

# RELATIVISTIC COSMOLOGY AND THE PIONEERS ANOMALY

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## Abstract

NASA spacecrafts has suffered from three anomalies. The Pioneers spacecrafts were decelerated, and their spin when not disturbed, was declining. On the other hand, fly-bys for gravity assists, appeared with extra speeds, relative to infinity. The Pioneers and fly-by anomalies are given now exact general relativistic full general solutions, in a rotating expanding Universe. We cite new evidence on the rotation of the Universe. Our solution seems to be the only one that solves the three anomalies.

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## 1. Introduction

Anderson et al.(2008), and Lämmerzahl et al.,2006), have alerted the scientific community about the fly-by anomaly: during Earth gravity assists, spacecraft has suffered from an extra-energization, characterized by a positive extra speed, such that, measured "at infinity", the hyperbolic orbiting object presented an empirically calculated  $\Delta V/V$  around  $10^{-6}$ .A formula was supplied,

$$\frac{\Delta V}{V} = \frac{2\omega R}{c} \quad (1.1)$$

where  $\omega$ ,  $R$  and  $c$  stand for the angular speed and radius of the central mass, and the speed of light in vacuo. T.L. Wilson, from NASA,(Houston) and H.-J. Blome (Aachen), delivered a lecture in Montreal, on July 17,2008, and called the attention to the fact that the most trusted cause for both this anomaly, and the Pioneers, would be "rotational dynamics"(Wilson and Blome ,(2008). One of us, had, by that time, published results on the Pioneers Anomaly, through the rotation of the Universe (Berman, 2007). Now, we shall address the three anomalies.

The Pioneers Anomaly is the deceleration of about  $-9.10^{-8}\text{cm.s}^{-2}$  suffered by NASA space-probes travelling towards outer space. It has no acceptable explanation within local Physics, but if we resort to Cosmology, it could be explained by the rotation of the Universe. Be cautious, because there is no center or axis of rotation. We are speaking either of a Machian or a General Relativistic cosmological vorticity. It could apply to each observed point in the Universe, observed by any observer. Another explanation, would be that our Universe obeys a variable speed of light Relativistic Cosmology, without vorticities.However, we shall see later that both models are equivalent. Thermal emission cannot be invoked, for it should also decelerate elliptical orbiters, but the deceleration only affects hyperbolic motion. It does not explain fly-bys, either. A secondary Pioneers anomaly refers to spinning down of the spacecraft, when they were not disturbed. Again, thermal emission cannot explain it.

In previous papers (Berman and Gomide, 2010; 2011; 2011a), by considering an exact but particular solution of Einstein's field equations for an expanding and rotating metric, found, by estimating the deceleration parameter of the present Universe, as  $q \approx -1/2$ , that the Universe appeared to possess a field of decelerations coinciding approximately with the Pioneers anomalous value (Anderson et al., 2002). We now shall consider the condition for an exact match with the Pioneers deceleration, with a large class of solutions in General Relativity. Sections 5 and 6, treat the second Pioneers anomaly, and the fly-by. In section 7, an alternative cosmological model will be presented, following an idea by Godlowski et al. (2004), which allows us to work with a non-modified RW's metric.

The key result for all these subjects, is that hyperbolic motion, extends towards infinity, and, thus, qualify for cosmological alternatives, and boundary conditions. The fly-bys, and the Pioneers, are in hyperbolic trajectories, when the anomalies appear, so that Cosmology needs to be invoked.

Ni (2008;2009), has reported observations on a possible rotation of the polarization of the cosmic background radiation, around 0.1 radians. As such radiation was originated at the inception of the Universe, we tried to estimate a possible angular speed or vorticity, by dividing 0.1 radians by the age of the Universe, obtaining about  $10^{-19}\text{rad.s}^{-1}$ . Compatible results were obtained by Chechin (2010) and Su and Chu (2009).

The numerical result is very close to the theoretical estimate, by Berman (2007),

$$\omega = c/R = 3.10^{-18}\text{rad.s}^{-1}. \tag{1.2}$$

where  $c$ ,  $R$  represent the speed of light in vacuum, and the radius of the causally related Universe.

