# Local structure of the set of steady-state solutions to the 2D incompressible Euler equations 

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#### Abstract

It is well known that the incompressible Euler equations can be formulated in a very geometric language. The geometric structures provide very valuable insights into the properties of the solutions. Analogies with the finite-dimensional model of geodesics on a Lie group with left-invariant metric can be very instructive, but it is often difficult to prove analogues of finite-dimensional results in the infinite-dimensional setting of Euler's equations. In this paper we establish a result in this direction in the simple case of steady-state solutions in two dimensions, under some non-degeneracy assumptions. In particular, we establish, in a non-degenerate situation, a local one-to-one correspondence between steady-states and co-adjoint orbits.


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## 1 Introduction

### 1.1 Background

We consider the 2d Euler equations for inviscid incompressible fluid in a smooth bounded domain $\Omega \subset \mathbb{R}^{2}$ :

$$
\left.\begin{array}{rlrl}
u_{t}+(u \cdot \nabla) u+\nabla p & =0 \\
\operatorname{div} u & =0
\end{array}\right\} \quad \begin{aligned}
& \text { in } \Omega,  \tag{1.2}\\
& u \cdot N=0
\end{aligned} \quad \text { at } \partial \Omega,
$$

where $N$ denotes a unit normal to $\partial \Omega$. It is well-known that the equations have a rich geometric structure, coming from their interpretation as equations for a geodesic flow in the group $\mathcal{D}_{\text {vol }}(\Omega)$ of volume preserving diffeomorphisms. The modern mathematical investigations exploring this geometric structure were initiated by the 1966 paper of V. I. Arnold [1]. The geometric point of view has lead to important insights about Euler's equations, often by analogies with the finite-dimensional situation of the geodesic flow on a Lie group equipped with a left-invariant metric. [1] The passage from such finite-dimensional models to the infinite-dimensional setting of Euler's equations is often impeded by a common difficulty in infinite dimensions: the assumptions which are needed for straightforward generalizations of basic results of finite-dimensional Calculus (such as the Implicit Function Theorem) are too strong to be satisfied in situations of interest. For important advances in this direction, see for example [4, 5, 13, 16].

Our goal in this paper is to establish rigorously, in certain cases, a geometric picture of the structure of the set of steady-states of Euler's equations (1.1), (1.2) suggested by the

[^1]finite-dimensional situation. The main theorem will be for the case when $\Omega$ is diffeomorphic to an annulus, but it seems reasonable to set up the problem in greater generality.

Our results are best described in the vorticity formulation of the equations. We assume that $\Omega \subset \mathbb{R}^{2}$ is a bounded smooth domain. The connected components of $\partial \Omega$ will be denoted by $\Gamma_{0}, \Gamma_{1}, \ldots, \Gamma_{l}$, with $\Gamma_{0}$ bordering the unbounded connected component of $\mathbb{R}^{2} \backslash \Omega$.

Any (smooth) divergence-free velocity field $u$ in $\Omega$ satisfying $u \cdot N=0$ at $\partial \Omega$ can be represented by a stream function $\psi$, defined by

$$
\begin{equation*}
u=\nabla^{\perp} \psi=\binom{-\psi_{x_{2}}}{\psi_{x_{1}}} . \tag{1.3}
\end{equation*}
$$

Clearly, $\psi$ is defined uniquely by $u$ up to a constant. Therefore, without loss of generality we set $\left.\psi\right|_{\Gamma_{0}}=0$. The vorticity $\omega$ is defined by

$$
\begin{equation*}
\omega=u_{2, x_{1}}-u_{1, x_{2}}=\Delta \psi . \tag{1.4}
\end{equation*}
$$

The stream function $\psi$ is determined by $\omega$ and suitable boundary conditions. To identify these boundary conditions, we note first from (1.2) that, for each fixed time, $\psi$ is constant also on any other boundary component. However, the constants can depend on time, i.e. $\left.\psi\right|_{\Gamma_{i}}$ may not, in general, be constant during the evolution for $i=1, \ldots, l$. But by Kelvin's theorem on conservation of circulation of $u, \gamma_{i}=\int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}$ are constant along the flow. (Note that, by the Gauss-Green theorem and the divergence-free condition in (1.1), the circulation $\gamma_{0}$ around $\Gamma_{0}$ is determined by $\gamma_{i}, 1 \leq i \leq l$.) The constants $\gamma_{i}, 1 \leq i \leq l$ will be considered as fixed parameters throughout the paper. Therefore, denoting by $\tau$ the unit tangent vector to the boundary given by rotating the normal $N$ by $\pi / 2$, our boundary conditions will be

$$
\begin{align*}
\left.\psi\right|_{\Gamma_{0}} & =0  \tag{1.5}\\
\left.\frac{\partial \psi}{\partial \tau}\right|_{\Gamma_{i}} & =0, \quad i=1, \ldots, l  \tag{1.6}\\
\int_{\Gamma_{i}} \frac{\partial \psi}{\partial N} & =\gamma_{i}, \quad i=1, \ldots, l \tag{1.7}
\end{align*}
$$

We introduce the subspace

$$
\begin{equation*}
\mathcal{U}_{\gamma_{i}}=\left\{\psi \in C_{\bar{\Omega}}^{\infty} \text { satisfying (1.5) - (1.7) }\right\}, \tag{1.8}
\end{equation*}
$$

of the space of stream functions

$$
\begin{equation*}
\mathcal{U}=\left\{\psi \in C_{\bar{\Omega}}^{\infty} \text { satisfying (1.5) }-(1.6)\right\} . \tag{1.9}
\end{equation*}
$$

