# An Explicit Harmonic Extension for The Constant-like Basis And Its Application to Domain Decompositions ${ }^{1}$ 

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#### Abstract

In this paper we are concerned with substructuring methods for the second-order elliptic problems in three-dimensional domains. We first design a simple and completely explicit nearly harmonic extension for the constant-like basis function (i.e., the face basis function), and then define a coarse subspace based on this nearly harmonic extension. Own to the resulting coarse solver, we develop a kind of substructuring preconditioner with inexact solvers. We show that the condition number of the preconditioned system grows only as the logarithm of the dimension of the local problem associated with an individual substructure, and is independent of possible jumps of the coefficient in the elliptic equation. Numerical experiments confirms the theoretical results.


Key words. domain decomposition, coarse subspace, constant-like basis, nearly harmonic extension, preconditioner, inexact solvers, condition number

AMS(MOS) subject classification 65F10, 65N30, 65N55

## 1 Introduction

Non-overlapping domain decomposition methods (DDMs) have been shown to be powerful techniques for solving partial differential equations, especially for the case with large jump coefficients. One's main task in non-overlapping DDMs is the construction of an efficient substructuring preconditioner for discretization system associated with the underlying partial differential equations. The construction of this preconditioner has been investigated from various ways and to various models in literature, see, for example, [1]-[11], [13]-[18], [20]-[24], [26]-[27], [29].

Most non-overlapping DDMs studied so far require exact subdomain solvers; we refer [28] and [32] (and the references cited therein). Such a requirement severely degrade the efficiency of the methods. There are only a few works studying substructuring methods with inexact subdomain solvers [3], [4], [11] and [6]. The essential difficulty is that discrete harmonic extensions on each subdomain are used in non-overlapping domain decomposition methods. In [3], analysis and numerical experiments with inexact algorithms of NeumannDirichlet type was done under the additional assumption of high accuracy of the inexact solvers. In [4], the harmonic extension on a subdomain was replaced by a simple average extension, and substructuring preconditioners with the average extension are constructed. Because of such average extension, nearly optimal convergence can not be gotten for these substructuring preconditioners. To avoid harmonic extensions, [11] considered so called approximate harmonic basis functions, which still involve high accuracy of the inexact solvers. Another way to construct substructuring preconditioner with inexact solvers was considered in [6] (mainly for two dimensions). In this preconditioner, overlapping face subspaces are used, and harmonic extension is used only in the definition of coarse subspace. How to design nearly harmonic extensions involved in the coarse subspace is the core problem in

[^0]the construction of such substructuring preconditioner. In the case with three dimensions, the coarse subspace may consist of the edge basis and face basis, which are constant-like functions. Then, one needs only to define suitable extensions for constant-like functions on faces. Such extensions, which depend on the geometric shapes of the subdomains, was first studied in [8].

In the present paper we introduce a new variant of substructuring preconditioner with inexact solvers considered in [6] (see also [32]). The main contribution of the paper is the design of a simple and completely explicit extension for the constant-like basis function. This extension plays a key role in the substructuring preconditioner with inexact solvers. We show that the new substructuring preconditioner possesses nearly optimal convergence, which is independent of possible large jumps of the coefficient across the interface. For the new method, no additional assumption is required.

The outline of the remainder of the paper is as follows. In Section 2, we introduce some notation and our motive. In Section 3, we present an explicit extension of the constantlike function. The results on the substructuring preconditioner with inexact solvers are described in Section 4. In Section 5, we prove the stability of the extension of constant-like function, which is used in Section 4. Some numerical results are reported in Section 6.

## 2 Preliminaries

### 2.1 Domain decomposition

Let $\Omega$ be a bounded polyhedron in $\mathcal{R}^{3}$. Consider the model problem

$$
\left\{\begin{array}{c}
-\operatorname{div}(\omega \nabla u)=f, \quad \text { in } \Omega,  \tag{2.1}\\
u=0, \quad \text { on } \partial \Omega
\end{array}\right.
$$

where $\omega \in L^{\infty}(\Omega)$ is a positive function.
Let $H_{0}^{1}(\Omega)$ denote the standard Sobolev space, and let $(\cdot, \cdot)$ denote the $L^{2}(\Omega)$-inner product. The weak formulation of $(2.1)$ in $H_{0}^{1}(\Omega)$ is then given by the following.

Find $u \in H_{0}^{1}(\Omega)$ such that

$$
\begin{equation*}
\mathcal{A}(u, v)=(f, v) \quad \forall v \in H_{0}^{1}(\Omega) \tag{2.2}
\end{equation*}
$$

where $(\cdot, \cdot)$ is the scalar product in $L^{2}(\Omega)$, and

$$
\mathcal{A}(u, v)=\int_{\Omega} \omega \nabla u \cdot \nabla v d p
$$

We will apply a kind of non-overlapping domain decomposition method to solving (2.2). For simplicity of exposition, we consider only the case with matching grids in this paper.

Let $\mathcal{T}_{h}=\left\{\tau_{i}\right\}$ be a regular and quasi-uniform triangulation of $\Omega$ with $\tau_{i}^{\prime} s$ being nonoverlapping simplexes of size $h(\in(0,1])$. The set of notes of $\mathcal{T}_{h}$ is denoted by $\mathcal{N}_{h}$. We then define $V_{h}(\Omega)$ to be the piecewise linear finite element subspace of $H_{0}^{1}(\Omega)$ associated with $\mathcal{T}_{h}$ :

$$
V_{h}(\Omega)=\left\{v \in H_{0}^{1}(\Omega):\left.v\right|_{\tau} \in \mathcal{P}_{1} \quad \forall \tau \in \mathcal{T}_{h}\right\}
$$

where $\mathcal{P}_{1}$ is the space of linear polynomials. Then the finite element approximation for (2.2) is to find $u_{h} \in V_{h}(\Omega)$ such that

$$
\begin{equation*}
\mathcal{A}\left(u_{h}, v_{h}\right)=\left(f, v_{h}\right), \quad \forall v_{h} \in V_{h}(\Omega) \tag{2.3}
\end{equation*}
$$

Let $\Omega$ be decomposed into the union of N polyhedrons $\Omega_{1}, \cdots, \Omega_{\mathrm{N}}$, which satisfy $\Omega_{i} \cap \Omega_{j}=$ $\emptyset$ when $i \neq j$. We assume that each $\partial \Omega_{k}$ can be written as a union of boundaries of
elements in $\mathcal{T}_{h}$, and all $\Omega_{k}$ are of size $H$ in the usual sense (see [5] and [32]). Without loss of generality, we assume that the coefficient $\omega(p)$ is piecewise constant, then each subdomain $\Omega_{k}$ is chosen such that $\omega(p)$ equals to a constant $\omega_{k}$ in $\Omega_{k}$. Note that $\left\{\Omega_{k}\right\}$ may not constitute a triangulation of $\Omega$.

The common part of two neighboring subdomains $\Omega_{i}$ and $\Omega_{j}$ may be a vertex, an edge or a face. In particular, we denote by $\Gamma_{i j}$ the common face of two neighboring subdomains $\Omega_{i}$ and $\Omega_{j}$ (i.e., $\Gamma_{i j}=\partial \Omega_{i} \cap \partial \Omega_{j}$ ). The union of all $\Gamma_{i j}$ is denoted by $\Gamma$, which is called the interface. In this paper, we choose Dirichlet data as the interface unknown.

Define the operator $A_{h}: V_{h}(\Omega) \rightarrow V_{h}(\Omega)$ by

$$
\left(A_{h} v, w\right)=\mathcal{A}(v, w)=\sum_{k=1}^{N} \omega_{k} \int_{\Omega_{k}} \nabla v \cdot \nabla w d x, \quad v \in V_{h}(\Omega), \forall w \in V_{h}(\Omega)
$$

The equation (2.3) can be written as

$$
\begin{equation*}
A_{h} u_{h}=f_{h}, \quad u_{h} \in V_{h}(\Omega) \tag{2.4}
\end{equation*}
$$

The goal of this paper is to construct a substructuring preconditioner for $A_{h}$ based on the domain decomposition described above.

### 2.2 Notations

To introduce the new method, we need some more notations. Throughout this paper, a subset $G$ of $\Omega$ are always understood as an open set.

- subdomain spaces

For subdomain $\Omega_{k}$, define

$$
V_{h}\left(\Omega_{k}\right)=\left\{\left.v\right|_{\Omega_{k}}: \forall v \in V_{h}(\Omega)\right\}
$$

and

$$
V_{h}^{p}\left(\Omega_{k}\right)=\left\{v_{h} \in V_{h}(\Omega): \operatorname{supp} v_{h} \subset \Omega_{k}\right\}
$$

Set

$$
\Omega_{i j}=\Omega_{i} \cup \Gamma_{i j} \cup \Omega_{j}
$$

and define

$$
V_{h}^{p}\left(\Omega_{i j}\right)=\left\{v_{h} \in V_{h}(\Omega): \operatorname{supp} v_{h} \subset \Omega_{i j}\right\}
$$

- interface space and face spaces

As usual, we define the (global) interface space by

$$
W_{h}(\Gamma)=\left\{\left.v\right|_{\Gamma}: \forall v \in V_{h}(\Omega)\right\}
$$

For each $\partial \Omega_{k}$, set

$$
W_{h}\left(\partial \Omega_{k}\right)=\left\{\left.v\right|_{\partial \Omega_{k}}: \forall v \in W_{h}(\Gamma)\right\}
$$

For a face $\mathrm{F}=\Gamma_{i j}$, define

$$
\tilde{W}_{h}(\mathrm{~F})=\left\{\phi_{h} \in W_{h}(\Gamma): \operatorname{supp} \phi_{h} \subset \mathrm{~F}\right\} .
$$

- interpolation-type operator and constant-like basis

For a subset $G$ of $\Gamma$, define the interpolation-type operator $I_{G}^{0}: W_{h}(\Gamma) \rightarrow W_{h}(\Gamma)$ as

$$
\left(I_{G}^{0} \phi_{h}\right)(p)=\left\{\begin{array}{cc}
\phi_{h}(p), & \text { if } p \in \mathcal{N}_{h} \cap G \\
0, & \text { if } p \in \mathcal{N}_{h} \cap(\Gamma \backslash G)
\end{array}\right.
$$

In particular, we have

$$
\left(I_{G}^{0} 1\right)(p)= \begin{cases}1, & \text { if } p \in \mathcal{N}_{h} \cap G \\ 0, & \text { if } p \in \mathcal{N}_{h} \cap(\Gamma \backslash G)\end{cases}
$$

If $G$ is an edge or a face generated by the domain decomposition, we call $\phi_{G}=I_{G}^{0} 1$ to be constant-like basis function on $G$, which will be used repeatedly.

- integration average and algebraic average

For a function $\varphi_{h} \in W_{h}(\Gamma)$, let $\gamma_{G}\left(\varphi_{h}\right)$ denote the integration average of $\varphi_{h}$ on $G$, and let $\gamma_{h, G}(\varphi)$ denote the algebraic average of the values of $\varphi$ on the nodes in $G$.

- sets of faces, edges, vertices and subdomains

For convergence, let $\mathcal{F}_{\Gamma}$ denote the set of all the faces $\Gamma_{i j}$. Besides, let $\mathcal{E}_{\Gamma}$ and $\mathcal{V}_{\Gamma}$ denote the set of the interior edges and the set of interior vertices generated by the decomposition

$$
\bar{\Omega}=\bigcup \bar{\Omega}_{k}
$$

respectively. For an edge $\mathrm{E} \in \mathcal{E}_{\Gamma}$, let $\mathcal{Q}_{\mathrm{E}}$ denote the set of the indices $k$ of the subdomains $\Omega_{k}$ which contain E as an edge. Namely,

$$
\mathcal{Q}_{\mathrm{E}}=\left\{k: \mathrm{E} \subset \partial \Omega_{k}\right\} .
$$

Define

$$
\Omega_{\mathrm{E}}=\bigcup_{k \in \mathcal{Q}_{\mathrm{E}}} \Omega_{k}, \quad \mathrm{E} \in \mathcal{E}_{\Gamma} .
$$

- face inner-products, scaling norm and interface norm

For a subset $G$ of $\Gamma$, let $\langle\cdot, \cdot\rangle_{G}$ denote the $L^{2}$ inner product on $G$. In particular, the $\langle\cdot, \cdot\rangle_{\Gamma}$ is abbreviated as $\langle\cdot, \cdot \cdot\rangle$. Let $\|\cdot\|_{0, G}$ denote the norm induced from $\langle\cdot, \cdot\rangle_{G}$.

For a sub-faces $G$ of $\Gamma$, let $H_{G}$ denote the "size" of $G$. Define the scaled norm

$$
\|\phi\|_{\frac{1}{2}, G}=\left(|\phi|_{\frac{1}{2}, G}^{2}+H_{G}^{-1}\|\phi\|_{0, G}^{2}\right)^{\frac{1}{2}}, \quad \forall \phi \in H^{\frac{1}{2}}(G)
$$

For convenience, define

$$
\left\|\phi_{h}\right\|_{*, \Gamma}=\left(\sum_{k=1}^{\mathrm{N}} \omega_{k}\left|\phi_{h}\right|_{\frac{1}{2}, \partial \Omega_{k}}^{2}\right)^{\frac{1}{2}}, \quad \forall \phi_{h} \in W_{h}(\Gamma) .
$$

- discrete norms

Discrete norms (or semi-norms) of finite element functions will be used repeatedly in this paper, since the discrete norms are defined on a set of nodes only, and do not depend on the geometric shape of the underlying domain.