When one introduces a metric temporal coefficient  $g_{00}$  which is not constant, the new metric includes rotational effects. The metric has a rotation of the tri-space (identical with RW's tri-space) around the orthogonal time axis. This will be our framework, except for Section 7.

## 2. On the four kinds of rotation in Relativistic Cosmology

Consider the metric line-element:

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu \quad . \quad (2.1)$$

If the observer is at rest,

$$dx^i = 0 \quad ( i = 1, 2, 3 ) \quad ,$$

while,

$$dx^0 = dt \quad . \quad (2.2)$$

This last equality defines a proper time; we called cosmic time, in Cosmology.

From the geodesics' equations, we shall have:

$$\frac{d^2 x^i}{ds^2} + \Gamma^i_{\alpha\beta} \frac{dx^\alpha}{ds} \frac{dx^\beta}{ds} = \Gamma^i_{00} \quad . \quad (2.3)$$

We then find:

$$g^{ij} \frac{\partial g_{i0}}{\partial t} = 0 \quad . \quad (2.4)$$

This defines a Gaussian coordinate system, which in general implies that:

$$\frac{\partial g_{i0}}{\partial t} = 0 \quad . \quad (2.5)$$

We must now reset our clocks, so that, the above condition is universal (valid for all the particles in the Universe), and then our metric will assume the form:

$$ds^2 = dt^2 - g_{ij}(\vec{x}, t) dx^i dx^j \quad . \quad (2.6)$$

If we further impose that, in the origin of time, we have:

$$g_{i0}(t = 0) = 0 \quad , \quad (2.7)$$

then by (2.5), we shall have:

$$g_{i0}(t) = 0 \quad . \quad (2.8)$$

The above defines a Gaussian normal coordinate system.

For a comoving observer, in a freely falling perfect fluid, the quadrivelocity  $u^\mu$  will obey:

$$u^i = 0 \quad , \quad (2.9)$$

while, if we normalize the quadrivelocity, we find, from the condition:

$$g_{\mu\nu}u^\mu u^\nu = 1 \quad , \quad (2.10)$$

that,

$$g_{00}u^0 = 1 \quad . \quad (2.11)$$

Though later we shall discuss the case  $g_{00} = g_{00}(t) \neq 1$  , it is usually imposed:

$$g_{00} = u^0 = 1 \quad . \quad (2.12)$$

When dealing with Robertson-Walker's metric, this is the usual procedure. By this means, we have a tri-space, orthogonal to the time axis.

Gaussian coordinate systems, in fact, imply that, with  $g_{0i} = 0$  , there are no rotations in the metric, and in each point we may define a locally inertial reference system.

Gaussian normal coordinates were called "synchronous" ; in an arbitrary spacetime, when we pick a spacelike hypersurface  $S_0$  , and we eject geodesic lines orthogonal to it, with constant coordinates  $x^1, x^2$  and  $x^3$  , while  $x^0 \equiv t + t_0$  , where  $t_0 = 0$  on  $S_0$  , then  $t$  is the proper time, whose origin is  $t = 0$  on  $S_0$  (see MTW, 1973) .

In the above treatment, cosmic time is "absolute", so that the measure of the age of the Universe, according to this "time", is not subject to a relative nature.

Now, we might ask whether the tri-space, orthogonal to the time axis, could rotate relative to this axis. Berman (2008a, 2008b), has exactly defined this original idea, by identifying this rotation, which is different from all others, as will shall show bellow, with a time-varying metric coefficient  $g_{00}(t)$  . In the next Section, we relate the angular speed of the tri-space, relative to the time axis, with  $g_{00}(t)$  by means of,

$$\omega = \frac{1}{2} \frac{\dot{g}_{00}}{g_{00}} \quad . \quad (2.13)$$

In the above, we still may have a perfect fluid model.

Other type of rotation, is Raychaudhuri's vorticity, which is attached to non-perfect fluids (see, for instance, Berman (2007)). A third type of rotation, is what we usually call rotation of the metric, and is defined by non-diagonal terms, in the metric. For instance, Kerr's metric.