Together with $\omega$ and the equation

$$
\begin{equation*}
\Delta \psi=\omega \tag{1.10}
\end{equation*}
$$

the boundary conditions (1.5)-(1.7) uniquely determine $\psi$ (see [10]). Denoting by

$$
\begin{equation*}
\{f, g\}=f_{x_{1}} g_{x_{2}}-f_{x_{2}} g_{x_{1}} \tag{1.11}
\end{equation*}
$$

the 2 d Poisson bracket, equation (1.1) can be rewritten as

$$
\begin{equation*}
\omega_{t}+\{\psi, \omega\}=0, \tag{1.12}
\end{equation*}
$$

where $\psi$ is determined by $\omega$ through (1.10) and the boundary conditions (1.5)- (1.7). Equation (1.12) describes the transport of $\omega=\omega(t)$ by the group $\mathcal{D}_{\text {vol }}(\Omega)$ : we have

$$
\begin{equation*}
\omega(t)=\omega(0) \circ \eta^{-1}(t), \tag{1.13}
\end{equation*}
$$

where $\eta(t) \in \mathcal{D}_{\text {vol }}(\Omega)$ represents the particle trajectories and $\eta^{-1}$ denotes the inverse of $\eta$. In other words, letting for each smooth function $\omega$ on $\Omega$

$$
\begin{equation*}
\mathcal{O}_{\omega}=\left\{\omega \circ \eta^{-1}: \eta \in \mathcal{D}_{\mathrm{vol}}(\Omega)\right\} \tag{1.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\omega_{0}=\omega(0), \tag{1.15}
\end{equation*}
$$

the solution of (1.12) always stays on $\mathcal{O}_{\omega_{0}}$. Moreover, equation (1.12) can be thought of (formally) as a Hamiltonian system on $\mathcal{O}_{\omega_{0}}$, with the Hamiltonian given by the usual energy

$$
\begin{equation*}
\mathcal{E}(\omega)=\int_{\Omega} \frac{1}{2}|\nabla \psi|^{2}=\int_{\Omega}-\frac{1}{2} \omega \psi+\left.\sum_{j=1}^{l} \frac{1}{2} \gamma_{j} \psi\right|_{\Gamma_{j}} . \tag{1.16}
\end{equation*}
$$

The (formal) symplectic structure on $\mathcal{O}_{\omega_{0}}$ is of course of great independent interest, but we will not be concerned with it in this work. We will only study the equilibria, and these are (formally) characterized as the critical points of the restriction of $\mathcal{E}$ to the orbits $\mathcal{O}_{\omega} \cdot{ }^{[2]}$

To summarize, we formally have the following situation: the space of vorticities is foliated by the orbits $\mathcal{O}_{\omega}$, and the equilibria are the critical points of $\mathcal{E}$ restricted to the orbits. In finite dimension a routine application of the Implicit Function Theorem would imply that if $\mathcal{O}_{\bar{\omega}}$ is smooth near $\bar{\omega}$ and $\bar{\omega}$ is a non-degenerate critical point of $\mathcal{E}$ on $\mathcal{O}_{\bar{\omega}}$, then, near $\bar{\omega}$, the equilibria form a manifold transversal to the foliation by the orbits, of dimension equal to

[^2]the co-dimension of the orbits. In other words, in a non-degenerate situation, the equilibria are locally in one-to-one correspondence with the orbits.

Our goal is to establish an analogue of this correspondence in the context of (1.12). Several obstacles have to be overcome: the orbits $\mathcal{O}_{\omega}$ are typically not sub-manifolds of the space of vorticities if we work with the usual Banach spaces used in PDEs; certain linearized operators suffer from loss of derivatives; it is not clear what are good "coordinates" in which to do calculations. Our main result is, roughly speaking, that the difficulties can be resolved for steady-states for which the vorticity $\omega$ has no critical points in $\bar{\Omega}$. This assumption is of course restrictive. However, it is likely that in some situations the critical points of $\omega$ can genuinely complicate the picture and lead to some degeneracies, especially in the case of hyperbolic critical points. Elliptic critical points seem to be less dangerous. They still lead to difficulties for our method (namely certain linear maps are no longer "tame", see Section (1.3), but these might perhaps be manageable.

We now outline the main points of our approach. We start with the classical observation that any function $\psi$ satisfying

$$
\begin{equation*}
\Delta \psi=F(\psi) \tag{1.17}
\end{equation*}
$$

with the boundary conditions (1.5)- (1.7) gives a steady-state. This is easily seen from (1.12). Moreover, if $\bar{\psi}$ and $\bar{F}$ solve (1.17), (1.5) - (1.7), and if $\bar{\omega}=\bar{F}(\bar{\psi})$ has no critical points, then any nearby steady-state can be obtained in this way. This is one reason for the restriction on the geometry of $\Omega$ : the condition that $\bar{\omega}$ has no critical points imply that $\Omega$ is diffeomorphic to an annulus, as $\bar{\omega}$ is constant on $\partial \Omega$. (For the case where $\omega$ has a single elliptic critical point in a simply connected domain $\Omega$, a refinement of our method seems necessary.) The boundary of $\Omega$ consists then of an inner and outer boundary components,

$$
\begin{equation*}
\partial \Omega=\Gamma_{i} \cup \Gamma_{o} \tag{1.18}
\end{equation*}
$$

This is one characterization of the steady-states in the situation we wish to investigate: we see that they are, in some sense, locally parametrized by the functions $F$ such that

$$
\begin{equation*}
F^{\prime} \neq 0 . \tag{1.19}
\end{equation*}
$$

The map $F \mapsto \psi$, defined via (1.17) and (1.5)- (1.7) is not quite one-to-one, since changes to $\bar{F}$ outside the range of $\bar{\psi}$ do not change the solution $\bar{\psi}$, but this is not a serious problem.

It is worth remarking that steady-states of the form (1.17) with $F^{\prime} \neq 0$ naturally arise in the statistical theories of 2D Euler flows, see [12, 18, 21, and in Shnirelman's theory of mixing, see [20].

Our plan is to establish the correspondence between the functions $F$ in (1.22) and the orbits $\mathcal{O}_{\omega}$. By contrast with the finite-dimensional case, there is a simple obstacle showing
that we cannot really consider all orbits $\mathcal{O}_{\omega}$ : for any steady-state solution satisfying (1.5)(1.7) the vorticity $\omega$ must be constant along each boundary component. We introduce the space

$$
\begin{equation*}
\mathcal{F}=\left\{\omega \in C_{\bar{\Omega}}^{\infty}\left|\frac{\partial \omega}{\partial \tau}\right| \partial \Omega=0\right\} . \tag{1.20}
\end{equation*}
$$

Our main result (Theorem [1) establishes a local correspondence between the functions $F$ and the orbits contained in $\mathcal{F}$, i.e. with constant values of $\omega$ at the boundary components.

