

## Review:

动力学普遍方程: 
$$\sum_{i=1}^n (\vec{F}_i - m_i \ddot{\vec{r}}_i) \cdot \delta \vec{r}_i = 0$$

or 
$$\sum_{i=1}^{3n} (F_i - m_i \ddot{x}_i) \cdot \delta x_i = 0$$

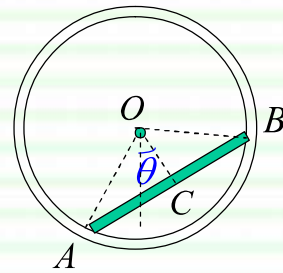
基本形式拉格朗日方程: 
$$\frac{d}{dt} \left( \frac{\partial T}{\partial \dot{q}_\alpha} \right) - \frac{\partial T}{\partial q_\alpha} = Q_\alpha$$

保守系的拉格朗日方程: 
$$Q_\alpha = - \frac{\partial V}{\partial q_\alpha}$$
$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}_\alpha} \right) - \frac{\partial L}{\partial q_\alpha} = 0$$

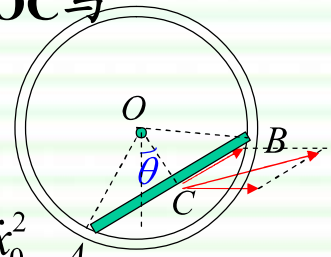
本次课:小振动和冲击运动

上次例题：

**例题2**、图中放在一直线轨道上的圆环O的质量为 $m_2$ ，半径为 $r$ 。圆环与轨道间有足够的摩擦力阻止圆环滑动。圆环内有一质量为 $m_1$ ，长为 $l$ 的匀质杆AB，其中点C与圆环的中心O的距离为 $r/2^{1/2}$ ，则杆长 $l=2^{1/2}r$ 。杆与圆环间摩擦不计。写出系统的运动微分方程。



解:  $i=2$ , 取圆环中心  $O$  的坐标  $x_0$  以及  $OC$  与垂直线的交角  $\theta$  为广义坐标。



圆环作纯滚动, 则:  $\omega = \frac{\dot{x}_0}{r}$

圆环动能:  $T_2 = \frac{1}{2} m_2 \dot{x}_0^2 + \frac{1}{2} (m_2 r^2) \omega^2 = m_2 \dot{x}_0^2$

杆 AB 动能:

$$T_1 = \frac{1}{2} m_1 (\dot{x}_0^2 + \frac{1}{2} r^2 \dot{\theta}^2 + \sqrt{2} r \dot{x}_0 \dot{\theta} \cos \theta) + \frac{1}{2} \cdot \frac{1}{12} m_1 (\sqrt{2} r)^2 \dot{\theta}^2$$

总动能:

$$T = T_1 + T_2 = (m_2 + \frac{m_1}{2}) \dot{x}_0^2 + \frac{1}{3} m_1 r^2 \dot{\theta}^2 + \frac{\sqrt{2}}{2} m_1 r \dot{x}_0 \dot{\theta} \cos \theta$$

零势面:  $O$  点的水平面:  $V = -\frac{\sqrt{2}}{2} m_1 r g \cos \theta$

$$\begin{aligned} \therefore L &= T - V \\ &= (m_2 + \frac{m_1}{2})\dot{x}_0^2 + \frac{1}{3}m_1r^2\dot{\theta}^2 + \frac{\sqrt{2}}{2}m_1rx_0\dot{\theta}\cos\theta + \frac{\sqrt{2}}{2}m_1rg\cos\theta \end{aligned}$$

代入拉格朗日方程中：

$$x_0: (2m_2 + m_1)\ddot{x}_0 + \frac{\sqrt{2}}{2}m_1r(\ddot{\theta}\cos\theta - \dot{\theta}^2\sin\theta) = 0$$

$$\theta: 4r\ddot{\theta} + 3\sqrt{2}\ddot{x}_0\cos\theta + 3\sqrt{2}g\sin\theta = 0$$