A fourth kind of rotation, is also attached to a perfect fluid model, like Berman's one: it is the Godolowski et al. (2004) idea, which is developed in Section 7 bellow.

### 3. Field equations for Gomide-Uehara-R.W.-metric

Consider first a temporal metric coefficient which depends only on  $t$  . The line element becomes:

$$ds^2 = -\frac{R^2(t)}{(1+kr^2/4)^2} [d\sigma^2] + g_{00}(t) dt^2 \quad . \quad (3.1)$$

The field equations, in General Relativity Theory (GRT) become:

$$3\dot{R}^2 = \kappa(\rho + \frac{\Lambda}{\kappa})g_{00}R^2 - 3kg_{00} \quad , \quad (3.2)$$

and,

$$6\ddot{R} = -g_{00}\kappa(\rho + 3p - 2\frac{\Lambda}{\kappa})R - 3g_{00}\dot{R}\dot{g}^{00} \quad . \quad (3.3)$$

Local inertial processes are observed through proper time, so that the four-force is given by:

$$F^\alpha = \frac{d}{d\tau}(mu^\alpha) = mg^{00}\ddot{x}^\alpha - \frac{1}{2}m\dot{x}^\alpha\left[\frac{\dot{g}_{00}}{g_{00}}\right] \quad . \quad (3.4)$$

Of course, when  $g_{00} = 1$  , the above equations reproduce conventional Robertson-Walker's field equations.

In order to understand equation (3.4) , it is convenient to relate the rest-mass  $m$  , to an inertial mass  $M_i$  , with:

$$M_i = \frac{m}{g_{00}} \quad . \quad (3.5)$$

It can be seen that  $M_i$  represents the inertia of a particle, when observed along cosmic time, i.e., coordinate time. In this case, we observe that we have two acceleration terms, which we call,

$$a_1^\alpha = \ddot{x}^\alpha \quad , \quad (3.6)$$

and,

$$a_2^\alpha = -\frac{1}{2g_{00}} (\dot{x}^\alpha \dot{g}_{00}) \quad . \quad (3.7)$$

The first acceleration is linear; the second, resembles rotational motion, and depends on  $g_{00}$  and its time-derivative.

If we consider  $a_2^\alpha$  a centripetal acceleration, we conclude that the angular speed  $\omega$  is given by,

$$\omega = \frac{1}{2} \left( \frac{\dot{g}_{00}}{g_{00}} \right) \quad . \quad (3.8)$$

The case where  $g_{00}$  depends also on  $r$  ,  $\theta$  and  $\phi$  was considered also by Berman (2008b) and does not differ qualitatively from the present analysis, so that, we refer the reader to that paper.

## 4. The exact solution to the Pioneers anomaly.

Consider the possible solution for the rotating case. We equate (1.2) and (3.8). We try a power-law solution for  $R$ , and find,

$$g_{00} = Ae^{t^{1-1/m}} \quad ( A = \text{constant} ) .$$

The scale-factor assumes a power-law , as in constant deceleration parameter models (Berman,1983;— and Gomide,1988),

$$R = (mDt)^{1/m} \quad , \quad (4.1)$$

where,  $m$  ,  $D = \text{constants}$ , and,

$$m = q + 1 > 0, \tag{4.2}$$

where  $q$  is the deceleration parameter. We may choose  $q$  as needed to fit the observational data.

We find,

$$H = (mt)^{-1}$$

If we now solve for energy-density of matter, and cosmic pressure, for a perfect fluid, the best way to present the calculation, and the most simple, is showing the matter energy-density  $\rho$  and the  $\sigma$ -or gravitational density parameter, to be defined below. We find,

$$\rho = \frac{3t^{-2}}{m^2 A \kappa} e^{-t^{1-1/m}} + 3k(mDt)^{-2/m} - \frac{\Lambda}{\kappa}$$

$$\sigma = \left(\rho + 3p - 2\frac{\Lambda}{\kappa}\right) = \left[\frac{6(1-m)t^{-2} + 3t^{-1/m}}{m^2 A \kappa}\right] e^{-t^{1-1/m}}$$

For the present Universe, the infinite time limit makes the above densities become zero.