We need to introduce some (local) parametrization of the "space of orbits" $\mathcal{O}_{\omega}$. The distribution functions $A_{\omega}$ defined by

$$
\begin{equation*}
A_{\omega}(\lambda)=|\{x \in \Omega \mid \omega(x)<\lambda\}| \tag{1.21}
\end{equation*}
$$

provide a good option for those orbits $\mathcal{O}_{\omega} \subset \mathcal{F}$ with constant values of $\omega$ on the boundary components, and for which $\omega$ has no critical points, see Proposition 8, (The assumption that $\omega$ has no critical points surfaces again and imposes on $\Omega$ to be diffeomorphic to an annulus.) In fact it will be better to work with the inverses $A_{\omega}^{-1}$, for several reasons. One is that their domain, which is the interval $[0,|\Omega|]$, does not change, and another is the identity

$$
\begin{equation*}
A_{\omega}^{-1}=F \circ A_{\psi}^{-1} \tag{1.22}
\end{equation*}
$$

satisfied by the solutions of (1.17) in the case when $F^{\prime}>0$ (when $F^{\prime}<0$ a similar identity holds; in the proof we will perform computations assuming $F^{\prime}>0$ for simplicity, the case $F^{\prime}<0$ being completely analogous). This identity will be crucial for the analysis for the following reason: it shows that there is a chance for establishing some correspondence between $F$ and $A_{\omega}^{-1}$. The right-hand side of (1.22) is non-linear in $F$ (as $\psi$ depends on $F$ ), but the non-linear part comes in only through $\psi$ and hence it is regularized by equation (1.17). In some sense, the leading part of the dependence of $A_{\omega}^{-1}$ on $F$ behaves as a composition of $F$ with a fixed function, which, for many purposes, is almost the same as identity. Heuristically, the function $A_{\omega}^{-1}$ can be thought of as being obtained from $F$ by applying a kind of non-linear Fredholm map to $F$. This "Fredholmness" of the map $F \mapsto A_{\omega}^{-1}$ is crucial for our approach. On the other hand, the linearization of $A_{\omega}^{-1}$ suffers from loss of derivatives, and (1.22) shows it quite clearly:

$$
\begin{equation*}
\delta\left(F \circ A_{\psi}^{-1}\right)=(\delta F) \circ A_{\psi}^{-1}+\left(F^{\prime} \circ A_{\psi}^{-1}\right) \cdot \delta\left(A_{\psi}^{-1}\right) . \tag{1.23}
\end{equation*}
$$

The first term on the right-hand side is very good, but the second term contains $F^{\prime}$. This does not seem to be easy to avoid, and it can be overcome by working with the Nash-Moser Implicit Function Theorem, which will enable us to establish a good local correspondence between $F$ and $A_{\omega}^{-1}$ mentioned above, see Theorem 11. The correspondence between $F$ and $A_{\omega}^{-1}$ cannot be one-to-one, as already noted above, but only for the trivial reason that
in (1.17) the behavior of $F$ outside the range of $\psi$ does not affect the solution. Theorem 1 does contain an injective part, though, which says that, if two nearby steady-states have same distribution functions, then they are identical. Also, some natural non-degeneracy assumptions are needed, in the form of transversality conditions for linearized operators. Some conditions of this form are needed even in the finite-dimensional situation.

### 1.2 Statement of main result

We first introduce the two non-degeneracy assumptions of Theorem 1 .
We will denote $\bar{\psi}, \bar{F}$, and $\bar{\omega}=\bar{F}(\bar{\psi})$ the quantities associated with a reference steadystate. The function $F$ in (1.17) gives a good (local) parametrization of steady-states when $F^{\prime} \neq 0$ (modulo the lack of injectivity mentioned above). For this reason $\Omega$ is assumed to be diffeomorphic to an annulus. However, a well-defined map $F \mapsto \psi$ returning a solution to

$$
\begin{equation*}
\Delta \psi=F(\psi), \quad \psi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \psi}{\partial \tau}{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}, \quad 1 \leq i \leq l \tag{1.24}
\end{equation*}
$$

can be constructed in a neighborhood of $\bar{F}$, modulo some non-degeneracy condition discussed shortly, for an arbitrary number $l$ of boundary components. Thus for the construction of the map $F \mapsto \psi$ we will make no restriction on the topology of $\Omega$.

The map $F \mapsto \psi$ is well-defined provided the linear map

$$
\begin{equation*}
\Delta \phi-F^{\prime}(\psi) \phi=k, \quad \phi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \phi}{\partial \tau}{ }_{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0 \tag{1.25}
\end{equation*}
$$

is invertible for each $k \in C_{\bar{\Omega}}^{\infty}$ and for each $F$ in a neighborhood of $\bar{F}$. In fact, it is enough to make this assumption at the reference steady-state only:

$$
\left\{\begin{array}{l}
\text { the reference steady-state } \bar{\omega}=\bar{F}(\bar{\psi}) \text { is non-degenerate in the sense that }  \tag{ND1}\\
\Delta \phi-\bar{F}^{\prime}(\bar{\psi}) \phi=0, \quad \phi_{\mid \Gamma_{o}}=0, \left.\quad \frac{\partial \phi}{\partial \tau} \right\rvert\, \Gamma_{i}
\end{array}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0,1 \leq i \leq l .15 .\right.
$$

By the Fredholm alternative, the operator $\Delta-\bar{F}^{\prime}(\bar{\psi})$ is invertible with the boundary conditions of (ND1). Since we do not work in a Banach-space setting but with Fréchet spaces, it is not automatic that (ND1) implies that $\Delta-F^{\prime}(\psi)$ is invertible in $C^{\infty}$ for $F$ near $\bar{F}$. See [7] for a discussion of this crucial point and counterexamples, in particular Section I.5.5 of Part I.