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只考虑微振动则：  $(2m_2 + m_1)\ddot{x}_0 + \frac{\sqrt{2}}{2}m_1r\ddot{\theta} = 0$

$$4r\ddot{\theta} + 3\sqrt{2}\ddot{x}_0 + 3\sqrt{2}g\theta = 0$$

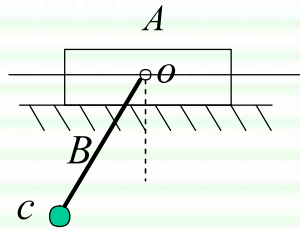
消去  $\ddot{x}_0$  得：

$$\frac{8m_2 + m_1}{2m_2 + m_1}r\ddot{\theta} + 3\sqrt{2}g\theta = 0$$

$$\Rightarrow \ddot{\theta} + \underbrace{3\sqrt{2}\frac{g}{r}\frac{2m_2 + m_1}{8m_2 + m_1}}_{\omega^2}\theta = 0$$

此乃小振动问题！

例题3、如图示，物块A沿水平面滑动，细杆B一端用铰链与A相连，另一端与球C相固连所组成椭圆摆。设细杆B长为 $l$ ，A和C的质量分别为 $m_1$ 和 $m_2$ ，细杆B重量不计，求：椭圆摆运动微分方程。



解: 1) A、C组成系统,  $i=2$

平面坐标系  $O_1xy$ , 广义坐标  $\theta$  和  $x_1$

2) 先找A和C( $x_2, y_2$ )的坐标:

$$\begin{cases} x_2 = x_1 + l \sin \theta \\ y_2 = l \cos \theta \end{cases} \Rightarrow \begin{cases} \dot{x}_2 = \dot{x}_1 + l \dot{\theta} \cos \theta \\ \dot{y}_2 = -l \dot{\theta} \sin \theta \end{cases}$$

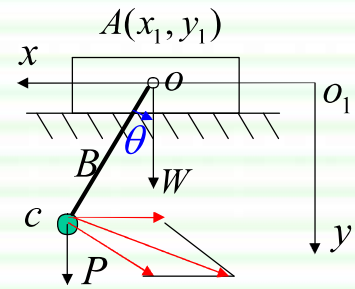
**A:**  $T_A = \frac{1}{2} m_1 \dot{x}_1^2$

**C:**  $v_c = (\dot{x}_2^2 + \dot{y}_2^2)^{1/2} = (\dot{x}_1^2 + l^2 \dot{\theta}^2 + 2l\dot{x}_1\dot{\theta} \cos \theta)^{1/2}$

**Or:**  $|\vec{v}_c|^2 = |\dot{x}_1 \vec{i} + l \dot{\theta} \vec{e}_{\perp oc}|^2 = \dot{x}_1^2 + l^2 \dot{\theta}^2 - 2\dot{x}_1 l \dot{\theta} \cos(\pi - \theta)$

$$\therefore T_c = \frac{1}{2} m_2 v_c^2 = \frac{1}{2} m_2 (\dot{x}_1^2 + l^2 \dot{\theta}^2 + 2l\dot{x}_1\dot{\theta} \cos \theta)$$

总动能:  $T = \frac{1}{2} (m_1 + m_2) \dot{x}_1^2 + \frac{1}{2} m_2 l^2 \dot{\theta}^2 + m_2 l \dot{x}_1 \dot{\theta} \cos \theta$



势能:  $V = -m_2 gl \cos \theta$

$\therefore L = T - V$

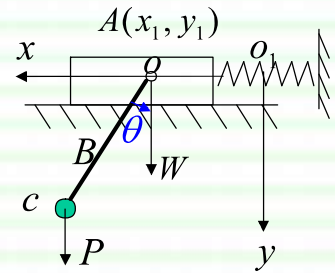
$$= \frac{1}{2}(m_1 + m_2)\dot{x}_1^2 + \frac{1}{2}m_2 l^2 \dot{\theta}^2 + m_2 l \dot{x}_1 \dot{\theta} \cos \theta + m_2 gl \cos \theta$$

代入拉格朗日方程中:  $\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) - \frac{\partial L}{\partial q} = 0 \quad i=2$

考虑微小摆动:

$$\ddot{\theta} + \frac{m_1 + m_2}{m_1} \cdot \frac{g}{l} \theta = 0$$

若A用弹簧连接在固定点处, ?