It is possible to define,

$$\rho_{grav} = -\frac{3H^2}{\kappa g_{00}} \quad (\text{negative energy-density of the gravitational field})$$

Now, let us obtain the gravitational energy of the field,

$$E_{grav} = \rho_{grav} V = -(4/3)\pi R^3 \left(\frac{3H^2}{\kappa g_{00}}\right) = -\frac{c^4 R^{4-3m}}{2Gm^2 A(mD)^{1/m-3}} \tag{4.3}$$

## 5. The Second Pioneers Anomaly

The universal angular acceleration, is given by,

$$\alpha_u = \dot{\omega} = -cH/R = -c^2/R^2 \tag{5.1}$$

The spins of the Pioneers were telemetered, and as a surprise, it shows that the on-board measurements yield a decreasing angular speed, when the space-probes were not disturbed. Turyshev and Toth (2010), published the graphs (Figures 2.16 and 2.17 in their paper), from which it is clear that there is an angular deceleration of about 0.1 RPM per three years, or,

$$\alpha \approx -1.2 \times 10^{-10} \text{rad/s}^2. \tag{5.2}$$



As the diameter of the space-probes is about 10 meters, the linear acceleration is practically the Pioneers anomalous deceleration value ,in this case,  $-6.10^{-8}$  cm.s<sup>-2</sup>.The present solution of the second anomaly, confirms our first anomaly explanation.

I have elsewhere pointed out that we are in face of an angular acceleration frame-dragging field, for it is our result (5.1) above, for the Universe, that causes the result (5.2), through the general formula,

$$\alpha = -\frac{cH}{l} \tag{5.3}$$

where  $l$  is the linear magnitude of the localized body suffering the angular acceleration frame-dragging. For sub-atomic matter, this angular acceleration can become important.

## 6. The Solution of the Fly-by Anomaly

Consider a two-body problem, relative to an inertial system. The additional speed, measured at infinity, relative to the total speed, measured at infinity, is proportional to twice the tangential speed of the earth,  $w_e R_e$ , divided by the total speed  $V \rightarrow wR \approx c$  taken care of the Universe angular speed. In fact, we write,

$$\frac{\Delta V}{V} = \frac{V + \omega_e R_e - (V - \omega_e R_e)}{c} = \frac{2\omega_e R_e}{c} \approx 3 \times 10^{-6} \tag{6.1}$$

The trick, is that infinity, in a rotating Universe, like ours, has a precise meaning, through the angular speed formula (1.2).

## 7. The Godlowski et al. Rotation

We, now, shall follow an idea by Godlowski et al.(2004), and supply another General Relativistic model, of an expanding and rotating Universe. Their idea, is that the homogeneous and isotropic models, may still rotate relative to the local gyroscope, by means of a simple replacement, in the Friedman-RW's equations, of the kinetic term, by the addition of a rotational kinetic one.

Einstein's field equations, for a perfect fluid with perfect gas equation of state, and RW's metric, are two ones. The first, is an energy-density equation, the second is a definition of cosmic pressure, which can be substituted by energy momentum conservation. But, upon

writing the  $\dot{R}^2$  term, we shall add an extra rotational term, namely  $(\omega R)^2$ , in order to account for rotation. If we keep (3.1), the field equations become, for a flat Universe,

$$6H^2 = \kappa\rho + \Lambda \quad (7.1)$$

with,

$$p = \beta\rho \quad (7.2)$$

and,

$$\dot{\rho} = -3\sqrt{2}H\rho(1 + \beta) \quad (7.3)$$

The usual solution, with Berman's constant deceleration parameter models, render (Berman, 1983;—and Gomide, 1986),