The second non-degeneracy condition says that the steady-states and the co-adjoint orbits intersect trivially. (Since now we are using the parametrization of the "space of
orbits" via the distribution functions $A_{\omega}$, we will need to impose that $\Omega$ is diffeomorphic to an annulus since then it is crucial that $\omega$ has no critical points.) More precisely, a tangent $\nu$ to the set of steady-states at $\omega$ is a solution to the linearized steady-state equation $\omega=F(\psi)$, i.e. $\nu=\Delta \phi$ where $\phi$ solves, for some $f$,

$$
\begin{equation*}
\Delta \phi=F^{\prime}(\psi) \phi+f(\psi), \quad \phi_{\mid \Gamma_{o}}=0, \left.\quad \frac{\partial \phi}{\partial \tau} \right\rvert\, \Gamma_{i}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0 . \tag{1.26}
\end{equation*}
$$

A function $\nu$ is tangent to the co-adjoint orbit $\mathcal{O}_{\omega}$ at $\omega$ if there exists a stream function $\alpha \in \mathcal{U}$ such that $\nu=\{\omega, \alpha\}$. (This is immediate by linearizing $\omega \circ \eta$ at $\eta=$ Id.) The second non-degeneracy condition is as follows:

$$
\left\{\begin{array}{l}
\text { if } \phi \text { satisfies the linearized steady-state equation (1.26) at } \bar{\omega}=\bar{F}(\bar{\psi})  \tag{ND2}\\
\text { and if } \nu=\Delta \phi \text { is tangent to } \mathcal{O}(\bar{\omega}) \text { at } \bar{\omega} \text {, i.e. } \nu=\{\bar{\omega}, \alpha\} \text { for some } \alpha \in \mathcal{U}, \\
\text { then } \phi=0 .
\end{array}\right.
$$

Again, we emphasize that this non-degeneracy assumption is made at the reference steadystate only and not in an entire neighborhood.

Theorem 1 Let $\Omega \subset \mathbb{R}^{2}$ be diffeomorphic to an annulus with inner and outer boundary components $\Gamma_{i}, \Gamma_{o}$ respectively. Consider a smooth steady-state solution to Euler's equation on $\Omega$ with vorticity $\bar{\omega}$ and stream function $\bar{\psi}$ without critical points. Let $\bar{F}$ and $\gamma_{i}$ such that

$$
\begin{equation*}
\Delta \bar{\psi}=\bar{F}(\bar{\psi}), \quad \bar{\psi}_{\mid \Gamma_{0}}=0, \quad \frac{\partial \bar{\psi}}{\partial \tau}{ }_{\mid \Gamma_{i}}=0, \quad \gamma_{i}=\int_{\Gamma_{i}} \frac{\partial \bar{\psi}}{\partial N} \tag{1.27}
\end{equation*}
$$

(in particular $\bar{F}^{\prime} \neq 0$ ). Assume further that (ND1) and (ND2) are satisfied. Then, there exists a neighborhood $\mathcal{W}$ of $\bar{\omega}$ in $C_{\bar{\Omega}}^{\infty} \cap \mathcal{F}$ such that each co-adjoint orbit intersecting $\mathcal{W}$ contains exactly one smooth steady-state solution there.

Remark $\quad \mathcal{W}$ can be taken to be a $\|\cdot\|_{11}$-neighborhood. See Proposition 19 (proving the injective part of Theorem (1) and the Remark after the statement of Theorem [2] in the Appendix (on the existence part of the Nash-Moser Inverse Function Theorem).

An important aspect of the proof of Theorem 1 is that both non-degeneracy assumptions (ND1) and (ND2), which are made at the reference steady-state, are sufficient to imply nondegeneracy for steady-states in an entire neighborhood. It turns out that (ND1) and (ND2) are of exactly the same type, and the proof that $D T(F) f$ has a tame right-inverse will parallel the proof that $\Delta \phi-F^{\prime}(\psi) \phi=k$ (with the boundary conditions of (1.26)) has a
tame inverse. There is a kind of Fredholmness at work in both cases. To understand (ND1), consider the linear elliptic equation

$$
\begin{equation*}
\Delta \phi+c \phi=k \tag{1.28}
\end{equation*}
$$

parametrized by $c$ and with the same boundary conditions as in (1.26). The second term $K(c) \phi=c \phi$ is a "compact" perturbation of $\Delta \phi$ in the sense that, for each $n \geq 0, c \phi$ is $C^{n+2, \alpha}$ when $\Delta \phi$ is $C^{n, \alpha}$ (assuming $c$ is smooth). Assume that $\bar{c}$ is non-degenerate in the sense that $\Delta+\bar{c}$ has trivial kernel in $C^{\infty}$. Then, $\Delta+c$ is also non-degenerate if $\|c-\bar{c}\|_{0, \alpha}$ is sufficiently small. This comes from the estimate $\|(c-\bar{c}) \phi\|_{n, \alpha} \leq C\|c-\bar{c}\|_{0, \alpha}\|\phi\|_{n, \alpha}+C\|c-\bar{c}\|_{n, \alpha}\|\phi\|_{0, \alpha}$. In addition, the elliptic estimates are easily converted into tame estimates for the inverse.

The main thrust in using the Nash-Moser theorem (see Theorem 3) is to show that the derivative $D T$ has a tame right-inverse. For the problem at hand, this boils down to showing that a map of the form

$$
\begin{equation*}
h=g+K(F) g \tag{1.29}
\end{equation*}
$$

has a tame inverse (see Proposition 18 of Section 4.3). This is possible since, here again, the second term $K(F) g$ is a "compact perturbation" in the sense that, for each $n \geq 0$, $K(F) g$ is $C^{n+2, \alpha}$ when $g$ is $C^{n, \alpha}$ (and $F$ is smooth). On the other hand, estimates on $\|K(F) g-K(\bar{F}) g\|_{n, \alpha}$ are considerably more difficult to establish than those on $\|(c-\bar{c}) \phi\|_{n, \alpha}$. Again, estimates for $h=g+K(F) g$ yield easily tame estimates for the inverse.