此时系统势能将附加弹性势能项:  $\frac{1}{2} k x_1^2$

解:  $i=2$  广义坐标  $\theta$  和  $x_1$

$$T_A = \frac{1}{2} m_1 \dot{x}_1^2 \quad \vec{v}_c = \dot{x}_1 \vec{i} + l \dot{\theta} \vec{e}_{cp}$$

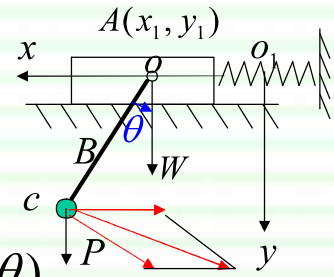
$$\therefore v_c^2 = \dot{x}_1^2 + l^2 \dot{\theta}^2 - 2\dot{x}_1 l \dot{\theta} \cos(\pi - \theta)$$

$$\therefore T_c = \frac{1}{2} m_2 v_c^2 = \frac{1}{2} m_2 (\dot{x}_1^2 + l^2 \dot{\theta}^2 + 2l\dot{x}_1 \dot{\theta} \cos \theta)$$

总动能:  $T = \frac{1}{2} (m_1 + m_2) \dot{x}_1^2 + \frac{1}{2} m_2 l^2 \dot{\theta}^2 + m_2 l \dot{x}_1 \dot{\theta} \cos \theta$

势能:  $V = -m_2 g l \cos \theta + \frac{1}{2} k x_1^2 \quad \therefore L = T - V$

代入拉格朗日方程中:  $\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) - \frac{\partial L}{\partial q} = 0$





例题4、两个相同的塔轮；半径为 $r$ 的鼓轮上绕有细绳，轮I绕有铅直弹簧( $k$ )，轮II挂一重物。塔轮对轴的转动惯量均为 $J$ ，重物质量为 $m$ 。求此振动系统的固有频率。

解：  $i=1$ ，取 $q=x$ ，  $\dot{x} = \omega r$

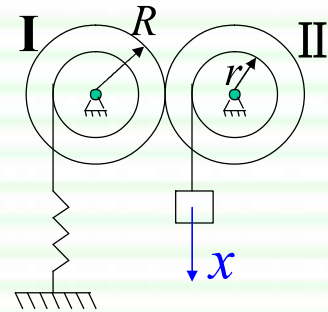
系统动能：  $T = \frac{1}{2}m\dot{x}^2 + 2 \times \frac{1}{2}J\left(\frac{\dot{x}}{r}\right)^2$

系统势能：  $V = \frac{1}{2}kx^2 - mgx$

$$\therefore L = T - V = \frac{1}{2}m\dot{x}^2 + J\left(\frac{\dot{x}}{r}\right)^2 - \frac{1}{2}kx^2$$

代入保守系的拉格朗日方程中：

$$\Rightarrow \left(m + \frac{2J}{r^2}\right) \ddot{x} + kx = 0 \quad \Rightarrow \omega_0 = \sqrt{\frac{kr^2}{2J + mr^2}}$$



例题5、质量为 $m$ 的质点悬挂在轻绳一端，绳的另一端绕在半径为 $r$ 的固定圆柱体上，构成一个摆。设平衡时绳的下垂部分长为 $l$ ，试求摆的运动微分方程。

解：  $i=1$ ，取  $q=\theta$

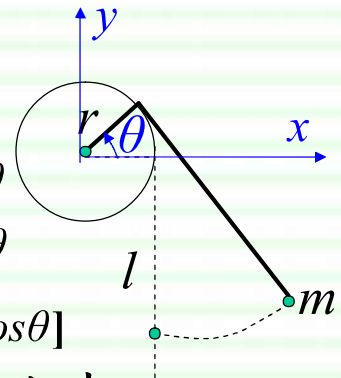
系统动能：  $T = \frac{1}{2}m[(l+r\dot{\theta})^2]$

$$\begin{cases} x = r \cos \theta + (l+r\theta) \sin \theta \\ y = r \sin \theta - (l+r\theta) \cos \theta \end{cases} \Rightarrow \begin{cases} \dot{x} = \dot{\theta}(l+r\theta) \cos \theta \\ \dot{y} = \dot{\theta}(l+r\theta) \sin \theta \end{cases}$$

系统势能：  $V = mg[(l+r \sin \theta) - (l+r\theta) \cos \theta]$

$\therefore L = T - V$  代入保守系的拉格朗日方程中：

$$\begin{aligned} (l+r\theta) \ddot{\theta} + r\dot{\theta}^2 + g \sin \theta &= 0 \\ \Rightarrow \ddot{\theta} + \frac{g}{l} \theta &= 0 \end{aligned}$$