$$R = (mDt)^{1/m} \quad (7.4)$$

$$H = (mt)^{-1} \quad (7.5)$$

$$\ddot{R} = -qH^2R = -(m - 1)H^2R \quad (7.6)$$

Notice that we may have a negative deceleration parameter, implying that the Universe accelerates, probably due to a positive cosmological "constant", but, nevertheless, it is subjected to a negative rotational deceleration, a kind of centripetal one, that acts on each observed point of the Universe, relative to each observer, given by relation (1.2), so that,

$$\ddot{R} = -qH^2R = qa_{cp} \quad (7.7)$$

We now supply the necessary relations among the constants, so that the above equations be observed, namely,

$$m = \frac{3}{2}\sqrt{2}(1 + \beta) = \pm \frac{\sqrt{6}}{\sqrt{\kappa\rho_0 + \Lambda_0}}$$

$$\rho = \rho_0 t^{-2} \quad (7.8)$$

$$\Lambda = \Lambda_0 t^{-2} \quad (7.9)$$

## 8. Final comments

If we calculate the centripetal acceleration corresponding to the above angular speed (1.2), we find, for the present Universe, with  $R \approx 10^{28}$ cm and  $c \simeq 3.10^{10}$ cm./s ,

$$a_{cp} = -\omega^2 R \cong -9.10^{-8} cm/s^2 \quad . \quad (8.1)$$

Our model of Section 4 has been automatically calculated alike with (1.2) and (8.1). This value matches the observed experimentally deceleration of the NASA Pioneers' space-probes. Equation (3.3) shows that one can have a positive cosmological lambda term accelerating the Universe, i.e.,  $\ddot{R} \geq 0$  along with a centripetal deceleration that is felt by any observer, relative any observed point, given by (8.1). Berman and Gomide (2010) had obtained a Machian General Relativistic solution, though particular. We call it Machian, because it parallels the semi-relativistic Machian solution by Berman (2007b).

A cosmologist has made very important criticisms on our work. First, he says why do not the planets in the solar system show the calculated deceleration on the Pioneers? The reason is that elliptical orbits are closed, and localized. You do not feel the expansion of the universe in the sizes of the orbits either. In General Relativity books, authors make this explicit. You do not include Hubble's expansion in Schwarzschild's metric. But, those space probes that undergo hyperbolic motion, which orbits extend towards infinity, they acquire cosmological characteristics, like, the given P.A. deceleration. Second objection, there are important papers (Rievers and Lämmerzahl, (2011); Francisco et al. (2011); Cuesta (2011)) which resolve the P.A. with non-gravitational Physics. Our answer, that is OK, we have now alternative explanations. However, in the Introduction of this paper, we have responded why thermal emission is no good an explanation because it does not explain the other two anomalies neither why the elliptical orbiters did not suffer the same deceleration; as to Cuesta (2011) he also has no explanation for the other two anomalies. This does not preclude ours. Third, cosmological reasons were discarded, including rotation of the Universe. The problem is that those discarded cosmologies, did not employ the correct metric. For instance, they discarded rotation by examining Godel model, which is non expanding, and with a strange metric. The two kinds of rotating and expanding metrics we employ now, were not discarded or discussed by the authors cited by this cosmologist. Then, the final question, is how come that a well respected author, dismissed planetary Coriolis forces induced by rotation of

distant masses, by means of the constraints in the solar system. The answer is the same above, and also that one needs to consider Mach's Principle on one side, and the theoretical meaning of vorticities, because one is not speaking in a center or an axis of rotation or so. When we say, in Cosmology, that the Universe rotates, we mean that there is a field of vorticities, just that. The whole idea is that Cosmology does not enter the Solar System except for non-closed orbits that extend to outer space. For the Gomide Uehara RW's metric, it is the tri-space that rotates relative to the orthogonal time-axis.

Another cosmologist pointed out a different "problem". He was discussing the prior paper, to the present one (Berman and Gomide, 2010). He objects, that the angular speed formula of ours, is coordinate dependent. Now, when you choose a specific metric, you do it thinking about the kind of problem you have to tackle. After you choose the convenient metric, you forget tensor calculus, and you work with coordinate-dependent relations. They work only for the given metric, of course.