The special case $F^{\prime}>0$
When $F^{\prime}>0$, it is clear that (ND1) is automatically satisfied at $\omega=F(\psi)$ (multiply (1.26) and integrate by parts). It turns out that (ND2) is also automatically satisfied. Let $\nu=\{\omega, \alpha\}$ satisfy (1.26). Then

$$
\begin{equation*}
\int_{\Omega} \frac{\nu^{2}}{F^{\prime}(\psi)}=\int_{\Omega} \phi \nu+\int_{\Omega} \frac{f(\psi)}{F^{\prime}(\psi)}\{\omega, \alpha\} . \tag{1.30}
\end{equation*}
$$

The second term of the right-hand side vanishes using (1.47). Integrating by parts,

$$
\begin{equation*}
\int_{\Omega}\left(\frac{\nu^{2}}{F^{\prime}(\psi)}+|\nabla \phi|^{2}\right)=0 \tag{1.31}
\end{equation*}
$$

which forces $\phi=0$ when $F^{\prime}>0$.
Theorem $\mathbb{1}$ is a statement about the local structure of the set of steady-states. However, the proof suggests that a global statement should hold in the case where $\bar{F}^{\prime}>0$. Namely one might speculate that the entire collection of steady-states satisfying $F^{\prime}>0$ is in one-to-one correspondence with their co-adjoint orbits. In other words, on any orbit containing a steady flow with $F^{\prime}>0$, this flow might be unique with this property. Note that this is certainly true for radial flows since the profile of $\omega$ and $A_{\omega}$ are the same up to a change of variables. In this case, (1.24) reduces to a second order ODE with two boundary conditions.

### 1.3 The Nash-Moser Inverse Function Theorem and tame estimates

The Nash-Moser Inverse Function Theorem for tame Fréchet spaces (see Theorem 3) bears certain important differences with the classical Inverse Function Theorem in Banach spaces. These differences concern in particular the notion of differentiability (which is in a sense weaker than the usual one for Banach spaces), the notion of tameness, and the fact that a right-inverse to the first derivative is assumed to exist in an entire neighborhood of $\bar{F}$ and not just at $\bar{F}$. The aim of this Section is to address these differences by giving precise definitions and clarifying certain assumptions. See Theorem 21 in the Appendix for the existence part of the Nash-Moser Inverse Function Theorem.

Our terminology and definitions mosly follow [7].

## Smooth maps of Fréchet spaces

Let the topology on a Fréchet space $\mathcal{X}$ be defined by a countable family of semi-norms $|\cdot|_{n}$, $n=0,1,2, \ldots$ We will be exclusively concerned with spaces $C_{\mathcal{K}}^{\infty}$ of smooth functions on a compact manifold $\mathcal{K}$ (possibly with boundary), and the semi-norms will be either the norms $\|\cdot\|_{n}$-norms (i.e. the sup-norms of derivatives up to $n$-th order) or the $\|\cdot\|_{n, \alpha}$-norms (i.e. the Hölder-norms of derivatives up to $n$-th order). Two gradings on $\mathcal{X}$ (i.e. two families of semi-norms $|\cdot|_{n}$ and $\left.|\cdot|_{n}^{\prime}, n=0,1,2, \ldots\right)$ are equivalent if they define the same topology.

Let now $\mathcal{X}, \mathcal{Y}$ be Fréchet spaces with semi-norms $|\cdot|_{n}, n=0,1,2, \ldots$ (for simplicity, we will use the same notation for the gradings of $\mathcal{X}$ and $\mathcal{Y}$ ). Let $\mathcal{B}$ be an open set and $P:(\mathcal{B} \subset \mathcal{X}) \rightarrow \mathcal{Y}$ a map between these Fréchet spaces. Continuity is defined as usual. In particular, when $\mathcal{X}$ and $\mathcal{Y}$ are of the form $C_{\mathcal{K}}^{\infty}$, then $P$ is continuous on $\mathcal{B}$ if, for each $n$, there exists $m=m(n)$ such that $P:\left(\mathcal{B},|\cdot|_{m(n)}\right) \rightarrow\left(\mathcal{Y},|\cdot|_{n}\right)$ is continuous. Also, two gradings are equivalent if the identity maps

$$
\begin{equation*}
\text { Id }:\left(\mathcal{X},\left\{|\cdot|_{n}\right\}_{n}\right) \rightarrow\left(\mathcal{X},\left\{|\cdot|_{n}^{\prime}\right\}_{n}\right), \quad \text { Id }:\left(\mathcal{X},\left\{|\cdot|_{n}^{\prime}\right\}_{n}\right) \rightarrow\left(\mathcal{X},\left\{|\cdot|_{n}\right\}_{n}\right) \tag{1.32}
\end{equation*}
$$

are both continuous. Clearly the $C^{k}$ - and the $C^{k, \alpha}$-gradings are equivalent on $C_{\mathcal{K}}^{\infty}$. The corresponding topology is called the $C^{\infty}$-topology. For the spaces of the form $C_{\mathcal{K}}^{\infty}$, we will prove continuity of maps using whichever grading is more convenient (the $C^{k, \alpha_{\text {- }} \text {-grading for }}$ operators involving elliptic equations, the $C^{k}$-grading otherwise).

The notion of differentiability, on the other hand, is in a sense weaker than the usual notion for maps of Banach spaces. The map $P$ is differentiable at $u \in \mathcal{B}$ if for each $v \in \mathcal{X}$ the limit

$$
\begin{equation*}
D P(u) v:=\lim _{t \rightarrow 0} \frac{P(u+t v)-P(u)}{t} \tag{1.33}
\end{equation*}
$$

exists in the Fréchet-topology, that is, there exists an element $D P(u) v \in \mathcal{Y}$ such that

$$
\begin{equation*}
\lim _{t \rightarrow 0}\left|\frac{P(u+t v)-P(u)}{t}-D P(u) v\right|_{n}=0 \quad \text { for each } \quad n=0,1,2, \ldots \tag{1.34}
\end{equation*}
$$

In other words, $P$ has Gâteaux-derivatives at $u$ in all directions. $P$ is continuously differentiable in $\mathcal{B}$ if the map

$$
\left\{\begin{array}{cccc}
(\mathcal{B} \subset \mathcal{X}) & \times & \mathcal{X} &  \tag{1.35}\\
u & & v & \\
D P(u) v
\end{array}\right\}
$$

is continuous jointly in the two variables $u$ and $v$. In the case the spaces are Banach spaces, this definition of differentiability is weaker than continuous Fréchet-differentiability, which is a usual assumption for the classical Inverse Function Theorem there. Partial derivatives for maps of several variables are defined in the usual way, as well as derivatives of higher order, e.g. (when they exist)

$$
\begin{equation*}
D^{2} P(u)\left(v_{1}, v_{2}\right)=\lim _{t \rightarrow 0} \frac{D P\left(u+t v_{2}\right) v_{1}-D P(u) v_{1}}{t} \tag{1.36}
\end{equation*}
$$

the limit again taken in the Fréchet-topology. A map is smooth if derivatives of all orders exist and are continuous. All maps will turn out to be smooth, but we will only need at most two derivatives in order to apply the Nash-Moser Inverse Function Theorem, see Theorem 3. (Specifically, we will prove that $T, D T, D^{2} T$, and a right-inverse $L$ to $D T$ are continuous.) Thus, we will call these maps smooth even though we only establish that they have continuous derivatives of order at most two.