We first give definitions of two well known discrete norms (refer to [32]), which are equivalent to their respective continuous norms. For $v_{h} \in V_{h}\left(\Omega_{k}\right)$, the discrete $H^{1}$ seminorm is defined by

$$
\left|v_{h}\right|_{1, h, \Omega_{k}}^{2}=h^{3} \sum_{p_{i}, p_{j} \in \mathcal{N}_{h} \cap \Omega_{k}}\left|v_{h}\left(p_{i}\right)-v_{h}\left(p_{j}\right)\right|^{2},
$$

where $p_{i}$ and $p_{j}$ denote two neighboring nodes. Similarly, the discrete $L^{2}$ norm on an edge E of $\Omega_{k}$ is defined by

$$
\left\|v_{h}\right\|_{0, h, \mathrm{E}}^{2}=h \sum_{p \in \mathcal{N}_{h} \cap \mathrm{E}}\left|v_{h}(p)\right|^{2}
$$

- $H_{00}^{\frac{1}{2}}$ norms

For $\varphi_{h} \in \tilde{W}_{h}(\mathrm{~F})$ with $\mathrm{F}=\Gamma_{i j}$, define

$$
\left\|\varphi_{h}\right\|_{H_{00}^{\frac{1}{2}(\mathrm{~F})}}^{2}=\left|\varphi_{h}\right|_{\frac{1}{2}, \mathrm{~F}}^{2}+\int_{\mathrm{F}} \frac{\left|\varphi_{h}(x)\right|^{2}}{\operatorname{dist}(x, \partial \mathrm{~F})} d s(x) .
$$

Hereafter, $\operatorname{dist}(x, \partial \mathrm{~F})$ denotes the shortest distance from a point $x \in \mathrm{~F}$ to the boundary $\partial F$. It is known that

$$
\left\|\varphi_{h}\right\|_{H_{00}^{\frac{1}{2}}(\mathrm{~F})}^{2} \equiv\left|\tilde{\varphi}_{h}\right|_{\frac{1}{2}, \Omega_{i}}^{2} \equiv\left|\tilde{\varphi}_{h}\right|_{\frac{1}{2}}^{2}, \Omega_{j}
$$

where $\tilde{\varphi}_{h} \in W_{h}(\Gamma)$ denotes the zero extension of $\varphi_{h}$. Moreover, we have

$$
\left\|\varphi_{h}\right\|_{H_{0}^{2}(\mathrm{~F})}^{2}=\int_{\mathrm{F}}^{\frac{1}{2}} \frac{\left|\varphi_{h}(x)\right|^{2}}{\operatorname{dist}(x, \partial \mathrm{~F})} d s(x) .
$$

The corresponding discrete semi-norm is defined by

$$
\left\|\varphi_{h}\right\|_{h, H_{00}^{2}}^{2}=h^{2} \sum_{p \in \mathcal{\mathcal { N } _ { h } \cap \mathrm { F }}} \frac{\left|\varphi_{h}(p)\right|^{2}}{\operatorname{dist}(p, \partial \mathrm{~F})} .
$$

- spectrally equivalences

For simplicity, we will frequently use the notations $\lesssim$ and $\equiv$. For any two non-negative quantities $x$ and $y, x \lesssim y$ means that $x \leq C y$ for some constant $C$ independent of mesh size $h$, subdomain size $d$ and the related parameters. $x \equiv y$ means $x \lesssim y$ and $y \lesssim x$.

### 2.3 Motivation

We first recall the main ideas of the existing substucturing preconditioners.
Let $E_{k}: W_{h}\left(\partial \Omega_{k}\right) \rightarrow V_{h}\left(\Omega_{k}\right)$ be the discrete harmonic extension. Define the harmonic subspace

$$
V_{h}^{\Gamma}(\Omega)=\left\{v_{h} \in V_{h}(\Omega):\left.v_{h}\right|_{\Omega_{k}}=E_{k}\left(\phi_{h} \mid \partial \Omega_{k}\right)(k=1, \cdots, \mathrm{~N}) \text { for some } \phi_{h} \in W_{h}(\Gamma)\right\} .
$$

Then, we have the initial space decomposition

$$
V_{h}(\Omega)=\sum_{k=1}^{\mathrm{N}} V_{h}^{p}\left(\Omega_{k}\right)+V_{h}^{\Gamma}(\Omega) .
$$

Let $A_{h, k}: V_{h}^{p}\left(\Omega_{k}\right) \rightarrow V_{h}^{p}\left(\Omega_{k}\right)$ be the restriction of the operator $A_{h}$ on the local space $V_{h}^{p}\left(\Omega_{k}\right)$, and let $B_{h, \Gamma}: V_{h}^{\Gamma}(\Omega) \rightarrow V_{h}^{\Gamma}(\Omega)$ be a symmetric and positive definite operator which is spectrally equivalent to the restriction of the operator $A_{h}$ on the harmonic subspace $V_{h}^{\Gamma}(\Omega)$. Then, the classical substructuring preconditioner (refer to [5]) can be defined in the rough form

$$
\begin{equation*}
B_{\text {old }}^{-1}=\sum_{k=1}^{N} A_{h, k}^{-1} Q_{k}+B_{h, \Gamma}^{-1} Q_{\Gamma}, \tag{2.5}
\end{equation*}
$$

where $Q_{k}$ and $Q_{\Gamma}$ denote the standard $L^{2}$ projectors into their respective subspaces. For the preconditioner $B_{\text {old }}$, we have (see [5])

$$
\begin{equation*}
\operatorname{cond}\left(B_{\text {old }}^{-1} A_{h}\right) \lesssim \log ^{2}(H / h) . \tag{2.6}
\end{equation*}
$$

In many applications, the subspaces $V_{h}^{p}\left(\Omega_{k}\right)$ still have high dimensions, so it is expensive to use the exact solvers $A_{h, k}^{-1}$.

It was shown in [3] that substructuring preconditioners with inexact solvers $B_{h, k}^{-1}$ still possess nearly optimal convergence, if each $B_{h, k}$ has some spectrally approximation to $A_{h, k}$ (the usual spectrally equivalence is not enough). Hereafter, "inexact" means that $B_{h, k}$ is only spectrally equivalent to $A_{h, k}$, for example, $B_{h, k}$ is a multigrid preconditioner for $A_{h, k}$. It seems difficult to design an efficient substructuring preconditioner with completely inexact solvers $B_{h, k}^{-1}$, instead of $A_{h, k}^{-1}$ itself or its approximation. In essence, one has to modify the harmonic subspace $V_{h}^{\Gamma}(\Omega)$ by replacing each harmonic extension $E_{k}$ with another extension.

In [4], a substructuring preconditioner $B_{b p v}^{-1}$ with inexact solvers was been designed by replacing each harmonic extension $E_{k}$ with a simple average extension. It has been shown that the condition number of the resulting preconditioned system can be estimated by

$$
\begin{equation*}
\operatorname{cond}\left(B_{b p v}^{-1} A_{h}\right) \lesssim H / h \tag{2.7}
\end{equation*}
$$

In [11], another substructuring preconditioner $B_{h l m}^{-1}$ with inexact solvers was been designed by replacing each harmonic extension $E_{k}$ with an approximate harmonic extension. If the approximate harmonic extension is exact enough, then

$$
\begin{equation*}
\operatorname{cond}\left(B_{b l m}^{-1} A_{h}\right) \lesssim \log ^{2}(H / h) \tag{2.8}
\end{equation*}
$$

The approximate harmonic extension can be defined by approximate harmonic basis functions. It is not practical to compute all the approximate harmonic basis functions. Because of this, an alternative method, which still require high accuracy of $B_{h, k}$, was considered in [11].

Another way to construct substructuring preconditioner with inexact solvers was considered in [6] (mainly for two dimensions).

Let $\hat{W}_{h}^{0}(\Gamma)$ be a subspace of $W_{h}(\Gamma)$, such that a function $\varphi \in \hat{W}_{h}^{0}(\Gamma)$ equals a constant $C_{\mathrm{F}}$ at every nodes in each face F of $\Gamma$. Define the coarse subspace

$$
\hat{V}_{h}^{0}(\Omega)=\left\{v_{h} \in V_{h}(\Omega):\left.v_{h}\right|_{\Omega_{k}}=E_{k}\left(\left.\phi_{h}\right|_{\partial \Omega_{k}}\right)(k=1, \cdots, \mathrm{~N}) \text { for some } \phi_{h} \in \hat{W}_{h}^{0}(\Gamma)\right\} .
$$

Then, we have the space decomposition

$$
V_{h}(\Omega)=\hat{V}_{h}^{0}(\Omega)+\sum_{\Gamma_{i j}} V_{h}^{p}\left(\Omega_{i j}\right)
$$

Let $B_{0}: \hat{V}_{h}^{0}(\Omega) \rightarrow \hat{V}_{h}^{0}(\Omega)$ and $B_{i j}: V_{h}^{p}\left(\Omega_{i j}\right) \rightarrow V_{h}^{p}\left(\Omega_{i j}\right)$ be symmetric and positive definite operators which are spectrally equivalent to the restrictions of the operator $A_{h}$ on the coarse subspace $\hat{V}_{h}^{0}(\Omega)$ and the local subspace $V_{h}^{p}\left(\Omega_{i j}\right)$, respectively. Then, we can define a substructuring preconditioner as (refer to [6] and [32])

$$
\begin{equation*}
B_{d w}^{-1}=B_{0}^{-1} Q_{0}+\sum_{\Gamma_{i j}} B_{i j}^{-1} Q_{i j} \tag{2.9}
\end{equation*}
$$

where $Q_{0}$ and $Q_{i j}$ denote the standard $L^{2}$ projectors into their respective subspaces. For the preconditioner $B_{d w}$, we have (refer to [6] and [32])

$$
\begin{equation*}
\operatorname{cond}\left(B_{d w}^{-1} A_{h}\right) \lesssim \log ^{2}(H / h) \tag{2.10}
\end{equation*}
$$

As pointed out in [32], the advantage of the preconditioner $B_{d w}$ is that one can use inexact solver $B_{i j}$ for the restriction of $A_{h}$ on $V_{h}^{p}\left(\Omega_{i j}\right)$, but an inexact solver $B_{0}^{-1}$ is hard to come by as harmonic extensions in the definition of $\hat{V}_{h}^{0}(\Omega)$ mean an exact solver. As an alternate method, one can use a nearly harmonic extension to replace the exact harmonic extension $E_{k}$. However, the design of such approximate harmonic extension is also difficult
for the general case that either the subdomain $\Omega_{k}$ is a general polyhedron or $\mathcal{T}_{h}$ is a general quasi-uniform triangulation (an attempt for two dimension is given in [19]). From the definition of the coarse subspace $\hat{V}_{h}^{0}(\Omega)$, we know that one needs only to design such extension for functions in the subspace $\hat{W}_{h}^{0}(\Gamma)$. In essence, one needs only to define a suitable extension for the constant-like basis function on each face of $\Gamma$. This kind of extension was first studied in [8]. For the purpose of applications, the definition of such extension would be sufficiently simple, otherwise, the action of $B_{0}^{-1}$ is difficult to implement.

In this paper, we construct a completely explicit extension for the constant-like basis function on each face of $\Gamma$. As we will see, the implementation of the new extensions is very convenient and cheap. We can define a new variant of the coarse subspace $\hat{V}_{h}^{0}(\Omega)$ based on the extensions. Then, we use the new coarse subspace to construct a substructuring preconditioner with inexact solvers.

### 2.4 Basic tools

In the analysis later, we will use some basic estimates on $H^{\frac{1}{2}}$ norms.
The following results are well-known (see, for example, [32]).
Lemma 2.1 Let E and F be an edge and a face of $\Omega_{k}$. Then,

$$
\begin{gather*}
\left\|\phi_{h}\right\|_{0, \partial \mathrm{~F}} \lesssim \log ^{\frac{1}{2}}(H / h)\left\|\phi_{h}\right\|_{\frac{1}{2}, \partial \Omega_{k}}, \quad \forall \phi_{h} \in W_{h}\left(\partial \Omega_{k}\right),  \tag{2.11}\\
\left\|I_{\mathrm{E}}^{0} \phi_{h}\right\|_{\frac{1}{2}, \partial \Omega_{k}} \lesssim \log ^{\frac{1}{2}}(H / h)\left\|\phi_{h}\right\|_{\frac{1}{2}, \partial \Omega_{k}}, \quad \forall \phi_{h} \in W_{h}\left(\partial \Omega_{k}\right) \tag{2.12}
\end{gather*}
$$

and

$$
\begin{equation*}
\left\|I_{\mathrm{F}}^{0} \phi_{h}\right\|_{H_{00}^{\frac{1}{2}}(\mathrm{~F})} \lesssim \log (d / h)\left\|\phi_{h}\right\|_{\frac{1}{2}, \partial \Omega_{k}}, \quad \forall \phi_{h} \in W_{h}\left(\partial \Omega_{k}\right) \tag{2.13}
\end{equation*}
$$

The following result can be proved as in Lemma 6.2 in [18], together with the standard technique.

