\alpha} - \frac{1}{T^\beta} = 0$$

$$\frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} = 0$$

$$\frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} = 0$$

热平衡条件  $T^\alpha = T^\beta$

力平衡条件  $p^\alpha = p^\beta$

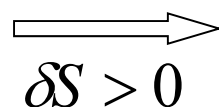
相平衡条件  $\mu^\alpha = \mu^\beta$

□ 复相系的平衡条件：整个单元二相系达到平衡时，两相的温度，压强和化学势必须分别相等。

此结论对于三相，四相等复相系均适用。

讨论:

非平衡



平衡

$$\delta S = \delta U^\alpha \left( \frac{1}{T^\alpha} - \frac{1}{T^\beta} \right) + \delta V^\alpha \left( \frac{p^\alpha}{T^\alpha} - \frac{p^\beta}{T^\beta} \right) - \delta n^\alpha \left( \frac{\mu^\alpha}{T^\alpha} - \frac{\mu^\beta}{T^\beta} \right)$$

① 若热平衡条件不满足, 即  $T^\alpha > T^\beta, p^\alpha = p^\beta, \mu^\alpha = \mu^\beta$

则  $\delta U_\alpha < 0$  即热量从高温( $\alpha$ )相传递到低温( $\beta$ )相

② 若力学平衡条件不满足, 即  $p^\alpha > p^\beta, T^\alpha = T^\beta, \mu^\alpha = \mu^\beta$

则  $\delta V_\alpha > 0$  压强大的相( $\alpha$ )将膨胀, 压强小的相( $\beta$ )将压缩。

③ 若相变平衡条件不满足, 即  $\mu^\alpha > \mu^\beta, p^\alpha = p^\beta, T^\alpha = T^\beta$

则  $\delta n_\alpha < 0$  物质由化学势高的相( $\alpha$ )转移到化学势低的相( $\beta$ )

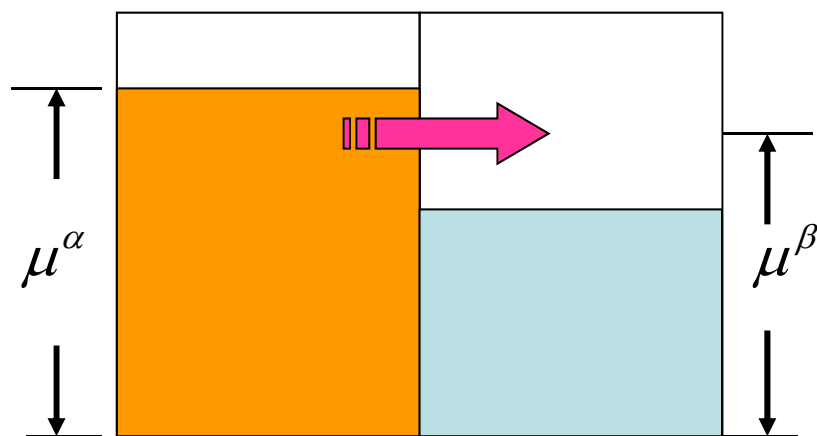
# 化学势

$$\mu^\alpha > \mu^\beta \Rightarrow \delta n^\alpha < 0$$



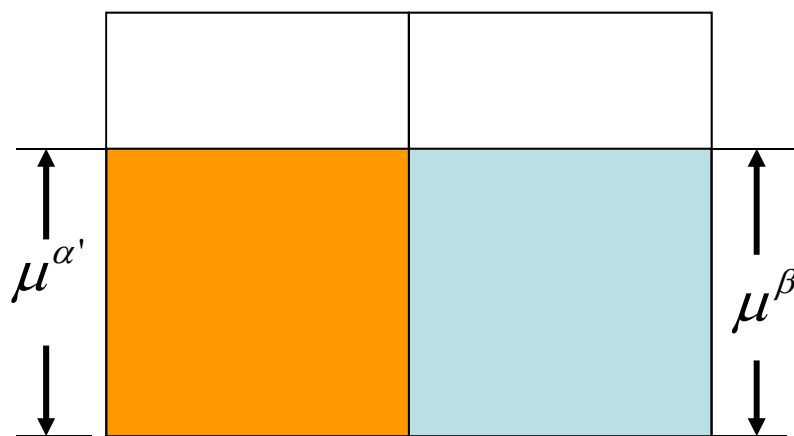
粒子方向

粒子从化学势高的相向低的相跑！！



化学不平衡

$$\mu^\alpha > \mu^\beta$$



化学平衡

$$\mu^{\alpha'} = \mu^{\beta'}$$

化学势差促使粒子流动。



作业： P106:

3.1; 3.2; 3.3; 3.4