A piece of terminology. If a map $P(u, v)$ is linear in $v$, then we will say that it is a family of linear maps and write it as $P(u) v$ to emphasize linearity in $v$. (Similar terminology and notations apply for maps of more than two variables.)

Many rules of the usual calculus apply. In particular, the first derivative $D P(u) v$ is linear in $v$ (see [7], Section 3.2 of Part I) the chain rule and the Fundamental Theorem of Calculus hold (see [7], Section 2 of Part I, for a definition and properties of Fréchet-space valued integrals), as well as Taylor's formula with integral remainder (see Theorem 3.5.6, p. 82 in Part I of [7]):

$$
\begin{equation*}
P(u+v)=P(u)+D P(u) v+\int_{0}^{1}(1-t) D^{2} P(u+t v)(v, v) d t \tag{1.37}
\end{equation*}
$$

Also, the Open Mapping Theorem holds: if a continuous linear map of Fréchet spaces is invertible, then it is a linear isomorphism, i.e. its inverse is again a continuous map of Fréchet spaces. On the other hand, if for a smooth family of linear maps, $P(\bar{u}) v$ is invertible, then it is not true in general that $P(u) v$ has an inverse for $u$ in a neighborhood of $\bar{u}$. This is in contrast to the Banach-space setting, where the set of invertible operators is open (this is related to the fact that the set of bounded linear maps on Banach spaces is itself a Banach space). This explains why the invertibility of the first derivative $D T(u) v$ for $u$ in an entire neighborhood must be assumed in order to apply the Nash-Moser Inverse Function Theorem.

## The tame Fréchet category

For our purposes (see [7], Section II.1, for a more general notion), a tame Fréchet space $\mathcal{X}$ is a Fréchet space, with semi-norms $|\cdot|_{n}, n=0,1,2, \ldots$, which comes with a family of smoothing operators $\{S(t)\}_{t>0}$ such that, for all $t>0$ and $u \in \mathcal{X}$,

$$
\begin{equation*}
|S(t) u|_{m} \leq C t^{m-l}|u|_{l}, \quad|u-S(t) u|_{l} \leq C t^{l-m}|u|_{m}, \quad m \geq l, \tag{1.38}
\end{equation*}
$$

where the constants depend on $m, l$, but not on $t$ nor $u$. Spaces of the form $C_{\mathcal{K}}^{\infty}$ are tame, see [7], Part II, Theorem 1.3.6, p. 137 and Corollary 1.3.7, p. 138. It is interesting, see [19], as well as Corollary 1.4.2, p. 176, Part II of [7], that the estimates (1.38) imply the interpolation inequalities

$$
\begin{equation*}
|u|_{i} \leq C|u|_{m}^{\frac{l-i}{l-m}}|u|_{l}^{\frac{i-m}{l-m}}, \quad m \leq i \leq l \tag{1.39}
\end{equation*}
$$

where the constants depend on $i, m, l$. These inequalities can otherwise be verified directly "by hand" in the $C^{k}-, C^{k, \alpha_{-}}$, or $H^{k}$-gradings, for example, see [3], [7] (Theorem 2.2.1, p. 143, Part II), (9).

A continuous map $P:(\mathcal{B} \subset \mathcal{X}) \rightarrow \mathcal{Y}$ of tame Fréchet spaces is tame if, for each $u_{0} \in \mathcal{B}$ there exist a neighborhood $\mathcal{V}$ of $u_{0}$ in $\mathcal{B}, r \in \mathbb{N}$ (the degree), $b \in \mathbb{N}$ (the base), and constants $C_{n}$ such that

$$
\begin{equation*}
|P(u)|_{n} \leq C_{n}\left(|u|_{n+r}+1\right), \quad n \geq b \tag{1.40}
\end{equation*}
$$

for any $u \in \mathcal{V}$. These are called tame estimates for $P$. We will usually suppress the dependence on $n$ for the constants and simply write $C$ (even though this dependence on $n$ is, of course, crucial). It can be proven, see Proposition 2 below, that, if a linear map $L$ is tame with degree $r$ and base $b$, then tame estimates can be derived in the form

$$
\begin{equation*}
|L u|_{n} \leq C|u|_{n+r}, \quad n \geq b \tag{1.41}
\end{equation*}
$$

for all $u \in \mathcal{X}$ without any restriction. A map is smooth tame if it is smooth and derivatives of all orders are tame.

We say that a grading $\left\{|\cdot|_{n}^{\prime}\right\}_{n}$ is tame equivalent to $\left\{|\cdot|_{n}\right\}_{n}$ if the identity maps (1.32) are both tame. (Obviously, this defines an equivalence relation.) In this case, the smoothing operators $S(t)$ satisfy again inequalities of the form (1.38) with $|\cdot|_{n}$ replaced by $|\cdot|_{n}^{\prime}$. A map $P:(\mathcal{B} \subset \mathcal{X}) \rightarrow \mathcal{Y}$ of tame Fréchet spaces remains tame if one replaces the gradings on $\mathcal{X}$ and $\mathcal{Y}$ with tame equivalent gradings. Note though that the degrees may be different in different gradings, and thus the choice of grading in which the tame estimates are derived should be made with care. We will derive tame estimates for all maps in the same grading, and our choice will be the $C^{n, \alpha}$-grading in order to take full advantage of elliptic regularity
afforded by the elliptic system (1.24). (Clearly, the $C^{n}$ - and $C^{n, \alpha}$-gradings on $C_{\mathcal{K}}^{\infty}$ are tame equivalent.)