Lemma 2.2 Let E be an edge of $\Omega_{k}$. Then,

$$
\begin{equation*}
\left\|\phi_{h}-\gamma_{h, \mathrm{E}}\left(\phi_{h}\right)\right\|_{\frac{1}{2}, \partial \Omega_{k}} \lesssim \log ^{\frac{1}{2}}(H / h)\left|\phi_{h}\right|_{\frac{1}{2}, \partial \Omega_{k}}, \quad \forall \phi_{h} \in W_{h}\left(\partial \Omega_{k}\right) \tag{2.14}
\end{equation*}
$$

## 3 Extensions for particular functions

As in Subsection 2.3, the desired coarse subspace of $V_{h}(\Omega)$ is always associated with an interface coarse subspace. Let $I_{\mathrm{F}}^{0} 1$ and $I_{\mathrm{E}}^{0} 1$ denote the constant-like basis functions on the face F and the coarse edge E, respectively (see Subsection 2.2). In this paper, we consider the interface coarse subspace

$$
W_{h}^{0}(\Gamma)=\operatorname{span}\left\{I_{F}^{0} 1, I_{\mathrm{E}}^{0} 1, \varphi_{p}: F \in \mathcal{F}_{\Gamma}, \mathrm{E} \in \mathcal{E}_{\Gamma}, p \in \mathcal{V}_{\Gamma}\right\}
$$

Hereafter, $\varphi_{p}$ denotes the nodal basis function on the node $p \in \mathcal{N}_{h} \cap \Omega$.
In this section, we construct an approximation for the restriction of the discrete harmonic extension $E_{k}$ on $W_{h}^{0}(\Gamma)$. For convenience, let F denote a general face $\Gamma_{i j}$, and let $\phi_{\mathrm{F}}=I_{\mathrm{F}}^{0} 1 \in \tilde{W}_{h}(F)$ be the constant-like basis function on F .

### 3.1 An explicit extension of $\phi_{F}$

In this subsection, we define a stable extension $E_{\mathrm{F}} \phi_{\mathrm{F}} \in V_{h}^{p}\left(\Omega_{i j}\right)$, which satisfies $\left(E_{\mathrm{F}} \phi_{\mathrm{F}}\right)(p)=$ $\phi_{\mathrm{F}}(p)$ when $p \in \overline{\mathrm{~F}}$.

For a node $p$ in $\Omega_{i j} \backslash F$, let $p^{\prime}$ denote the projection of $p$ on the plane containing $F$. Besides, when $p^{\prime} \in \mathrm{F}$, we use $p^{\prime \prime}$ to denote a point on the boundary $\partial \mathrm{F}$, such that $\left|p^{\prime} p^{\prime \prime}\right|=\operatorname{dist}\left(p^{\prime}, \partial \mathrm{F}\right)$, which is the shortest distance from $p^{\prime}$ to the boundary $\partial \mathrm{F}$ (see Figure 1).


Figure 1: illustration for the notations
In short words, we define the extension $E_{\mathrm{F}} \phi_{\mathrm{F}}$ such that the values $\left(E_{\mathrm{F}} \phi_{\mathrm{F}}\right)(p)$ decrease gradually when the lengths $\left|p p^{\prime}\right|$ increase, or the lengths $\left|p^{\prime} p^{\prime \prime}\right|$ decrease. Although this property holds also for the extension designed in [8], we will use a different idea from [8].

For simplicity of exposition, we make an assumption:
Assumption 3.1: when $p \in \partial \Omega_{i j}$ with $p^{\prime} \in \mathrm{F},\left|p p^{\prime}\right| \geq\left|p^{\prime} p^{\prime \prime}\right|$.
This assumption means that each subdomain $\Omega_{k}$ is not too thin. As we will see in Appendix, we need only to revise slightly the definition of the extension for the general situation without this assumption. To give the exact definition of $E_{\mathrm{F}} \phi_{\mathrm{F}}$ with $\mathrm{F}=\Gamma_{i j}$, we set

$$
\Lambda_{\mathrm{F}}=\left\{p \in \mathcal{N}_{h} \cap\left(\Omega_{i j} \backslash \mathrm{~F}\right): p^{\prime} \in \mathrm{F},\left|p p^{\prime}\right| \leq\left|p^{\prime} p^{\prime \prime}\right|\right\}
$$

For a face $\mathrm{F}=\Gamma_{i j}$, define the extension $E_{\mathrm{F}} \phi_{\mathrm{F}} \in V_{h}^{p}\left(\Omega_{i j}\right)$ as follows:

$$
\left(E_{\mathrm{F}} \phi_{\mathrm{F}}\right)(p)=\left\{\begin{array}{cc}
1, \quad \text { if } p \in \mathrm{~F},  \tag{3.1}\\
1-\frac{\left|p p^{\prime}\right|}{\left|p^{\prime} p^{\prime \prime}\right|}, & \text { if } p \in \Lambda_{\mathrm{F}}, \\
0, & \text { otherwise } .
\end{array}\right.
$$

In particular, we have $\left(E_{\mathrm{F}} \phi_{\mathrm{F}}\right)(p)=\phi_{\mathrm{F}}(p)$ when $p \in \overline{\mathrm{~F}}$.
For rectangular face F with uniform triangulation, the values of $E_{\mathrm{F}} \phi_{\mathrm{F}}$ at some nodes are given in Figure 2.


Figure 2: extension of $\phi_{F}$ for particular case

Set

$$
d_{\mathrm{F}}=\max _{q \in \mathrm{~F}} \operatorname{dist}(q, \partial \mathrm{~F})
$$

It is clear that $d_{\mathrm{F}}$ is just the radium of the largest circle contained in F , and $d_{\mathrm{F}} \equiv H$.
It is easy to see that the extension operator $E_{\mathrm{F}}$ possesses the properties:
(i) the support set $\Omega_{\mathrm{F}}$ of $E_{\mathrm{F}} \phi_{\mathrm{F}}$ is a simply connected domain with the size $d_{\mathrm{F}}$. In fact, $\Omega_{\mathrm{F}}$ like a cone with the bottom F .
(ii) the calculation of $E_{\mathrm{F}} \phi_{\mathrm{F}}$ possesses the optimal complexity $O\left(n_{\mathrm{F}}\right)$ with $n_{\mathrm{F}}=\left(d_{\mathrm{F}} / h\right)^{3}$ being the number of the nodes in $\Omega_{\mathrm{F}}$.

Besides, the extension $E_{\mathrm{F}}$ is stable in the following sense
Theorem 3.1 The extension $E_{\mathrm{F}}$ defined by (3.15) satisfies the stability condition

$$
\begin{equation*}
\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2},\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{j}}^{2} \lesssim H \log (H / h) \equiv\left|\phi_{\mathrm{F}}\right|_{H_{00}}^{2}(\mathrm{~F}) \quad\left(\mathrm{F}=\Gamma_{i j}\right) . \tag{3.16}
\end{equation*}
$$

Since proof of the result is technical, we prove this theorem in Section 5.
Remark 3.1 The definition of the extension $E_{\mathrm{F}}$ not only is simple and completely explicit, but also is identical to different geometric shape (tetrahedron, hexahedron or other polyhedrons) of the two subdomains containing F as a face. Thus, the action of $E_{\mathrm{F}}$ is easy to implement.

### 3.2 A nearly harmonic extension of $\phi \in W_{h}^{0}(\Gamma)$

Let $R_{k}^{0}: W_{h}\left(\partial \Omega_{k}\right) \rightarrow V_{h}\left(\Omega_{k}\right)$ denote the zero extension operator in the sense that $R_{k}^{0} \phi=\phi$ on $\partial \Omega_{k}$, and $R_{k}^{0} \phi$ vanishes at all internal nodes of $\Omega_{k}$ for $\phi \in W_{h}\left(\partial \Omega_{k}\right)$. For $\psi \in W_{h}^{0}(\Gamma)$, we have

$$
\psi=\gamma_{h, \partial \Omega_{k}}(\psi)+I_{\mathcal{W}_{k}}^{0}\left(\psi-\gamma_{h, \partial \Omega_{k}}(\psi)\right)+\sum_{\mathrm{F} \subset \partial \Omega_{k}} \gamma_{h, \mathrm{~F}}\left(\psi-\gamma_{h, \partial \Omega_{k}}(\psi)\right) \phi_{\mathrm{F}}, \quad \text { on } \quad \partial \Omega_{k}
$$

Hereafter, $\mathcal{W}_{k}$ denotes the wire-basket set of $\Omega_{k}$, i.e., the union of all the edges of $\Omega_{k}$. Note that $\gamma_{h, \mathrm{~F}}\left(\psi-\gamma_{h, \partial \Omega_{k}}(\psi)\right)$ is just the (constant) value of $\psi-\gamma_{h, \partial \Omega_{k}}(\psi)$ at the interior nodes of F . Then, we define on each $\Omega_{k}$

$$
\begin{equation*}
E_{0} \psi=\gamma_{h, \partial \Omega_{k}}(\psi)+R_{k}^{0}\left[I_{\mathcal{W}_{k}}^{0}\left(\psi-\gamma_{h, \partial \Omega_{k}}(\psi)\right)\right]+\sum_{\mathrm{F} \subset \partial \Omega_{k}} \gamma_{h, \mathrm{~F}}\left(\psi-\gamma_{h, \partial \Omega_{k}}(\psi)\right) E_{\mathrm{F}} \phi_{\mathrm{F}} . \tag{3.17}
\end{equation*}
$$

It is easy to see that $E_{0} \psi=\psi$ on $\Gamma$.
The extension $E_{0}$ is nearly harmonic in the following sense
Proposition 3.1 The global coarse extension $E_{0}$ satisfies

$$
\begin{equation*}
\sum_{k=1}^{\mathrm{N}} \omega_{k}\left|E_{0} \phi_{0}\right|_{1, \Omega_{k}}^{2} \lesssim \log (H / h)\left\|\phi_{0}\right\|_{*, \Gamma}^{2}, \quad \forall \phi_{0} \in W_{h}^{0}(\Gamma) . \tag{3.18}
\end{equation*}
$$

Proof. By the definition of $E_{0}$, we have

$$
\begin{aligned}
\left|E_{0} \phi_{0}\right|_{1, \Omega_{k}}^{2} & \lesssim\left|R_{k}^{0}\left[I_{\mathcal{W}_{k}}^{0}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right)\right]\right|_{1, \Omega_{k}}^{2} \\
& +\sum_{\mathrm{F} \subset \partial \Omega_{k}}\left|\gamma_{h, \mathrm{~F}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right)\right|^{2} \cdot\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{k}}^{2} \\
& \lesssim\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \mathcal{W}_{k}}^{2}
\end{aligned}
$$

$$
\begin{equation*}
+\sum_{\mathrm{F} \subset \partial \Omega_{k}}\left(H^{-2}\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \partial \Omega_{k}}^{2} \cdot\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{k}}^{2}\right) . \tag{3.19}
\end{equation*}
$$

Using (2.11) and Friedrichs' inequality, yields

$$
\begin{equation*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \mathcal{W}_{k}}^{2} \lesssim \log (H / h)\left|\phi_{0}\right|_{\frac{1}{2}, \partial \Omega_{k}}^{2} . \tag{3.20}
\end{equation*}
$$

On the other hand, we deduce by (3.16) and Friedrichs' inequality

$$
\begin{aligned}
H^{-2}\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \partial \Omega_{k}}^{2} \cdot\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{k}}^{2} & \lesssim H^{-1} \log (H / h)\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \partial \Omega_{k}}^{2} \\
& \lesssim \log (H / h)\left|\phi_{0}\right|_{\frac{1}{2}, \partial \Omega_{k}}^{2} .
\end{aligned}
$$

Substituting (3.20) and the above inequality into (3.19), we get (3.18).

## 4 A substructuring method with inexact solvers

This section is devoted to construction of a substructuring preconditioner with inexact solvers. The new preconditioner is based on the extension designed in the last section.

### 4.1 Space decomposition for $V_{h}(\Omega)$

In this subsection, we define a space decomposition of $V_{h}(\Omega)$.
For each subdomain $\Omega_{i j}$ containing $\Gamma_{i j}$, let $V_{h}^{p}\left(\Omega_{i j}\right)$ be the subspace given in Subsection 2.2. Namely,

$$
V_{h}^{p}\left(\Omega_{i j}\right)=\operatorname{span}\left\{\varphi_{p}: p \in \mathcal{N}_{h} \cap \Omega_{i j}\right\} .
$$

For the extension $E_{0}: W_{h}^{0}(\Gamma) \rightarrow V_{h}(\Omega)$ described in the last section, define

$$
V_{h}^{0}(\Omega)=\left\{v_{h} \in V_{h}(\Omega): v_{h}=E_{0} \phi \text { for some } \phi_{h} \in W_{h}^{0}(\Gamma)\right\} .
$$

For $\mathrm{E} \in \mathcal{E}_{\Gamma}$, define

$$
V_{h}\left(\Omega_{\mathrm{E}}\right)=\operatorname{span}\left\{\varphi_{p}: p \in \mathrm{E} \cap \mathcal{N}_{h}\right\} .
$$

Then, we have the space decomposition

$$
V_{h}(\Omega)=V_{h}^{0}(\Omega)+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} V_{h}\left(\Omega_{\mathrm{E}}\right)+\sum_{\Gamma_{i j}} V_{h}^{p}\left(\Omega_{i j}\right) .
$$

Remark 4.1 Since the new extension $E_{0}$ described in the last section is used in $V_{h}^{0}(\Omega)$, the coarse subspace $V_{h}^{0}(\Omega)$ is different from the one in [8].

### 4.2 A substructuring preconditioner

Based on the space decomposition in the last subsection, we can define a preconditioner in the standard way.