The solutions of Section 4, and Section 7, are in fact a large class of solutions, for they embrace any possible deceleration parameter value, or, any power-law scale-factor. Our solution with the rotation of the Universe, is the only unified explanation that applies to the three NASA anomalies.

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## References

Adler, R.J.; Bazin, M.; Schiffer, M. (1975) - *Introduction to General Relativity*, 2<sup>nd</sup> Edition, McGraw-Hill, New York.

Anderson, J.D. et al. (2002) - "*Study of the anomalous acceleration of Pioneer 10 and 11*". - Phys. Rev. D 65, 082004 (2002)

Anderson, J.D. et al. (2008) - PRL100,091102.

Berman, M.S. (1983) - *Special Law of Variation for Hubble's Parameter*, Nuovo Cimento **74B**, 182-186.

Berman, M.S. (2007) - *Introduction to General Relativistic and Scalar-Tensor Cosmologies*, Nova Science Publisher, New York. (see Section 7.12)

Berman, M.S. (2007b) - *The Pioneer Anomaly and a Machian Universe* - Astrophysics and Space Science, **312**, 275. Los Alamos Archives, <http://arxiv.org/abs/physics/0606117>.

Berman, M. S. (2008a) - *A General Relativistic Rotating Evolutionary Universe*, Astrophysics and Space Science, **314**, 319-321.

Berman, M. S. (2008b) - *A General Relativistic Rotating Evolutionary Universe - Part II*, Astrophysics and Space Science, **315**, 367-369.

Berman, M.S.; Gomide, F.M. (1988) - *Cosmological Models with Constant Deceleration Parameter* - GRG **20**,191-198.

Berman, M.S.; Gomide, F.M. (2010) - Los Alamos Archives arXiv:1011.4627v3. Submitted.

Berman, M.S.; Gomide, F.M. (2011) - *On the Rotation of the Zero-Energy Expanding Universe*, Chapter 12, in "The Big-Bang: Theory, Assumptions and Problems", ed. by O'Connell and Hale, Nova Science, N.Y., to be published.

Berman, M.S.; Gomide, F.M. (2011a) - *Relativistic Cosmology and the Pioneers Anomaly*, Los Alamos Archives, arXiv:1106.5388 .

Chechin, L.M. (2010) - Astron. Rep. **54**, 719.

Godlowski, W. et al. (2004) - Los Alamos Archives, arxiv:astro-ph/0404329 .

Francisco, F. et al. (2011) - Los Alamos Archives, arXiv:1103.5222 .

Cuesta, H.J.M. (2011) - Los Alamos Archives, arXiv:1105.2832 .

Gomide, F.M.; Uehara, M. (1981) - Astronomy and Astrophysics, **95**, 362.

Lämmerzahl, C. et al. (2008) - in *Lasers, Clocks, and dragg-free Control*, ed. by Dittus, Lammerzahl and Turyshev, Springer, Heidelberg.

MTW - Misner, C.W.; Thorne, K.S.; Wheeler, J.A. (1973) - *Gravitation*, Freeman, San Francisco.

Ni, Wei-Tou (2008) - Progress Theor. Phys. Suppl. **172**, 49-60.

Ni, Wei-Tou (2009) - Int. Journ. Modern Phys. **A24**, 3493-3500.

Rievers, B.; Lammerzahl, C. (2011) - Los Alamos Archives, arXiv:1104.3985 .

Sabbata, de V.; Gasperini, M. (1979) - *Lettere al Nuovo Cimento*, **25**, 489.

So, L.L.; Vargas, T. (2006) - Los Alamos Archives, gr-qc/0611012 .

Su, S.C.; Chu, M.C. (2009) - Ap.J. 703, 354

Turyshev, S.G.; Toth, V.T. (2010) - *The Pioneer Anomaly*, Los Alamos Archives,  
arXiv:1001.3686

Weinberg, S. (1972) - *Gravitation and Cosmology*, Wiley. New York.

Wilson, T.L.; Blome, H.-J. (2008) - arXiv:0508.4067