### 冲击运动的拉格朗日方程

冲击运动的特点：

- 1、冲力大；作用时间短；
- 2、可忽略冲击造成的位形变化；

3、 $\lim_{\Delta t \rightarrow 0} \int_t^{t+\Delta t} \vec{F} dt$  是有限的。

一、矢量力学关于质点冲击运动的动力学方程:

$$\therefore m \frac{d\vec{v}}{dt} = \vec{F}$$

$$\therefore \lim_{\Delta t \rightarrow 0} \int_t^{t+\Delta t} m \frac{d\vec{v}}{dt} \cdot dt = \lim_{\Delta t \rightarrow 0} \int_t^{t+\Delta t} \vec{F} dt$$

$$m \left[ \lim_{\Delta t \rightarrow 0} \vec{V}(t + \Delta t) - \vec{V}(t) \right] = \vec{I}$$

$$\Rightarrow \Delta \vec{P} = \vec{I}$$

二、分析力学关于系统的冲击运动的动力学方程:

$$\Rightarrow \Delta P_k = I_k$$

$$\therefore \frac{d}{dt} \left( \frac{\partial T}{\partial \dot{q}_k} \right) - \frac{\partial T}{\partial q_k} = Q_k \quad (k = 1, 2 \dots s)$$

以  $dt$  乘上式且对碰撞时间  $\Delta t$  积分:

$$\therefore \int_t^{t+\Delta t} d\left(\frac{\partial T}{\partial \dot{q}_k}\right) - \int_t^{t+\Delta t} \frac{\partial T}{\partial q_k} dt = \int_t^{t+\Delta t} Q_k dt$$

$\because \Delta t \rightarrow 0$

$$?? = \lim_{\Delta t \rightarrow 0} \int_t^{t+\Delta t} \left( \sum_{i=1}^{3n} F_i \frac{\partial x_i}{\partial q_k} \right) dt = \sum_{i=1}^{3n} \left[ \lim_{\Delta t \rightarrow 0} \left( \int_t^{t+\Delta t} F_i dt \right) \frac{\partial x_i}{\partial q_k} \right]$$

$$= \sum_{i=1}^{3n} I_i \frac{\partial x_i}{\partial q_k} = I_k \quad \text{广义冲量}$$

$$?? = \int_t^{t+\Delta t} dP_k = P_{k2} - P_{k1}$$

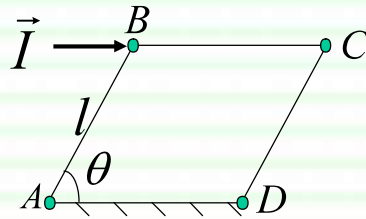
$$\therefore P_{k2} - P_{k1} = I_k \quad (k = 1, 2, \dots, s)$$

$$\text{即: } \left[ \frac{\partial T}{\partial \dot{q}_k} \right]_{t_2} - \left[ \frac{\partial T}{\partial \dot{q}_k} \right]_{t_1} = I_k \quad (k = 1, 2, \dots, s)$$

一般取零

关于系统冲击运动的拉格朗日方程。

**例题1:** 铰链平行四边形ABCD置于光滑水平面上，AB，CD杆的质量均为 $m_1$ ，BC杆质量为 $m_2$ ，AD杆固定不动， $\angle BAD = \theta$ 。现于B点处作用一冲量 $\vec{I}$ ，如图示。求B点速度及整个系统所获得的动能。



解:  $i=1$ . 取  $q=\theta$

杆AB和杆BC的质心速度分别为:

$$u_1 = \frac{l}{2} \dot{\theta}, \quad u_2 = l \dot{\theta}$$

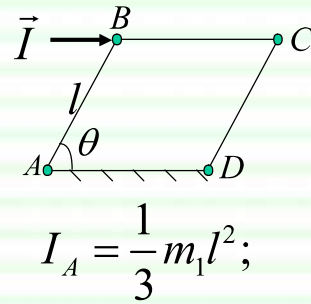
系统的动能为:

$$T = 2 \times \frac{1}{2} I_A \dot{\theta}^2 + \frac{1}{2} m_2 u_2^2 = \left( \frac{m_1}{3} + \frac{m_2}{2} \right) l^2 \dot{\theta}^2$$