Define symmetric and positive definite operators as follows:

- the global coarse solver $B_{0}: V_{h}^{0}(\Omega) \rightarrow V_{h}^{0}(\Omega)$ satisfies:

$$
\begin{equation*}
\left(B_{0} v_{h}, v_{h}\right) \equiv\left(A_{h} v_{h}, v_{h}\right), \quad \forall v_{h} \in V_{h}^{0}(\Omega) ; \tag{4.1}
\end{equation*}
$$

- the edge solver $B_{\mathrm{E}}: V_{h}\left(\Omega_{\mathrm{E}}\right) \rightarrow V_{h}\left(\Omega_{\mathrm{E}}\right)$ satisfies:

$$
\begin{equation*}
\left(B_{\mathrm{E}} v_{h}, v_{h}\right) \equiv\left(A_{h} v_{h}, v_{h}\right), \quad \forall v_{h} \in V_{h}\left(\Omega_{\mathrm{E}}\right) ; \tag{4.2}
\end{equation*}
$$

- the interface solver $B_{i j}: V_{h}^{p}\left(\Omega_{i j}\right) \rightarrow V_{h}^{p}\left(\Omega_{i j}\right)$ satisfies:

$$
\begin{equation*}
\left(B_{i j} v_{h}, v_{h}\right) \equiv\left(A_{h} v_{h}, v_{h}\right), \quad \forall v_{h} \in V_{h}^{p}\left(\Omega_{i j}\right) . \tag{4.3}
\end{equation*}
$$

In applications, the interface solver $B_{i j}$ is usually chosen as a symmetric multigrid preconditioner for the restriction of $A_{h}$ on $V_{h}^{p}\left(\Omega_{i j}\right)$. Since all the subspaces $V_{h}^{0}(\Omega)$ and $V_{h}\left(\Omega_{\mathrm{E}}\right)$ have low dimensions, the solvers $B_{0}$ and $B_{\mathrm{E}}$ can be simply defined as the restriction operators of $A_{h}$ on their respective subspaces.

Now, the desired domain decomposition preconditioner for $A_{h}$ is defined as

$$
\begin{equation*}
B_{h}^{-1}=B_{0}^{-1} Q_{0}+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} B_{\mathrm{E}}^{-1} Q_{\mathrm{E}}+\sum_{\Gamma_{i j}} B_{i j}^{-1} Q_{i j} \tag{4.4}
\end{equation*}
$$

where $Q_{0}: V_{h}(\Omega) \rightarrow V_{h}^{0}(\Omega), Q_{\mathrm{E}}: V_{h}(\Omega) \rightarrow V_{h}\left(\Omega_{\mathrm{E}}\right)$ and $Q_{i j}: V_{h}(\Omega) \rightarrow V_{h}^{p}\left(\Omega_{i j}\right)$ denote $L^{2}$ projectors.

### 4.3 Implementation

Before describing our algorithms, we give exact definitions of the edge solver $B_{\mathrm{E}}$ and the coarse solver $B_{0}$.

- the edge solver

It is known that

$$
|\nabla v|_{0, \Omega_{k}} \equiv\|v\|_{0, \mathrm{E}}, \quad \forall v \in V_{h}\left(\Omega_{\mathrm{E}}\right) \quad\left(\mathrm{E} \subset \partial \Omega_{k}\right) .
$$

Then,

$$
\left(A_{h} v, v\right) \equiv \sum_{k \in \mathcal{Q}_{\mathrm{E}}} \omega_{k}\|v\|_{0, \mathrm{E}}^{2}, \quad \forall v \in V_{h}\left(\Omega_{\mathrm{E}}\right) .
$$

We can define $B_{\mathrm{E}}: V_{h}\left(\Omega_{\mathrm{E}}\right) \rightarrow V_{h}\left(\Omega_{\mathrm{E}}\right)$ by

$$
\left(B_{\mathrm{E}} v, w\right)=\left(\sum_{k \in \mathcal{Q}_{\mathrm{E}}} \omega_{k}\right)\langle v, w\rangle_{\mathrm{E}}, \quad \forall w \in V_{h}\left(\Omega_{\mathrm{E}}\right) .
$$

The stiffness of $B_{\mathrm{E}}$ is just a scaling of the mass matrix associated with E .

- the coarse solver

As we will see in Section 5, we have

$$
\left|E_{0} \phi_{\mathrm{F}}\right|_{1, \Omega_{k}}^{2} \equiv H \log (H / h) .
$$

As in Proposition 3.1, one can verify that

$$
\left|v_{h}^{0}\right|_{1, \Omega_{k}}^{2} \equiv\left\|v_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(v_{h}^{0}\right)\right\|_{0, \mathcal{W}_{k}}^{2}+H \log (H / h) \sum_{\mathrm{F} \subset \partial \Omega_{k}}\left|\gamma_{h, \mathrm{~F}}\left(v_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(v_{h}^{0}\right)\right)\right|^{2}, \quad \forall v_{h}^{0} \in V_{h}^{0}(\Omega) .
$$

Hence, we have for any $v_{h}^{0} \in V_{h}^{0}(\Omega)$

$$
\left(A_{h} v_{h}^{0}, v_{h}^{0}\right) \equiv \sum_{k=1}^{N} \omega_{k}\left\{\left\|v_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(v_{h}^{0}\right)\right\|_{0, \mathcal{w}_{k}}^{2}+H \log (H / h) \sum_{\mathrm{F} \subset \partial \Omega_{k}}\left|\gamma_{h, \mathrm{~F}}\left(v_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(v_{h}^{0}\right)\right)\right|^{2}\right\} .
$$

This means that $B_{0}$ can be defined as

$$
\begin{gathered}
\left(B_{0} v_{h}^{0}, w_{h}^{0}\right)=\sum_{k=1}^{N} \omega_{k}\left\{\left\langlev_{h}^{0}-\gamma_{\left.h, \partial \Omega_{k}\left(v_{h}^{0}\right), w_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(w_{h}^{0}\right)\right\rangle_{h, \mathcal{W}_{k}}}^{\left.+H \log (H / h) \sum_{\mathrm{F} \subset \partial \Omega_{k}} \gamma_{h, \mathrm{~F}}\left(v_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(v_{h}^{0}\right)\right) \cdot \gamma_{h, \mathrm{~F}}\left(w_{h}^{0}-\gamma_{h, \partial \Omega_{k}}\left(w_{h}^{0}\right)\right)\right\},}\right.\right. \\
v_{h}^{0} \in V_{h}^{0}(\Omega), \forall w_{h}^{0} \in V_{h}^{0}(\Omega)
\end{gathered}
$$

Then, the stiffness matrix of $B_{0}$ is calculated by $\mathbf{B}_{0}=\left(b_{r l}\right)_{L_{0} \times L_{0}}$ with

$$
\begin{gathered}
b_{r l}=\left(B_{0}\left(E_{0} \psi_{r}^{0}\right), E_{0} \psi_{l}^{0}\right)=\sum_{k=1}^{N} \omega_{k}\left\{\left\langle\psi_{r}^{0}-\gamma_{h, \partial \Omega_{k}}\left(\psi_{r}^{0}\right), \psi_{l}^{0}-\gamma_{h, \partial \Omega_{k}}\left(\psi_{l}^{0}\right)\right\rangle_{h, \mathcal{W}_{k}}\right. \\
\left.+H \log (H / h) \sum_{\mathrm{F} \subset \partial \Omega_{k}} \gamma_{h, \mathrm{~F}}\left(\psi_{r}^{0}-\gamma_{h, \partial \Omega_{k}}\left(\psi_{r}^{0}\right)\right) \cdot \gamma_{h, \mathrm{~F}}\left(\psi_{l}^{0}-\gamma_{h, \partial \Omega_{k}}\left(\psi_{l}^{0}\right)\right)\right\} \\
\left(r, l=1, \cdots, L_{0}\right)
\end{gathered}
$$

Remark 4.2 Note that our coarse solver $B_{0}$ is different from the one described in Algorithm 6.10 of [8]. The coarse solver in [8] involves $N$ extra unknowns (see also [5]), which need to be solved by a special technique (see [32]). The coarse solver $B_{0}$ is similar with the one given in Subsection 4.3 of [7], in which stability of the coarse solver was not proved (We do not know why the article [7] cited directly the stability result of the coarse solver in [8]).

The action of $B_{h}^{-1}$ can be described by the following algorithm
Algorithm 4.1. For $g \in V_{h}(\Omega)$, the solution $u_{g} \in V_{h}(\Omega)$ satisfying

$$
\left(B_{h} u_{g}, v_{h}\right)=\left(g, v_{h}\right), \quad \forall v_{h} \in V_{h}(\Omega)
$$

can be gotten as follows:
Step 1. Computing $u_{0} \in V_{h}^{0}(\Omega)$ by

$$
\left(B_{0} u_{0}, v_{h}\right)=\left(g, v_{h}\right), \quad \forall v_{h} \in V_{h}^{0}(\Omega) ;
$$

Step 2. Computing $u_{\mathrm{E}} \in V_{h}\left(\Omega_{\mathrm{E}}\right)$ in parallel by

$$
\left(B_{\mathrm{E}} u_{\mathrm{E}}, v_{h}\right)=\left(g, v_{h}\right), \quad \forall v_{h} \in V_{h}\left(\Omega_{\mathrm{E}}\right) ;
$$

Step 3. Computing $u_{i j}^{p} \in V_{h}^{p}\left(\Omega_{i j}\right)$ in parallel by

$$
\left(B_{i j} u_{i j}^{p}, v_{h}\right)=\left(g, v_{h}\right), \quad \forall v_{h} \in V_{h}^{p}\left(\Omega_{i j}\right) ;
$$

Step 4. Set

$$
u_{g}=u_{0}+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} u_{\mathrm{E}}+\sum_{\Gamma_{i j}} u_{i j}^{p} .
$$

Remark 4.3 The efficiency of the preconditioner $B_{h}$ strongly depends on the action of the coarse solver $B_{0}$. The face basis extensions $E_{\mathrm{F}} \phi_{\mathrm{F}}$ play a key role in the definition of $B_{0}$. This is why we focus on the design of the extension $E_{\mathrm{F}}$ in this paper.

In the rest of this subsection, we consider an acceleration of Algorithm 4.1 (see [16] for the detailed discussion on this kind of acceleration). For convenience, set $V=V_{h}(\Omega)$, and define

$$
V_{1}=V_{h}^{0}(\Omega)+\sum_{\Gamma_{i j}} V_{h}^{p}\left(\Omega_{i j}\right) \quad \text { and } \quad V_{2}=\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} V_{h}\left(\Omega_{\mathrm{E}}\right) .
$$

Let $V_{2}^{\perp}$ denote the orthogonal complement of $V_{2}$ with respect to the inner product $\left(A_{h} \cdot, \cdot\right)$. It is clear that

$$
V=V_{1}+V_{2} .
$$

Let $A_{\mathrm{E}}: V_{h}\left(\Omega_{\mathrm{E}}\right) \rightarrow V_{h}\left(\Omega_{\mathrm{E}}\right)$ denote the restriction of $A_{h}$ on $V_{h}\left(\Omega_{\mathrm{E}}\right)$. Define

$$
B_{1}^{-1}=B_{0}^{-1} Q_{0}+\sum_{\Gamma_{i j}} B_{i j}^{-1} Q_{i j}
$$

The following algorithm can be viewed as a combination between the CG method and the multiplicative Schwarz method with the above decomposition for solving (2.4).
Algorithm 4.2 (multiplicative Schwarz-CG) Let $u_{1} \in V$ be an initial guess such that the error $u_{h}-u_{1} \in V_{2}^{\perp}$ (for example, $u_{1}$ can be chosen as $u_{1}=\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} A_{\mathrm{E}}^{-1} Q_{\mathrm{E}} f_{h} \in V_{2}$ ). When an approximation $u_{n} \in V$ has been gotten, we look for $u_{n+1} \in V$ as follows $(n \geq 1)$ :
Step 1. Solve $\varepsilon_{n 1} \in V_{1}$ by

$$
\varepsilon_{n 1}=B_{1}^{-1}\left(f_{h}-A_{h} u_{n}\right)
$$

If $\varepsilon_{n 1}=0$, then the iteration is terminated; otherwise, goto Step 2:
Step 2. Compute $u_{\mathrm{E}} \in V_{h}\left(\Omega_{\mathrm{E}}\right)$ in parallel by

$$
\left(A_{\mathrm{E}} u_{\mathrm{E}}, v\right)=\left(f_{h}-A_{h}\left(u_{n}+\varepsilon_{n 1}\right), v\right), \quad \forall v \in V_{h}\left(\Omega_{\mathrm{E}}\right) .
$$

Define $\varepsilon_{n 2} \in V_{2}$ by

$$
\varepsilon_{n 2}=\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} u_{\mathrm{E}}
$$

and set

$$
\varepsilon_{n}=\varepsilon_{n 1}+\varepsilon_{n 2} .
$$

Step 3. Compute

$$
\alpha_{n}=\varepsilon_{n}-\frac{\left(\varepsilon_{n}, A_{h} \alpha_{n-1}\right)}{\left\|\alpha_{n-1}\right\|_{A_{h}}^{2}} \alpha_{n-1}, \quad\left(\alpha_{0}=0\right)
$$

and

$$
u_{n+1}=u_{n}+\frac{\left(f_{h}-A_{h} u_{n}, \alpha_{n}\right)}{\left\|\alpha_{n}\right\|_{A_{h}}^{2}} \alpha_{n}
$$

Remark 4.4 Step 1 in Algorithm 4.2 would be implemented by Step 1 and Step 3 in Algorithm 4.1. Since each $V_{h}\left(\Omega_{\mathrm{E}}\right)$ has a very low dimension, which equals the number of nodes in the coarse edge E , Step 2 in Algorithm 4.2 is also very cheap.

### 4.4 Convergence

The following result gives an estimate of $\operatorname{cond}\left(B_{h}^{-1} A_{h}\right)$.
Theorem 4.1 Let the extension be defined by Subsection 3.1. For the preconditioner $B_{h}$ defined in Subsection 4.3, we have

$$
\begin{equation*}
\operatorname{cond}\left(B_{h}^{-1} A_{h}\right) \leq C \log ^{2}(H / h) \tag{4.5}
\end{equation*}
$$

where $C$ is a constant independent of $h, H$ and the jumps of the coefficient $\omega$ across the faces $\Gamma_{i j}$.