## Further remarks on tame estimates

Clearly, a tame map with degree $r$ also has degree $r^{\prime} \geq r$. In turn, it is possible to choose the neighborhood $\mathcal{V}$ in which (1.40) holds to be a $|\cdot|_{r+b}$-neighborhood (by making $r$ larger and $\mathcal{V}$ smaller). Composition of tame maps is again a tame map.

On the other hand, the Open Mapping Theorem does not hold in the tame Fréchet category: if a tame linear map is invertible, then an inverse exists and is continuous, but it need not be tame. See [7], Section 1.5.5, Part I, for counterexamples.

For a map of several variables, these are allowed to have different degrees, e.g.

$$
\begin{equation*}
\left|P\left(u_{1}, u_{2}\right)\right|_{n} \leq C_{n}\left(\left|u_{1}\right|_{n+r_{1}}+\left|u_{2}\right|_{n+r_{2}}+1\right), \quad|u|_{r_{1}+b}<\delta_{1}, \quad|u|_{r_{2}+b}<\delta_{2}, \quad n \geq b \tag{1.42}
\end{equation*}
$$

For a family of linear maps, it is possible to do away with the restriction on the variables in which the map is linear (see Lemma 2.1.7, p. 143, Part II of [7)

Proposition 2 (Tame estimates for families of linear maps) Let $P(u) v$ be a family of linear maps. Then there exist constants $C_{n}$ such that

$$
\begin{equation*}
|P(u) v|_{n} \leq C_{n}\left(|u|_{n+r}+|v|_{n+s}+1\right), \quad n \geq b \tag{1.43}
\end{equation*}
$$

for $u$ in $a|\cdot|_{r+b}$-neighborhood and $v$ in $a|\cdot|_{s+b}$-neighborhood if and only if there exist constants $C_{n}^{\prime}$ such that

$$
\begin{equation*}
|P(u) v|_{n} \leq C_{n}^{\prime}\left(|u|_{n+r}|v|_{s}+|v|_{n+s}\right), \quad n \geq b \tag{1.44}
\end{equation*}
$$

for $u$ in $a|\cdot|_{r+b}$-neighborhood and any $v$ (without restriction).
This generalizes to maps linear in more than one variable. In our proof, we will systematically derive tame estimates in the format (1.44) when relevant.

## The Nash-Moser Inverse Function Theorem

Theorem 3 (Surjective part of the Nash-Moser theorem) Let $T:(\mathcal{B} \subset \mathcal{X}) \rightarrow \mathcal{Y}$ be a map of tame Fréchet spaces. Suppose that $T$ possesses first and second derivatives DT and $D^{2} T$, that $D T$ has a right-inverse $L$, and that all these maps are continuous and satisfy tame estimates. Then, for any $x_{0} \in \mathcal{B}$, there exists a neighborhood $\mathcal{V}\left(x_{0}\right)$ of $x_{0}$, a neighborhood $\mathcal{V}\left(y_{0}\right)$ of $y_{0}=T\left(x_{0}\right)$, and a map $R: \mathcal{V}\left(y_{0}\right) \rightarrow \mathcal{V}\left(x_{0}\right)$ such that $T(R(y))=y$ for $y \in \mathcal{V}\left(y_{0}\right)$. Furthermore, $R$ is continuous and tame, and if $T$ and $L$ are smooth tame, then so is $R$.

See the proof of Theorem 21 for the existence part of the surjective Nash-Moser Inverse Function Theorem. We refer to [7], Section III. 1 for proofs of further properties of $R$ (smoothness and tameness).

The proof Theorem 1 will heavily rely on smoothness and tameness of elementary maps of Fréchet spaces. The necessary lemmas are given in the Appendix and will be used countless times, often without explicit reference.

### 1.4 Examples of the "orbit calculus"

A rigorous interpretation of the orbits $\mathcal{O}_{\omega}$ as symplectic leaves would require some care. Instead, we give in this Section two examples of the "orbit calculus". Both results go back to Arnold [1], 2], but our calculations here are slightly different and do not resort to Lie-group theoretical interpretations.

Let $\omega=\Delta \psi$ be a critical point of the kinetic energy

$$
\begin{equation*}
\mathcal{E}(\omega)=\frac{1}{2} \int_{\Omega}|\nabla \psi|^{2} \mathrm{dvol}=-\frac{1}{2} \int_{\Omega} \omega \psi \mathrm{dvol}+\frac{1}{2} \sum_{i=1}^{l} \gamma_{i} \psi_{\mid \Gamma_{i}} \tag{1.45}
\end{equation*}
$$

restricted to its co-adjoint orbit (the $\gamma_{i}^{\prime}$ 's are fixed). Let $\nu=\{\alpha, \omega\}$ be tangent to the orbit $\mathcal{O}_{\omega}$ at $\omega$, and $\phi \in \mathcal{U}_{0}$ solving (1.26). Then the first derivative of the energy is given by

$$
\begin{equation*}
D \mathcal{E}(\omega) \nu=\int_{\Omega} \nabla \psi \cdot \nabla \psi \mathrm{dvol}=-\int_{\Omega} \psi \nu \mathrm{dvol}+\int_{\partial \Omega} \psi \frac{\partial \phi}{\partial N} d l=-\int_{\Omega} \psi\{\alpha, \omega\} \mathrm{dvol} . \tag{1.46}
\end{equation*}
$$

The identity

$$
\begin{equation*}
\int_{\Omega} f\{g, h\} \mathrm{dvol}=\int_{\Omega} g\{h, f\} \mathrm{dvol}-\int_{\partial \Omega} f g \frac{\partial h}{\partial N} d l \tag{1.47}
\end{equation*}
$$

gives

$$
\begin{equation*}
D \mathcal{E}(\omega) \nu=\int_{\Omega} \alpha\{\psi, \omega\} \mathrm{dvol} \tag{1.48}
\end{equation*}
$$

(the boundary terms vanish). Since $\alpha$ is arbitrary, we conclude that $\{\psi, \omega\}=0$.
Next we compute the second variation of $\mathcal{E}$ at a critical point:

$$
\begin{equation*}
D^{2} \mathcal{E}(\omega)(\nu, \nu)=\int_{\Omega}\left(|\nabla \phi|^{2}+\frac{\nu^{2}}{F^{\prime}(\psi)}\right) \mathrm{dvol} \tag{1.49}
\end{equation*}
$$

where $\psi$ solves (1.24) and $\phi$ solves (1.26).