$\therefore dx_B = -l d\theta \sin \theta$  广义冲量:  $I_k = I \frac{\partial x_B}{\partial \theta} = -Il \sin \theta$

代入  $\left( \frac{\partial T}{\partial \dot{\theta}} \right)_2 - \left( \frac{\partial T}{\partial \dot{\theta}} \right)_1 = I_k$  中:

$$\Rightarrow u_2 = \frac{-3I \sin \theta}{2m_1 + 3m_2} \quad T = \frac{3I^2 \sin^2 \theta}{2(2m_1 + 3m_2)}$$



**例题2:** 两个长均为 $l$ ，质量均为 $m$ 的均质杆在 $A$ 处铰链后悬挂在 $O$ 轴上，并在 $B$ 端受到冲量 $I$ 的作用，求碰后两杆的角速度。

**解:**  $i=2$ 。 取 $q_1=\varphi_1$ 和 $q_2=\varphi_2$

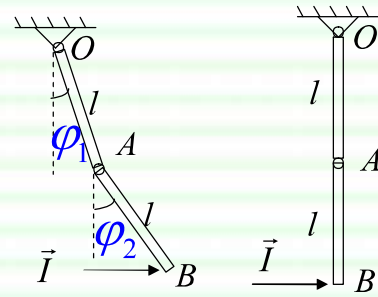
两杆质心速度分别为:

$$u_1 = \frac{1}{2}l\dot{\varphi}_1, \quad u_2 = l\dot{\varphi}_1 + \frac{l}{2}\dot{\varphi}_2$$

$$\therefore T = \frac{1}{2}I_0\dot{\varphi}_1^2 + \frac{1}{2}mu_2^2 + \frac{1}{2}I_{C_2}\dot{\varphi}_2^2$$

其中:  $I_0 = \frac{1}{3}ml^2; \quad I_{C_2} = \frac{1}{12}ml^2$

$$\therefore T = \frac{1}{2}ml^2 \left( \frac{4}{3}\dot{\varphi}_1^2 + \dot{\varphi}_1\dot{\varphi}_2 + \frac{1}{3}\dot{\varphi}_2^2 \right)$$



$$\left( \frac{\partial T}{\partial \dot{q}_k} \right)_2 - \left( \frac{\partial T}{\partial \dot{q}_k} \right)_1 = I_k$$



$$T = \frac{1}{2} ml^2 \left( \frac{4}{3} \dot{\varphi}_1^2 + \dot{\varphi}_1 \dot{\varphi}_2 + \frac{1}{3} \dot{\varphi}_2^2 \right)$$

现求广义冲量:  $I_k = \sum_{i=1}^n \vec{I}_i \cdot \frac{\partial \vec{r}_i}{\partial q_k} = \vec{I} \cdot \frac{\partial \vec{r}_B}{\partial q_k} = I \frac{\partial x_B}{\partial q_k}$

$$\because x_B = l \sin \varphi_1 + l \sin \varphi_2$$

$$\therefore I_{\varphi_1} = I \frac{\partial x_B}{\partial \varphi_1} = Il \cos \varphi_1 \xrightarrow{\varphi_1 \rightarrow 0} Il \quad \text{同理: } I_{\varphi_2} = Il$$

$$\Rightarrow \dot{\varphi}_1 = -\frac{6I}{7ml}; \quad \dot{\varphi}_2 = \frac{30I}{7ml}$$

? 若I作用点在AB杆质心处, 情况又如何?

$$I'_{\varphi_1} \text{ 不变 } \frac{l}{2} \delta \varphi_2 = \frac{Il}{2} \Rightarrow \dot{\varphi}_1 = \frac{3I}{7ml}; \quad \dot{\varphi}_2 = \frac{6I}{7ml}$$

**小论文：**图示水平面内行星轮系中，长为 $l$ 的曲柄 $OA$ 绕 $O$ 轴转动，其端点 $A$ 由铰链连接一 $r_2$ 的齿轮 $\text{II}$ ，已知 $r_2=1.5r_1$ ，曲柄 $OA$ 质量为 $m$ ，齿轮 $\text{I}$ 、 $\text{II}$ 的质量分别为 $m_1$ 、 $m_2$ 。在曲柄上作用一力偶 $M$ ，齿轮 $\text{I}$ 上有阻力偶 $M_1$ ，求曲柄运动方程。