The convergence of Algorithm 4.2 involves the subspace $V_{2}^{\perp}$. Define $T^{*}=\left.\left(B_{1}^{-1} A_{h}\right)\right|_{V_{2}}$. It is clear that $T^{*}$ is symmetric and positive definite with respect to the inner product $\left(A_{h} \cdot, \cdot\right)$.

Theorem $4.2{ }^{[16]}$ Let the sequence $u_{n+1}$ be defined by Algorithm 4.2. Then

$$
\begin{equation*}
\left\|u_{h}-u_{n+1}\right\|_{A_{h}}<2\left(\frac{\sqrt{\kappa\left(T^{*}\right)}-1}{\sqrt{\kappa\left(T^{*}\right)}+1}\right)^{n}\left\|u_{h}-u_{1}\right\|_{A_{h}} \tag{4.6}
\end{equation*}
$$

where $\kappa\left(T^{*}\right)$ denotes the condition number of the operator $T^{*}$. Moreover, we have $\kappa\left(T^{*}\right)<$ $\operatorname{cond}\left(B_{h}^{-1} A_{h}\right)$.

Remark 4.5 Algorithm 4.2 is as cheap as PCG method for solving (2.4) with the preconditioner $B_{h}$. However, Theorem 4.2 indicates that the former has a faster convergence speed than the later.

Before proving Theorem 4.1, we need to derive a new estimate of the extension $E_{0}$. For $\phi \in W_{h}(\Gamma)$, define $\phi_{0} \in W_{h}^{0}(\Gamma)$ by

$$
\begin{equation*}
\phi_{0}=\sum_{\mathrm{F} \in \mathcal{F}_{\Gamma}} \gamma_{h, \mathrm{~F}}(\phi) I_{\mathrm{F}}^{0} 1+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} \gamma_{h, \mathrm{E}}(\phi) I_{\mathrm{E}}^{0} 1+\sum_{p \in \mathcal{V}_{\Gamma}} \phi(p) \varphi_{p} . \tag{4.7}
\end{equation*}
$$

Lemma 4.1 Let $\phi \in W_{h}(\Gamma)$, and let $\phi_{0}$ be defined by (4.7). Then,

$$
\begin{equation*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2} \lesssim\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2} \tag{4.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right|^{2} \lesssim H^{-2}\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2} \tag{4.9}
\end{equation*}
$$

Proof. Since

$$
1=\sum_{\mathrm{F} \subset \partial \Omega_{k}} I_{\mathrm{F}}^{0} 1+\sum_{\mathrm{E} \subset \partial \Omega_{k}} I_{\mathrm{E}}^{0} 1+\sum_{p \in \partial \Omega_{k} \cap \mathcal{V}_{\Gamma}} \phi_{p} \quad \text { on } \partial \Omega_{k},
$$

we have

$$
\gamma_{h, \partial \Omega_{k}}(\phi)=\sum_{\mathrm{F} \subset \partial \Omega_{k}} \gamma_{h, \partial \Omega_{k}}(\phi) I_{\mathrm{F}}^{0} 1+\sum_{\mathrm{E} \subset \partial \Omega_{k}} \gamma_{h, \partial \Omega_{k}}(\phi) I_{\mathrm{E}}^{0} 1+\sum_{p \in \partial \Omega_{k} \cap \mathcal{V}_{\Gamma}} \gamma_{h, \partial \Omega_{k}}(\phi) \varphi_{p} \quad \text { on } \partial \Omega_{k} .
$$

Then,

$$
\begin{align*}
\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi) & =\sum_{\mathrm{F} \subset \partial \Omega_{k}}\left(\gamma_{h, \mathrm{~F}}(\phi)-\gamma_{h, \partial \Omega_{k}}(\phi)\right) I_{\mathrm{F}}^{0} 1+\sum_{\mathrm{E} \subset \partial \Omega_{k}}\left(\gamma_{h, \mathrm{E}}(\phi)-\gamma_{h, \partial \Omega_{k}}(\phi)\right) I_{\mathrm{E}}^{0} 1 \\
& +\sum_{p \in \partial \Omega_{k} \cap \mathcal{V}_{\Gamma}}\left(\phi(p)-\gamma_{h, \partial \Omega_{k}}(\phi)\right) \varphi_{p} \\
& =\sum_{\mathrm{F} \subset \partial \Omega_{k}} \gamma_{h, \mathrm{~F}}\left(\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right) I_{\mathrm{F}}^{0} 1+\sum_{\mathrm{E} \subset \partial \Omega_{k}} \gamma_{h, \mathrm{E}}\left(\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right) I_{\mathrm{E}}^{0} 1 \\
& +\sum_{p \in \partial \Omega_{k} \cap \mathcal{V}_{\Gamma}}\left(\phi(p)-\gamma_{h, \partial \Omega_{k}}(\phi)\right) \varphi_{p} \quad \text { on } \partial \Omega_{k} \tag{4.10}
\end{align*}
$$

Furthermore, we get by the direct calculation and the inverse estimates

$$
\begin{align*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2} & \lesssim H \sum_{\mathrm{E} \subset \partial \Omega_{k}}\left|\gamma_{h, \mathrm{E}}\left(\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right|^{2}+h\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \infty, \mathcal{W}_{k}}^{2} \\
& \lesssim\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2} \tag{4.11}
\end{align*}
$$

and

$$
\begin{align*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2} & \lesssim H^{2} \sum_{\mathrm{F} \subset \partial \Omega_{k}}\left|\gamma_{h, \mathrm{~F}}\left(\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right|^{2}+h^{2}\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \infty, \partial \Omega_{k}}^{2} \\
& \lesssim\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2} \tag{4.12}
\end{align*}
$$

It can be verified that

$$
\left|\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right|^{2} \lesssim H^{-2}\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2}
$$

which, together with (4.12), gives (4.9).

Lemma 4.2 Let $\phi \in W_{h}(\Gamma)$, and let $\phi_{0}$ be defined by (4.7). Then,

$$
\begin{equation*}
\sum_{k=1}^{\mathrm{N}} \omega_{k}\left|E_{0} \phi_{0}\right|_{1, \Omega_{k}}^{2} \lesssim \log (H / h)\|\phi\|_{*, \Gamma}^{2} \tag{4.13}
\end{equation*}
$$

Proof. It suffices to estimate (3.19) more carefully. It is easy to see that

$$
\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)=\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right) .
$$

By the triangle inequality, we deduce

$$
\begin{aligned}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \mathcal{W}_{k}}^{2} & \lesssim\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2}+\left\|\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right\|_{0, \mathcal{W}_{k}}^{2} \\
& \lesssim\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2}+H\left|\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right|^{2}
\end{aligned}
$$

and

$$
\begin{aligned}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \partial \Omega_{k}}^{2} & \lesssim\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2}+\left\|\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right\|_{0, \partial \Omega_{k}}^{2} \\
& \lesssim\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2}+H^{2}\left|\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}-\gamma_{h, \partial \Omega_{k}}(\phi)\right)\right|^{2}
\end{aligned}
$$

Substituting (4.8)-(4.9) and (4.12) into the above two inequalities, yields

$$
\begin{equation*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \mathcal{W}_{k}}^{2} \lesssim\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \mathcal{W}_{k}}^{2}+H^{-1}\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2}, \tag{4.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \partial \Omega_{k}}^{2} \lesssim\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{0, \partial \Omega_{k}}^{2} \tag{4.15}
\end{equation*}
$$

Combining (4.14) with (2.11), leads to

$$
\begin{equation*}
\left\|\phi_{0}-\gamma_{h, \partial \Omega_{k}}\left(\phi_{0}\right)\right\|_{0, \mathcal{W}_{k}}^{2} \lesssim \log (H / h) H^{-1}\left\|\phi-\gamma_{h, \partial \Omega_{k}}(\phi)\right\|_{\frac{1}{2}, \partial \Omega_{k}}^{2} \tag{4.16}
\end{equation*}
$$

Plugging (4.16) and (4.15) in (3.19), and using Friedrichs' inequality, leads to (4.13).

Proof of Theorem 4.2. The idea of the proof is standard. But, for readers' convenience, we still give a complete proof of this theorem below. One needs to establish a suitable decomposition for $v_{h} \in V_{h}(\Omega)$

$$
\begin{equation*}
v_{h}=v_{0}+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} v_{\mathrm{E}}+\sum_{\Gamma_{i j}} v_{i j} \tag{4.17}
\end{equation*}
$$

with

$$
v_{0} \in V_{h}^{0}(\Omega), \quad v_{\mathrm{E}} \in V_{h}\left(\Omega_{\mathrm{E}}\right) \text { and } v_{i j} \in V_{h}^{p}\left(\Omega_{i j}\right)
$$

This decomposition should satisfy the stability condition

$$
\begin{equation*}
\left(B_{0} v_{0}, v_{0}\right)+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}}\left(B_{\mathrm{E}} v_{\mathrm{E}}, v_{\mathrm{E}}\right)+\sum_{\Gamma_{i j}}\left(B_{i j} v_{i j}, v_{i j}\right) \lesssim \log ^{2}(H / h)\left(A_{h} v_{h}, v_{h}\right) \tag{4.18}
\end{equation*}
$$

Set $\phi_{h}=\left.v_{h}\right|_{\Gamma}$, and define $\phi_{0} \in W_{h}^{0}(\Gamma)$ as in (4.7). For $\mathrm{E} \in \mathcal{E}_{\Gamma}$ and $\mathrm{F}=\Gamma_{i j} \in \mathcal{F}_{\Gamma}$, set

$$
\phi_{\mathrm{E}}=I_{\mathrm{E}}^{0}\left[\phi_{h}-\gamma_{h, \mathrm{E}}\left(\phi_{h}\right)\right], \text { and } \phi_{i j}=I_{\mathrm{F}}^{0}\left[\phi_{h}-\gamma_{h, \mathrm{~F}}\left(\phi_{h}\right)\right] .
$$

It is clear that $\phi_{\mathrm{E}} \in \tilde{W}_{h}(\mathrm{E})$ and $\phi_{i j} \in \tilde{W}_{h}(\mathrm{~F})$. It is easy to see that

$$
\begin{equation*}
\phi_{h}=\phi_{0}+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} \phi_{\mathrm{E}}+\sum_{\Gamma_{i j}} \phi_{i j} \tag{4.19}
\end{equation*}
$$

Let $v_{\mathrm{E}}$ be the zero extension of $\phi_{\mathrm{E}}$, and set $v_{0}=E_{0} \phi_{0}$. Then,

$$
v_{\mathrm{E}} \in V_{h}\left(\Omega_{\mathrm{E}}\right) \quad \text { and } \quad v_{0} \in V_{h}^{0}(\Omega)
$$

Let $v_{i j}^{H} \in V_{h}^{p}\left(\Omega_{i j}\right)$ be defined such that $\left.v_{i j}^{H}\right|_{\Gamma_{i j}}=\phi_{i j}$, and $v_{i j}^{H}$ is discrete harmonic on $\Omega_{i}$ and $\Omega_{j}$. Set

$$
\begin{equation*}
v_{k}^{p}=\left.\left(v_{h}-v_{0}-\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}} v_{\mathrm{E}}-\sum_{\Gamma_{i j}} v_{i j}^{H}\right)\right|_{\Omega_{k}} . \tag{4.20}
\end{equation*}
$$

Then, we have $v_{k} \in V_{h}^{p}\left(\Omega_{k}\right)$ by (4.19). For each $k$, let $m_{k}$ be the number of faces that belong to $\partial \Omega_{k}$. Define

$$
v_{i j}=v_{i j}^{H}+v_{i}^{p} / m_{i}+v_{j}^{p} / m_{j} .
$$

It suffices to verify (4.18) for the functions defined above. For convenience, set

$$
\mathcal{G}\left(v_{h}\right)=\left(B_{0} v_{0}, v_{0}\right)+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}}\left(B_{\mathrm{E}} v_{\mathrm{E}}, v_{\mathrm{E}}\right)+\sum_{\Gamma_{i j}}\left(B_{i j} v_{i j}^{H}, v_{i j}^{H}\right) .
$$

Using (4.1), (4.13) and the trace Theorem, we get

$$
\begin{equation*}
\left(B_{0} v_{0}, v_{0}\right) \lesssim \log (H / h)\left(A_{h} v_{h}, v_{h}\right) \tag{4.21}
\end{equation*}
$$

Moreover, we have by (4.2), (2.12) and (2.14)

$$
\begin{equation*}
\left(B_{\mathrm{E}} v_{\mathrm{E}}, v_{\mathrm{E}}\right) \lesssim\left\|v_{\mathrm{E}}\right\|_{0, \mathrm{E}}^{2} \lesssim \log ^{2}(H / h) \sum_{k \in \mathcal{Q}_{\mathrm{E}}} \omega_{k} \int_{\Omega_{k}}\left|\nabla v_{h}\right|^{2} d p \tag{4.22}
\end{equation*}
$$

Here, we have used the discrete norm to derive the first inequality for $v_{\mathrm{E}}$ vanishing at all nodes outside E. Besides, we deduce by (4.3), the definition of $v_{i j}^{H}$ and (2.13)

$$
\begin{equation*}
\left(B_{i j} v_{i j}^{H}, v_{i j}^{H}\right) \lesssim\left(\omega_{i}+\omega_{j}\right)\left\|\phi_{i j}\right\|_{H_{00}^{\frac{1}{2}}\left(\Gamma_{i j}\right)}^{2} \lesssim \log ^{2}(H / h)\left(\omega_{i}\left|v_{h}\right|_{1, \Omega_{i}}^{2}+\omega_{j}\left|v_{h}\right|_{1, \Omega_{j}}^{2}\right) . \tag{4.23}
\end{equation*}
$$

This, together with (4.21) and (4.22), leads to

$$
\begin{equation*}
\mathcal{G}\left(v_{h}\right) \lesssim \log ^{2}(H / h)\left(A_{h} v_{h}, v_{h}\right) \tag{4.24}
\end{equation*}
$$

On the other hand, it follows by (4.20) that

$$
\left|v_{k}^{p}\right|_{1, \Omega_{k}}^{2} \lesssim\left|v_{h}\right|_{1, \Omega_{k}}^{2}+\left|v_{0}\right|_{1, \Omega_{k}}^{2}+\sum_{\mathrm{E} \in \mathcal{E}_{\Gamma}}\left|v_{\mathrm{E}}\right|_{1, \Omega_{k}}^{2}+\sum_{\Gamma_{i j}}\left|v_{i j}^{H}\right|_{1, \Omega_{k}}^{2} .
$$

Combing this with (4.1)-(4.3), yields

$$
\sum_{k=1}^{\mathrm{N}} \omega_{k}\left|v_{k}^{p}\right|_{1, \Omega_{k}}^{2} \lesssim\left(A_{h} v_{h}, v_{h}\right)+\mathcal{G}\left(v_{h}\right) .
$$

This, together with (4.24), leads to

$$
\sum_{\Gamma_{i j}}\left(B_{i j}\left(v_{i}^{p} / m_{i}+v_{j}^{p} / m_{j}\right), v_{i}^{p} / m_{i}+v_{j}^{p} / m_{j}\right) \lesssim \sum_{k=1}^{\mathrm{N}} \omega_{k}\left|v_{k}^{p}\right|_{1, \Omega_{k}}^{2} \lesssim \log ^{2}(H / h)\left(A_{h} v_{h}, v_{h}\right) .
$$

By the definition of $v_{i j},(4.23)$ and the above inequality, we further get

$$
\sum_{\Gamma_{i j}}\left(B_{i j} v_{i j}, v_{i j}\right) \lesssim \log ^{2}(H / h)\left(A_{h} v_{h}, v_{h}\right) .
$$

Now, (4.18) is a direct consequence of the above inequality, together with (4.24).

## 5 Analysis for the stability of the extension $E_{\mathrm{F}}$

This section is devoted to verification of the stability condition (3.16), which has been used in Lemma 4.2. Since quasi-uniform meshes are considered, the analysis is a bit technical. Our basic idea is to introduce a suitable "approximate" extension of $E_{\mathrm{F}}$. This auxiliary extension is defined by positive integers, so that its stability can be verified more easily by estimating two finite sums associated with the discrete $H^{1}$ semi-norm.

### 5.1 An auxiliary extension

For face $\mathrm{F}=\Gamma_{i j}$ and a node in $\Omega_{i j}$, let $p^{\prime}, p^{\prime \prime}, d_{\mathrm{F}}$ and $\Lambda_{\mathrm{F}}$ be defined as in Subsection 4.1. Without loss of generality, we assume that $h \leq \min \left\{\left|p p^{\prime}\right|: p \in \mathcal{N}_{h} \cap\left(\Omega_{i j} \backslash \mathrm{~F}\right)\right\}$. For a positive number $x$, let $[x]$ denote the integer part of $x$. Set $n_{\mathrm{F}}=\left[\frac{d_{\mathrm{F}}}{h}\right]$. For any node $p \in \Omega_{i j} \backslash \mathrm{~F}$ with $p^{\prime} \in \mathrm{F}$, define

$$
m_{p}=\left[\frac{\left|p p^{\prime}\right|}{h}\right] \quad \text { and } \quad n_{p}=\left[\frac{\left|p^{\prime} p^{\prime \prime}\right|}{h}\right] .
$$

To understand the meaning of the integers $m_{p}$ and $n_{p}$ more intuitively, we image that the segments $p p^{\prime}$ and $p^{\prime} p^{\prime \prime}$ are divided into some smaller segments with the size $h$. Then, $m_{p}$ and $n_{p}$ can be viewed roughly as the numbers of the division points on the segments $p p^{\prime}$ and $p^{\prime} p^{\prime \prime}$, respectively.

By the definition of the set $\Lambda_{\mathrm{F}}$, it is easy to see that the positive integers $m_{p}$ and $n_{p}$ possess the properties:

Property A. For each node $p \in \Lambda_{\mathrm{F}}$, we have $1 \leq m_{p} \leq n_{p} \leq n_{\mathrm{F}}$;
Property B. For two integers $r$ and $k$ satisfying $1 \leq r \leq k \leq n_{\mathrm{F}}$, there are at most $O\left(n_{\mathrm{F}}\right)$ nodes $p \in \Lambda_{\mathrm{F}}$, such that these nodes $p$ define the same $m_{p}=r$ and $n_{p}=k$;

Property C. For an integer $k$ satisfying $1 \leq k \leq n_{\mathrm{F}}$, there are at most $O\left(m_{p} n_{\mathrm{F}}\right)$ nodes $p \in \Lambda_{\mathrm{F}}$, such that these nodes $p$ define the same $n_{p}=k$.

For the constant-like basis function $\phi_{\mathrm{F}}=I_{\mathrm{F}}^{0} 1$, define the auxiliary extension

$$
\left(E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)(p)=\left\{\begin{array}{cl}
1, & \text { if } p \in \mathrm{~F},  \tag{5.1}\\
1-\frac{m_{p}}{n_{p}}, & \text { if } p \in \Lambda_{\mathrm{F}}, \\
0, & \text { otherwise } .
\end{array}\right.
$$

For the verification of the stability (3.16), one needs to prove the following two inequalities

$$
\begin{equation*}
\left|\left(E_{\mathrm{F}}-E_{\mathrm{F}}^{\prime}\right) \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) \tag{5.2}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) \tag{5.3}
\end{equation*}
$$

### 5.2 Auxiliary results

In this subsection, we estimate two finite sums, which will be used to verify the inequality (5.3).

In the rest of this section, $p_{1}$ and $p_{2}$ always denote two neighboring nodes. Since the triangulation $\mathcal{T}_{h}$ is quasi-uniform, there exists a (fixed) positive integer $k_{0}$, such that $\left|p_{1} p_{2}\right| \leq k_{0} h$ for any two neighboring nodes $p_{1}$ and $p_{2}$. Then, it can be verified, by the definitions of $m_{p}$ and $n_{p}$, that

$$
\begin{equation*}
\left|m_{p_{1}}-m_{p_{2}}\right|, \quad\left|n_{p_{1}}-n_{p_{2}}\right| \leq k_{0}+1 . \tag{5.4}
\end{equation*}
$$

Lemma 5.1 Let $\Lambda_{\mathrm{F}}^{\partial}$ be the set of nodes $p$, at which the extension $E_{\mathrm{F}} \phi_{\mathrm{F}}$ vanish. Namely,

$$
\Lambda_{\mathrm{F}}^{\partial}=\left\{p \in \mathcal{N}_{h} \cap \bar{\Omega}_{i j}: \quad p \notin \mathrm{~F}, \quad p \notin \Lambda_{\mathrm{F}}\right\}
$$

Then,

$$
\begin{equation*}
h \sum_{p_{1} \in \Lambda_{\mathrm{F}}^{\partial}} \sum_{p_{2} \in \Lambda_{\mathrm{F}}}\left(1-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \lesssim d_{\mathrm{F}} \tag{5.5}
\end{equation*}
$$

Note that $p_{1}$ and $p_{2}$ above denote two neighboring nodes.
Proof. For a node $p_{2} \in \Lambda_{\mathrm{F}}$, there are at most finite nodes $p_{1} \in \Lambda_{\mathrm{F}}^{\partial}$, such that $p_{1}$ and $p_{2}$ are neighboring each other. Then,

$$
\begin{equation*}
h \sum_{p_{1} \in \Lambda_{\mathrm{F}}^{\partial}} \sum_{p_{2} \in \Lambda_{\mathrm{F}}}\left(1-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \lesssim h \sum_{p_{2} \in \tilde{\Lambda}_{\mathrm{F}}}\left(1-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \tag{5.6}
\end{equation*}
$$

where

$$
\tilde{\Lambda}_{\mathrm{F}}=\left\{p \in \Lambda_{\mathrm{F}}: \text { there is } p_{*} \in \Lambda_{\mathrm{F}}^{\partial} \text { such that } p \text { and } p_{*} \text { are neighboring }\right\}
$$

For $p \in \tilde{\Lambda}_{\mathrm{F}}$, let $p_{*} \in \Lambda_{\mathrm{F}}^{\partial}$ denote a neighboring node with $p$, and let $p_{*}^{\prime}$ be the projection of $p_{*}$ on the plane containing F. Set

$$
\tilde{\Lambda}_{\mathrm{F}}^{(1)}=\left\{p \in \tilde{\Lambda}_{\mathrm{F}}: \quad p_{*}^{\prime} \in \mathrm{F}\right\} \quad \text { and } \quad \tilde{\Lambda}_{\mathrm{F}}^{(2)}=\left\{p \in \tilde{\Lambda}_{\mathrm{F}}: \quad p_{*}^{\prime} \notin \mathrm{F}\right\}
$$

Then,

$$
\tilde{\Lambda}_{\mathrm{F}}=\tilde{\Lambda}_{\mathrm{F}}^{(1)} \cup \tilde{\Lambda}_{\mathrm{F}}^{(2)}
$$

It is easy to see that the set $\tilde{\Lambda}_{\mathrm{F}}^{(2)}$ contains $O\left(n_{\mathrm{F}}\right)$ nodes at most. It follows by (5.6) that

$$
\begin{align*}
h \sum_{p_{1} \in \Lambda_{\mathrm{F}}^{\partial}} \sum_{p_{2} \in \Lambda_{\mathrm{F}}}\left(1-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} & \lesssim h \sum_{p \in \tilde{\Lambda}_{\mathrm{F}}^{(1)}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2}+h \sum_{p \in \tilde{\Lambda}_{\mathrm{F}}^{(2)}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2} \\
& \lesssim h \sum_{p \in \tilde{\Lambda}_{\mathrm{F}}^{(1)}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2}+d_{\mathrm{F}} \tag{5.7}
\end{align*}
$$

Let $p \in \tilde{\Lambda}_{\mathrm{F}}^{(1)}$. From the definition of $\Lambda_{\mathrm{F}}^{\partial}$, we know that either $p_{*} \in \partial \Omega_{i j} \backslash \partial \mathrm{~F}$ or $\left|p_{*} p_{*}^{\prime}\right|>\left|p_{*}^{\prime} p_{*}^{\prime \prime}\right|$, where $p_{*}^{\prime \prime}$ is defined as in Subsection 3.1. By Assumption 3.1, we have $m_{p_{*}} \geq n_{p_{*}}$ in any case. Then, we deduce, by (5.4), that

$$
m_{p} \geq m_{p_{*}}-\left(k_{0}+1\right) \geq n_{p_{*}}-\left(k_{0}+1\right) \geq n_{p}-2\left(k_{0}+1\right)
$$

This implies that

$$
\max \left\{1, n_{p}-k^{*}\right\} \leq m_{p} \leq n_{p}, \quad \forall p \in \tilde{\Lambda}_{\mathrm{F}}^{(1)}
$$

with $k^{*}=2\left(k_{0}+1\right)$. Thus, we have from Property A and Property B

$$
\begin{align*}
\sum_{p \in \tilde{\Lambda}_{\mathrm{F}}^{(1)}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2} & \lesssim n_{\mathrm{F}}\left\{\sum_{n_{p}=1}^{k^{*}} \sum_{m_{p}=1}^{k^{*}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2}+\sum_{n_{p}=k^{*}+1}^{n_{\mathrm{F}}} \sum_{m_{p}=n_{p}-k^{*}}^{n_{p}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2}\right\} \\
& \lesssim n_{\mathrm{F}}\left\{\left(k^{*}\right)^{2}+\sum_{n_{p}=k^{*}+1}^{n_{\mathrm{F}}} \sum_{m_{p}=n_{p}-k^{*}}^{n_{p}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2}\right\} . \tag{5.8}
\end{align*}
$$

It is easy to see that

$$
\sum_{m_{p}=n_{p}-k^{*}}^{n_{p}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2} \leq \sum_{m_{p}=n_{p}-k^{*}}^{n_{p}}\left(1-\frac{n_{p}-k^{*}}{n_{p}}\right)^{2} \leq\left(k^{*}\right)^{3} \cdot \frac{1}{n_{p}^{2}}
$$

Plugging this in (5.8), and note that $k^{*}$ is a constant, leads to

$$
h \sum_{p \in \tilde{\Lambda}_{\mathrm{F}}^{(1)}}\left(1-\frac{m_{p}}{n_{p}}\right)^{2} \lesssim d_{\mathrm{F}}\left(\left(k^{*}\right)^{2}+\left(k^{*}\right)^{3} \sum_{n_{p}=k^{*}+1}^{n_{\mathrm{F}}} \frac{1}{n_{p}^{2}}\right) \lesssim d_{\mathrm{F}}
$$

Combining (5.7) with the above inequality, gives the desired result.