**解：** $i=1$ ，选曲柄转角 $\phi$ 为广义坐标，轮心 $A$ 也是曲柄端点 $A$ 。

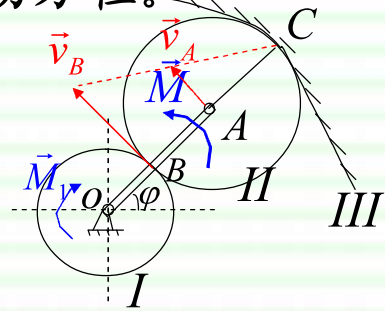
$$\therefore v_A = (r_1 + r_2)\omega = l\dot{\phi} \quad v_C = 0$$

$$\text{而：} \quad v_C = v_A - r_2\dot{\phi}_{\text{II}}$$

$$\text{则轮II的角速度：} \quad \dot{\phi}_{\text{II}} = \frac{l\dot{\phi}}{r_2} = \frac{5}{3}\dot{\phi}$$

又轮 $\text{I}$ 与轮 $\text{II}$ 的啮合处 $B$ 的速度： $v_B = 2v_A = 2l\dot{\phi}$

$$\text{则轮I的角速度：} \quad \dot{\phi}_{\text{I}} = \frac{v_B}{r_1} = \frac{2l\dot{\phi}}{r_1} = 5\dot{\phi}$$



曲柄的动能:

$$T_{qb} = \frac{1}{2} J_{qbo} \dot{\phi}^2 = \frac{1}{2} \cdot \frac{1}{3} ml^2 \dot{\phi}^2 = \frac{1}{6} ml^2 \dot{\phi}^2$$

齿轮I的动能:

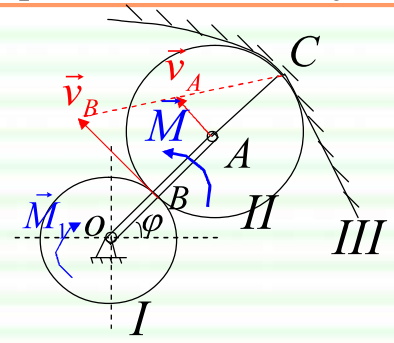
$$T_I = \frac{1}{2} J_{Io} \dot{\phi}_I^2 = \frac{1}{2} \cdot \frac{1}{2} m_1 r_1^2 \dot{\phi}_I^2 = m_1 l^2 \dot{\phi}^2$$

齿轮II的动能:

$$\begin{aligned} T_{II} &= \frac{1}{2} m_2 v_A^2 + \frac{1}{2} J_{IIA} \dot{\phi}_{II}^2 \\ &= \frac{1}{2} m_2 (l\dot{\phi})^2 + \frac{1}{2} \cdot \frac{1}{2} m_2 r_2^2 \left(\frac{l\dot{\phi}}{r_2}\right)^2 = \frac{3}{4} m_2 l^2 \dot{\phi}^2 \end{aligned}$$

系统的总动能:

$$T = \left(\frac{1}{6} m + m_1 + \frac{3}{4} m_2\right) l^2 \dot{\phi}^2$$



计算广义力:

给曲柄一虚位移  $\delta \varphi$ , 则  $\delta \varphi_I = 5 \delta \varphi$

则主动力的虚功为:

$$\delta W = M\delta\varphi - M_1\delta\varphi_I = M\delta\varphi - 5M_1\delta\varphi = (M - 5M_1)\delta\varphi$$

代入拉格朗日方程中得:

$$\left(\frac{1}{3}m + 2m_1 + \frac{3}{2}m_2\right)l^2\ddot{\varphi} = M - 5M_1$$

$$\Rightarrow \ddot{\varphi} = \frac{M - 5M_1}{\left(\frac{1}{3}m + 2m_1 + \frac{3}{2}m_2\right)l^2}$$

设  $t=0$  时,  $\varphi = \varphi_0$ ;  $\dot{\varphi} = \dot{\varphi}_0$

$$\Rightarrow \varphi = \frac{M - 5M_1}{\left(\frac{1}{3}m + 2m_1 + \frac{3}{2}m_2\right)l^2}t^2 + \dot{\varphi}_0 t + \varphi_0$$

作业：一、习题5.11、5.12、5.15



下次课:初积分问题