Lemma 5.2 The following inequality holds for two neighboring nodes $p_{1}$ and $p_{2}$

$$
\begin{equation*}
h \sum_{p_{1}, p_{2} \in \Lambda_{\mathrm{F}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) . \tag{5.9}
\end{equation*}
$$

Proof. For $p_{1} \in \Lambda_{\mathrm{F}}$, define

$$
\Lambda_{\mathrm{F}, p_{1}}=\left\{p \in \Lambda_{\mathrm{F}}: p \text { is neighboring with } p_{1}\right\}
$$

Let $k_{0}$ be defined by (5.4). Then, we have from Property A and Property B

$$
\begin{align*}
h \sum_{p_{1}, p_{2} \in \Lambda_{\mathrm{F}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} & \lesssim h \cdot \frac{d_{\mathrm{F}}}{h} \sum_{n_{p_{1}}=1}^{n_{\mathrm{F}}} \sum_{m_{p_{1}}=1}^{n_{p_{1}}} \sum_{p_{2} \in \Lambda_{\mathrm{F}, p_{1}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \\
& =d_{\mathrm{F}} \sum_{n_{p_{1}}=1}^{k_{0}+1} \sum_{m_{p_{1}}=1}^{n_{p_{1}}} \sum_{p_{2} \in \Lambda_{\mathrm{F}, p_{1}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \\
& +d_{\mathrm{F}} \sum_{n_{p_{1}}=k_{0}+2}^{n_{\mathrm{F}}} \sum_{m_{p_{1}}=1}^{k_{0}} \sum_{p_{2} \in \Lambda_{\mathrm{F}, p_{1}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \\
& +d_{\mathrm{F}} \sum_{\mathrm{F}}^{n_{\mathrm{F}}} \sum_{n_{p_{1}}=k_{0}+2}^{n_{p_{1}}} \sum_{m_{p_{1}}=k_{0}+1}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \\
& =I_{1}+I_{2}+I_{3} . \tag{5.10}
\end{align*}
$$

It is clear that $I_{1} \lesssim d_{\mathrm{F}}$. When $m_{p_{1}} \leq k_{0}$, we deduce from (5.4) that $m_{p_{2}} \leq 2 k_{0}+1$ for $p_{2} \in \Lambda_{\mathrm{F}, p_{1}}$. Note that $k_{0}$ is a constant, we get

$$
\begin{align*}
I_{2} & \lesssim d_{\mathrm{F}} \sum_{n_{p_{1}}=k_{0}+2}^{n_{\mathrm{F}}} \sum_{p_{2} \in \Lambda_{\mathrm{F}, p_{1}}}\left(\frac{1}{n_{p_{1}}^{2}}+\frac{1}{n_{p_{2}}^{2}}\right) \\
& \lesssim d_{\mathrm{F}} \sum_{n_{p_{1}}=k_{0}+2}^{n_{\mathrm{F}}}\left(\frac{1}{n_{p_{1}}^{2}}+\frac{1}{\left(n_{p_{1}}-k_{0}-1\right)^{2}}\right) \lesssim d_{\mathrm{F}} . \tag{5.11}
\end{align*}
$$

It follows by (5.4) that

$$
I_{3} \lesssim d_{\mathrm{F}} \sum_{n_{p_{1}}=k_{0}+2}^{n_{\mathrm{F}}} \sum_{m_{p_{1}}=k_{0}+1}^{n_{p_{1}}} \sum_{k=-k_{0}-1}^{k_{0}+1} \sum_{r=-k_{0}-1}^{k_{0}+1}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{1}}+r}{n_{p_{1}}+k}\right)^{2}
$$

$$
\begin{aligned}
& =d_{\mathrm{F}} \sum_{\substack{n_{p_{1}}=k_{0}+2}}^{n_{\mathrm{F}}} \sum_{m_{p_{1}}=k_{0}+1}^{n_{p_{1}}} \sum_{k=-k_{0}-1}^{k_{0}+1} \sum_{r=-k_{0}-1}^{k_{0}+1}\left(\frac{k m_{p_{1}}-r n_{p_{1}}}{n_{p_{1}}\left(n_{p_{1}}+k\right)}\right)^{2} \\
& \lesssim d_{\mathrm{F}} \sum_{\substack{n_{p_{1}}=k_{0}+2}}^{\sum_{n_{p_{1}}=k_{0}+1}^{n_{p_{1}}}} \frac{m_{p_{1}}^{2}+n_{p_{1}}^{2}}{n_{p_{1}}^{2}\left(n_{p_{1}}-k_{0}-1\right)^{2}} \\
& \lesssim d_{\mathrm{F}} \sum_{n_{p_{1}}=k_{0}+2}^{n_{\mathrm{F}}} \frac{n_{p_{1}}}{\left(n_{p_{1}}-k_{0}-1\right)^{2}} \lesssim d_{\mathrm{F}} \log n_{\mathrm{F}} .
\end{aligned}
$$

Plugging (5.11) and the above inequality in (5.10), leads to

$$
h \sum_{p_{1}, p_{2} \in \Lambda_{\mathrm{F}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) .
$$

### 5.3 Proof of Theorem 4.1

It suffices to verify that

$$
\begin{equation*}
\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2}, \quad\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{j}}^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) \equiv\left|\phi_{\mathrm{F}}\right|_{H_{00}(\mathrm{~F})}^{2} . \tag{5.12}
\end{equation*}
$$

Let $E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}$ be the auxiliary extension defined by Subsection 5.1.
Step 1. Verify the inequality (5.2)
Let $p \in \Lambda_{\mathrm{F}}$. Then, we have

$$
\begin{equation*}
\left|\left(E_{\mathrm{F}} \phi_{\mathrm{F}}-E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)(p)\right|=\left|\frac{\left|p p^{\prime}\right|}{\left|p^{\prime} p^{\prime \prime}\right|}-\left[\frac{\left|p p^{\prime}\right|}{h}\right] /\left[\frac{\left|p^{\prime} p^{\prime \prime}\right|}{h}\right]\right| . \tag{5.13}
\end{equation*}
$$

Since

$$
\frac{\left|p p^{\prime}\right|}{h} /\left(\frac{\left|p^{\prime} p^{\prime \prime}\right|}{h}+1\right) \leq\left[\frac{\left|p p^{\prime}\right|}{h}\right] /\left[\frac{\left|p^{\prime} p^{\prime \prime}\right|}{h}\right] \leq\left(\frac{\left|p p^{\prime}\right|}{h}+1\right) / \frac{\left|p^{\prime} p^{\prime \prime}\right|}{h},
$$

we get

$$
\frac{\left|p p^{\prime}\right|}{\left|p^{\prime} p^{\prime \prime}\right|+h} \leq\left[\frac{\left|p p^{\prime}\right|}{h}\right] /\left[\frac{\left|p^{\prime} p^{\prime \prime}\right|}{h}\right] \leq \frac{\left|p p^{\prime}\right|+h}{\left|p^{\prime} p^{\prime \prime}\right|} .
$$

Plugging this in (5.13), leads to

$$
\begin{equation*}
\left|\left(E_{\mathrm{F}} \phi_{\mathrm{F}}-E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)(p)\right| \leq \frac{h}{\left|p^{\prime} p^{\prime \prime}\right|} \leq \frac{1}{n_{p}} . \tag{5.14}
\end{equation*}
$$

By the inverse estimate and the discrete $L^{2}$-norm, we get

$$
\begin{aligned}
\left|\left(E_{\mathrm{F}}-E_{\mathrm{F}}^{\prime}\right) \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} & \lesssim h^{-2} \|\left.\left(E_{\mathrm{F}}-E_{\mathrm{F}}^{\prime}\right) \phi_{\mathrm{F}}\right|_{0, \Omega_{i}} ^{2} \\
& \equiv h \sum_{p \in \Lambda_{\mathrm{F}}}\left(E_{\mathrm{F}} \phi_{\mathrm{F}}-E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)^{2}(p) .
\end{aligned}
$$

This, together with (5.14), yields

$$
\begin{equation*}
\left|\left(E_{\mathrm{F}}-E_{\mathrm{F}}^{\prime}\right) \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} \lesssim h \sum_{p \in \Lambda_{\mathrm{F}}} \frac{1}{n_{p}^{2}} . \tag{5.15}
\end{equation*}
$$

By Property A and Property C, we get

$$
\sum_{p \in \Lambda_{\mathrm{F}}} \frac{1}{n_{p}^{2}} \lesssim n_{\mathrm{F}} \sum_{n_{p}=1}^{n_{\mathrm{F}}} \frac{m_{p}}{n_{p}^{2}} \leq n_{\mathrm{F}} \sum_{n_{p}=1}^{n_{\mathrm{F}}} \frac{n_{p}}{n_{p}^{2}} \lesssim n_{\mathrm{F}} \log n_{\mathrm{F}} .
$$

Plugging this in (5.15), gives (5.2).
Step 2. Verify the inequality (5.3)
For simplicity of exposition, let $\Lambda_{\mathrm{F}}^{*}$ denote the set of neighboring node paring $\left(p_{1}, p_{2}\right)$, which satisfies $\left(E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)\left(p_{1}\right)-\left(E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)\left(p_{2}\right) \neq 0$. By the discrete $H^{1}$ semi-norm, we have

$$
\begin{equation*}
\left|E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} \equiv h \sum_{\left(p_{1}, p_{2}\right) \in \Lambda_{\mathrm{F}}^{*}}\left|\left(E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)\left(p_{1}\right)-\left(E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right)\left(p_{2}\right)\right|^{2} . \tag{5.16}
\end{equation*}
$$

For ease of notation, define

$$
\Lambda_{\mathrm{F}, b}=\left\{p \in \mathcal{N}_{h} \cap \mathrm{~F}: p \text { closes } \partial \mathrm{F}\right\} .
$$

It is easy to see that the set $\Lambda_{\mathrm{F}}^{*}$ can be decomposed into several groups: (a) $p_{1} \in \Lambda_{\mathrm{F}, b}$ and $p_{2} \in \partial \mathrm{~F}$; (b) $p_{1} \in \mathrm{~F}$ and $p_{2} \in \Lambda_{\mathrm{F}}$; (c) $p_{1} \in \Lambda_{\mathrm{F}}^{\partial}$ and $p_{2} \in \Lambda_{\mathrm{F}}$; (d) $p_{1}, p_{2} \in \Lambda_{\mathrm{F}}$. It is certain that one can also consider the inverse situation with exchanging the positions of $p_{1}$ and $p_{2}$, but this will not affect the result.

It follows by (5.16) that

$$
\begin{align*}
\left|E_{\mathrm{F}}^{\prime} \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} & \equiv h \sum_{p_{1} \in \Lambda_{\mathrm{F}, b}} \sum_{p_{2} \in \partial \mathrm{~F}}(1-0)^{2}+h \sum_{p_{1} \in \mathrm{~F}} \sum_{p_{2} \in \Lambda_{\mathrm{F}}}\left(\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \\
& +h \sum_{p_{1} \in \Lambda_{\mathrm{F}}^{\partial}} \sum_{p_{2} \in \Lambda_{\mathrm{F}}}\left(1-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2}+h \sum_{p_{1}, p_{2} \in \Lambda_{\mathrm{F}}}\left(\frac{m_{p_{1}}}{n_{p_{1}}}-\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} . \tag{5.17}
\end{align*}
$$

It is clear that the set $\Lambda_{\mathrm{F}, b}$ contains only $O\left(n_{\mathrm{F}}\right)$ nodes. Then, we get for two neighboring nodes $p_{1}$ and $p_{2}$

$$
\begin{equation*}
h \sum_{p_{1} \in \Lambda_{\mathrm{F}, b}} \sum_{p_{2} \in \partial \mathrm{~F}}(1-0)^{2} \lesssim h \cdot \frac{d_{\mathrm{F}}}{h}=d_{\mathrm{F}} . \tag{5.18}
\end{equation*}
$$

When $p_{2} \in \Lambda_{\mathrm{F}}$ is neighboring with some $p_{1} \in \mathrm{~F}$, we have

$$
m_{p_{2}} \leq\left|p_{2} p_{2}^{\prime}\right| / h \leq\left|p_{1} p_{2}\right| / h \leq k_{0} .
$$

Besides, for any $p_{2} \in \Lambda_{\mathrm{F}}$, there are at most finite nodes $p_{1} \in \mathrm{~F}$, such that $p_{1}$ and $p_{2}$ are neighboring each other. Thus, we deduce by Property A and Property B

$$
\begin{equation*}
h \sum_{p_{1} \in \mathrm{~F}} \sum_{p_{2} \in \Lambda_{\mathrm{F}}}\left(\frac{m_{p_{2}}}{n_{p_{2}}}\right)^{2} \lesssim h \cdot \frac{d_{\mathrm{F}}}{h} \sum_{l=1}^{k_{0}} \sum_{n_{p_{2}}=1}^{n_{\mathrm{F}}}\left(\frac{l}{n_{p_{2}}}\right)^{2} \lesssim d_{\mathrm{F}} . \tag{5.19}
\end{equation*}
$$

Substituting (5.18)-(5.19), (5.5) and (5.9) into (5.17), yields (5.3).
Step 3. Prove the desired result (5.12)
Combining (5.2) and (5.3), yields

$$
\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{i}}^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) .
$$

In the same way, we can prove that

$$
\left|E_{\mathrm{F}} \phi_{\mathrm{F}}\right|_{1, \Omega_{j}}^{2} \lesssim d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) .
$$

On the other hand, it can be verified, by the discrete semi-norm in Subsection 2.2, that

$$
\left|\phi_{\mathrm{F}}\right|_{H_{00}}^{2}{ }_{\frac{1}{2}(\mathrm{~F})} \equiv d_{\mathrm{F}} \log \left(d_{\mathrm{F}} / h\right) .
$$

## 6 Numerical experiments

In this section, we give some numerical results to confirm our theoretical results described in section 4.

Consider the elliptic problem (2.1) with $\Omega$ being the cube $\Omega=[0,1]^{3}$, and the coefficient $a(x, y, z)$ being defined by

$$
\begin{array}{lr}
a(x, y, z)=10^{-5}, & \text { if } x, y \leq 0.5 \text { or } x, y \geq 0.5 \\
a(x, y, z)=1, & \text { otherwise } .
\end{array}
$$

The source function $f$ is chosen in a suitable manner.
Let $\Omega$ be decomposed into N cube subdomains with the edge length $H$. To illustrate wide practicality of the new method, we consider tetrahedron elements instead of hexahedron elements. Let each subdomain be divided into tetrahedron elements with the size $h$ in the standard way, and use the usual $P_{1}$ finite element approximate space.

We solve the algebraic system associated with the equation (2.4) by PCG iteration with the preconditioner $B_{h}$ defined in Section 4 or Algorithm 4.2. Here, each local solver $B_{i j}$ is chosen as the symmetric multigrid preconditioner for the restriction of $A_{h}$ on the subspace $V_{h}^{0}\left(\Omega_{i j}\right)$. The iteration terminates when the relative remainder is less than $1.0 D-5$. The iteration counts are listed as the following tables.

Table 6.1
iteration counts for PCG with $B_{h}$

| $H / h$ | $H=1 / 4$ | $H=1 / 6$ | $H=1 / 8$ |
| :---: | :---: | :---: | :---: |
| 8 | 29 | 32 | 32 |
| 16 | 40 | 42 | 42 |
| 32 | 50 | 51 | 52 |

Table 6.2
iteration counts for Algorithm 4.2

| $H / h$ | $H=1 / 4$ | $H=1 / 6$ | $H=1 / 8$ |
| :---: | :---: | :---: | :---: |
| 8 | 24 | 26 | 26 |
| 16 | 31 | 33 | 33 |
| 32 | 39 | 41 | 41 |

These numerical results indicates that the convergence of the new preconditioner is stable with the subdomain number $N$, and depends slightly on the ratio $H / h$. These are just predicted by Theorem 4.1 and Theorem 4.2. Besides, the numerical results show that Algorithm 4.2 is indeed more efficient than the standard PCG.

## 7 Appendix

In this Appendix, we revise the definition of the extension $E_{\mathrm{F}}$ for the general case without Assumption 3.1.

Let $o_{\mathrm{F}} \in \mathrm{F}$ denote the center of the largest circle ${ }^{2}$ contained in F . Then, $o_{\mathrm{F}}$ satisfies $\operatorname{dist}\left(o_{\mathrm{F}}, \partial \mathrm{F}\right)=d_{\mathrm{F}}$. Through the center $o_{\mathrm{F}}$, we draw the perpendicular line $L$ of F . Let $q_{1}^{\partial}, q_{2}^{\partial} \in \partial \Omega_{i j}$ denote the two intersection points of $L$ with $\partial \Omega_{i j}$. Set

$$
s_{\mathrm{F}}=\min \left\{\left|o_{\mathrm{F}} q_{1}^{\partial}\right|,\left|o_{\mathrm{F}} q_{2}^{\partial}\right|\right\}
$$

Set $r_{\mathrm{F}}=d_{\mathrm{F}} / s_{\mathrm{F}}$. Since $\Omega_{i}$ and $\Omega_{j}$ are regular and convex domains, the ratio $r_{\mathrm{F}}$ is uniformly bounded for every F. With the same notations in Subsection 3.1, set

$$
\Lambda_{\mathrm{F}}=\left\{p \in \mathcal{N}_{h} \cap\left(\Omega_{i j} \backslash \mathrm{~F}\right): p^{\prime} \in \mathrm{F}, r_{\mathrm{F}}\left|p p^{\prime}\right| \leq\left|p^{\prime} p^{\prime \prime}\right|\right\}
$$

For a face $\mathrm{F}=\Gamma_{i j}$, we revise the extension $E_{\mathrm{F}}$ as :

$$
\left(E_{\mathrm{F}} \phi_{\mathrm{F}}\right)(p)=\left\{\begin{array}{cl}
1, & \text { if } p \in \mathrm{~F}  \tag{7.1}\\
1-\frac{r_{\mathrm{F}}\left|p p^{\prime}\right|}{\left|p^{\prime} p^{\prime \prime}\right|}, & \text { if } p \in \Lambda_{\mathrm{F}} \\
0, & \text { otherwise }
\end{array}\right.
$$

For such revision, the proof of Theorem 3.1 is almost the same with that in Section 5. We revise the definition of the positive integer $m_{p}$ for $p \in \Omega_{i j} \backslash \mathrm{~F}\left(p^{\prime} \in \mathrm{F}\right)$ as

$$
m_{p}=\left[\frac{r_{\mathrm{F}}\left|p p^{\prime}\right|}{h}\right]
$$

but reserve the definition of the auxiliary extension $E_{\mathrm{F}}^{\prime}$ given in (5.1). We need only to revise slightly the proof of Lemma 5.1. Below is the changes needed to make.

The inequality (5.4) now becomes to be

$$
\begin{equation*}
\left|m_{p_{1}}-m_{p_{2}}\right| \leq r_{F} k_{0}+1,\left|n_{p_{1}}-n_{p_{2}}\right| \leq k_{0}+1 \tag{7.2}
\end{equation*}
$$

for two neighboring nodes $p_{1}$ and $p_{2}$.
Let $\Lambda_{\mathrm{F}}^{\partial}$ and $\tilde{\Lambda}_{\mathrm{F}}$ be defined in Lemma 5.1. For $p \in \tilde{\Lambda}_{\mathrm{F}}$, let $p_{*} \in \Lambda_{\mathrm{F}}^{\partial}$ be some neighboring node with $p$. When $p_{*} \in \partial \Omega_{i j}$, we deduce that $r_{\mathrm{F}}\left|p_{*} p_{*}^{\prime}\right| \geq\left|p_{*}^{\prime} p_{*}^{\prime \prime}\right|$ by the definition of $r_{\mathrm{F}}$ and the convexity of $\Omega_{i}$ and $\Omega_{j}$. Then, we have $m_{p_{*}} \geq n_{p_{*}}$ for $p_{*} \in \Lambda_{\mathrm{F}}^{\partial}$. For $p \in \tilde{\Lambda}_{\mathrm{F}}$, we get by

$$
\begin{equation*}
m_{p} \geq m_{p_{*}}-\left(r_{\mathrm{F}} k_{0}+1\right) \geq n_{p_{*}}-\left(r_{\mathrm{F}} k_{0}+1\right) \geq n_{p}-\left(r_{\mathrm{F}}+1\right) k_{0}-2 \tag{7.2}
\end{equation*}
$$

This implies that

$$
\max \left\{1, n_{p}-k^{*}\right\} \leq m_{p} \leq n_{p}, \quad \forall p \in \tilde{\Lambda}_{\mathrm{F}}
$$

with $k^{*}=\left(r_{\mathrm{F}}+1\right) k_{0}+2$.
Remark 7.1 It is easy to see that the above revision is needed only for the case with $r_{\mathrm{F}}>1$, since $r_{\mathrm{F}} \leq 1$ implies that Assumption 3.1 is satisfied. When $r_{\mathrm{F}} \leq 1$, the use of the original definition of $E_{\mathrm{F}}$ (see Subsection 3.1) can reduce the cost for calculating $E_{\mathrm{F}} \phi_{\mathrm{F}}$.

## 8 Conclusions

We have developed a kind of simple nearly harmonic extension for the constant-like basis function. This extension is used to define a coarse subspace, and then is applied to construct a substructuring preconditioner with inexact subdomain solvers. The simplicity of the extension guarantees that the resulting preconditioner is easy to implement.

[^1]
## References

[1] Y. Achdou, Y. Maday, and O. Widlund, Iterative substructuring preconditioners for mortar element methods in two dimensions. SIAM J. Numer. Anal., 36(1999), 551580.
[2] C. Bernardi, Y. Maday and A. Patera, A new nonconforming approach to domain decomposition:the mortar element method. In: Nonlinear partial differential equations and their applications (H. Brezis and J. Lions (eds)), Pitman, 1994.
[3] C. Borgers, The Neumann-Dirichlet domain decomposition method with inexact solvers on the subdomains, Numer. Math., 55(1989), 123-136
[4] J. Bramble, J. Pasciak and A. Vassilev, Analysis of non-overlapping domain decomposition algorithms with inexact solvers, Math. Comput., 67(1998), 1-19
[5] J. Bramble, J. Pasciak and A. Schatz, The construction of preconditioners for elliptic problems by substructuring. IV. Math. Comp., 53(1989), 1-24.
[6] M. Dryja and O. Widlund, Towards a unified theory of domain decomposition algorithms for elliptic problems. Third International Symposium on Domain Decomposition Methods for Partial Differential Equations (Houston, TX, 1989), 3-21, SIAM, Philadelphia, PA, 1990.
[7] M. Dryja, M. Sarkis and O. Widlund, Multilevel Schwarz methods for elliptic problems with discontinuous coefficients in three dimentions, Numer. Math., 72(1996), No.3, 313-348
[8] M. Dryja, F. Smith and O. Widlund, Schwarz analysis of iterative substructuring algorithms for elliptic problems in three dimensions, SIAM J. Math. Anal., 31(1994), No.6, 1662-1694
[9] Q. Du, Optimization based nonoverlapping domain decomposition algorithms and their convergence, SIAM J. Numer. Anal., 39(2001), 1056-1077
[10] C. Farhat, J. Mandel and F. Roux, Optimal convergence properties of the FETI Domain Decomposition Method. Comput. Methods. Appl. Mech. Engrg., 115 (1994), 367-388.
[11] G. Haase, U. Langer and A. Meyer, The approximate Dirichlet domain decomposition method. Part II: applications to 2nd-order B. V.P.s, Computing, 47(1991), 153-167
[12] G. Haase, S. V. Nepomnyaschikh, Explicit extension operators on hierarchical grids, East-West J. Numer. Math., 5(1997), 231-248
[13] Q. Hu, Preconditioning Poincare-Steklov operators arising from domain decompositions with mortar multipliers, IMA J. Numer. Anal., 24(2004), 643-669
[14] Q. Hu, A new kind of multilevel solver for second order Steklov-Poincare operators, Domain Decomposition Methods in Science and Engineering XVII, Springer, 2008, 391-398
[15] Q. Hu, Non-overlapping Domain Decomposition Methods with A New Class of Multilevel Solvers, research report, 2005, No.ICM-05-22, ICMSEC, China
[16] Q. Hu and G. Liang, Acceleration of the non-symmetrized two-level iteration, Appl. Numer. Math., 41(2002), 305-323
[17] Q. Hu, Z. Shi and D. Yu, Efficient solvers for saddle-point problems arising from domain decompositions with Lagrange multipliers, SIAM J. Numer. Anal., 42(2004), No. 3, 905-933
[18] Q. Hu and J. Zou, A non-overlapping domain decomposition method for Maxwell's equations in three dimensions, SIAM J. Numer. Anal., 41(2003), 1682-1708
[19] J. Huang and J. Zou, Construction of explicit extension operators on general finite element grids, Appl. Numer. Math., 43(2002), 211-227
[20] A. Klawonn, O. Widlund and M. Dryja, Dual-Primal FETI methods for three- dimensional elliptic problems with Heterogeneous coefficients. SIAM J. Numer. Anal., 40(2002), 159-179.
[21] Y. Kuznetsov, Efficient iterative solvers for elliptic finite element problems on nonmatching grids. Russian J. Numer. Anal. and Math. Modeling, 10(1995), 3, 187-211.
[22] U. Langer and O. Steinbach, Boundary element tearing and interconnecting methods. Computing, 71(2003), 205-228.
[23] J. Mandel and M. Brezina, Balancing domain decomposition for problems with large jumps in coefficients. Math. Comp., 65(1996), 1387-1401.
[24] J. Mandel and R. Tezaur, On the convergence of a dual-primal substructuring method. Numer. Math., 88(2001), 543-558.
[25] B. Smith, P. Bjorstad and W. Gropp, Domain Decomposition: Parallel Multilevel Methods for Elliptic Partial Differential Equations. Cambridge University Press, 1996.
[26] B. Smith and O. Widlund, A domain decomposition algorithm using hierarchical basis, SIAM J. Sci. Stat. Comput., 11(1990), 1212-1220
[27] C. Tong, T. Chan and C. Kuo, A domain decomposition preconditioner based on a change to a multilevel nodal basis. SIAM J. Sci. Stat. Comput., 12(1991), 1486-1495.
[28] A. Toselli and O. Widlund, Domain Decomposition Methods--Algorithms and Theory, Springer, 2004.
[29] A. Toselli, O. Widlund and B. Wohlmuth, An iterative substructuring method for Maxwell's equations in two dimensions. Math. Comp., 70(2001), 935-949.
[30] P. Le Tallec, Domain Decomposition Methods in Computational Mechanies. Comput. Mech. Adv., 2: 121-220, 1994.
[31] J. Xu, Iterative methods by space decomposition and subspace correction. SIAM Review, 34(1992), 581-613.
[32] J. Xu and J. Zou, Some non-overlapping domain decomposition methods. SIAM Review, 40(1998), 857-914.
[33] H. Yserentant, On the multi-level splitting of finite element spaces, Numer. Math., 49(1986), 379-412


[^0]:    ${ }^{1}$ The work was supported by The Key Project of Natural Science Foundation of China G10531080, National Basic Research Program of China G2005CB321702 and Natural Science Foundation of China G10771178. (email: hqy@lsec.cc.ac.cn)

[^1]:    ${ }^{2}$ If F possesses two parallel edges, then there may exist infinite largest circles contained in F. For this case, one needs to consider two such circles, which are tangent with three or more edges of F . Besides, the point $o_{\mathrm{F}}$ can be also chosen as any "central" point of F . Different choices of $o_{\mathrm{F}}$ will result in different values of the constant in the right side of (5.5).