Proof of $(\mathbf{1 . 4 9}) \quad$ Let $\omega_{\epsilon}=\omega \circ \eta_{\epsilon}$ with $\omega_{0}=\mathrm{id}$, and denote $\psi_{\epsilon}$ the corresponding stream functions. From the first derivative

$$
\begin{equation*}
\frac{d}{d \epsilon} \mathcal{E}\left(\omega_{\epsilon}\right)=\int_{\Omega} \nabla \psi_{\epsilon} \cdot \frac{\partial}{\partial \epsilon} \nabla \psi_{\epsilon} \mathrm{dvol} \tag{1.50}
\end{equation*}
$$

we find that the second derivative at $\epsilon=0$ is, posing $\dot{\psi}=\frac{\partial \psi_{\epsilon}}{\partial \epsilon}{ }_{\mid \epsilon=0}$ and $\ddot{\psi}=\frac{\partial^{2} \psi_{\epsilon}}{\partial \epsilon^{2}}{ }_{\mid \epsilon=0}$,

$$
\begin{equation*}
{\frac{d^{2}}{d \epsilon^{2}}}_{\mid \epsilon=0} \mathcal{E}\left(\omega_{\epsilon}\right)=\int_{\Omega}\left(|\nabla \dot{\psi}|^{2}+\nabla \psi \cdot \nabla \ddot{\psi}\right) \text { dvol. } \tag{1.51}
\end{equation*}
$$

The first term is $\int_{\Omega}|\nabla \phi|^{2}$ dvol. As for the second, integrating by parts we find

$$
\begin{equation*}
\int_{\Omega} \nabla \psi \cdot \nabla \ddot{\psi} \mathrm{dvol}=-\int_{\Omega} \psi \ddot{\omega} \mathrm{dvol} \tag{1.52}
\end{equation*}
$$

(the $\gamma_{i}$ 's are fixed). To calculate $\ddot{\omega}$, we can take $\eta_{\epsilon}$ as the flow corresponding to some $v=\nabla^{\perp} \alpha \in \mathcal{U}$. Then $\eta_{\epsilon}(x)=x+\epsilon v(x)+\frac{\epsilon^{2}}{2} \nabla_{v} v(x)+\ldots$ Taking second derivatives at $\epsilon=0$,

$$
\begin{align*}
{\frac{d^{2} \omega_{\epsilon}(x)}{d \epsilon^{2}}}_{\mid \epsilon=0} & =\frac{d^{2}}{d \epsilon^{2}}{ }_{\mid \epsilon=0} \omega\left(x+\epsilon v(x)+\frac{\epsilon^{2}}{2} \nabla_{v} v(x)+\ldots\right)  \tag{1.53}\\
& =\omega_{, k} v_{l} v_{k, l}+\omega_{, k l} v_{k} v_{l}=\left(\omega_{, k} v_{k} v_{l}\right)_{, l}  \tag{1.54}\\
& =\operatorname{div}(\{\alpha, \omega\} v)  \tag{1.55}\\
& =\operatorname{div}(\nu v) . \tag{1.56}
\end{align*}
$$

Integrating by parts, $-\int_{\Omega} \psi \operatorname{div}(\nu v) \mathrm{dvol}=\int_{\Omega} \nu \nabla \psi \cdot v$ dvol. Taking gradients of $\omega=F(\psi)$,

$$
\begin{equation*}
-\int_{\Omega} \psi \ddot{\omega} \mathrm{dvol}=\int_{\Omega} \nu \frac{\nabla \omega \cdot v}{F^{\prime}(\psi)} \mathrm{dvol}=\int_{\Omega} \frac{\nu^{2}}{F^{\prime}(\psi)} \mathrm{dvol} . \tag{1.57}
\end{equation*}
$$

This completes the proof.

### 1.5 Notation

Constants will generally be denoted by the same letter $C$, even in the derivation of tame estimates where it is important that they depend on the regularity index $n$. If $\mathcal{K}$ is the closure of a smooth, bounded region in Euclidean space (e.g. $\mathcal{K}=\bar{\Omega}$ or $[0,|\Omega|]$ ), then $C_{\mathcal{K}}^{n}$ is the space of $n$-times continuously differentiable functions on $\mathcal{K}$. The norm is given by

$$
\begin{equation*}
\|f\|_{n}:=\|f\|_{C_{\mathcal{K}}^{n}}:=\sup _{0 \leq j \leq n} \sup _{\mathcal{K}}\left|\nabla^{j} f\right| . \tag{1.58}
\end{equation*}
$$

$C_{\mathcal{K}}^{n, \alpha}$ denotes the subspace of functions in $C_{\mathcal{K}}^{n}$ whose derivatives up to order $n$ are Hölder continuous with exponent $\alpha$. Throughout, $\alpha$ will be a fixed constant in $(0,1)$. The norm is denoted

$$
\begin{equation*}
\|f\|_{n, \alpha}:=\|f\|_{C_{\mathcal{K}}^{n, \alpha}}:=\|f\|_{n}+\sum_{j=1}^{n}\left[\nabla^{j} f\right]_{\alpha} \quad[f]_{\alpha}:=\sup _{x \neq y \in \mathcal{K}} \frac{|f(x)-f(y)|}{|x-y|^{\alpha}} . \tag{1.59}
\end{equation*}
$$

We introduce the following spaces of functions with various regularity:

$$
\begin{align*}
\mathcal{F}^{n} & :=\left\{\omega \in C_{\bar{\Omega}}^{n}\left|\frac{\partial \omega}{\partial \tau}\right| \partial \Omega\right.  \tag{1.60}\\
\mathcal{U}^{n} & :=\left\{\psi \in \mathcal{F}^{n} \mid \psi_{\mid \Gamma_{0}}=0\right\},  \tag{1.61}\\
\mathcal{U}_{\gamma_{i}}^{n} & :=\left\{\psi \in \mathcal{U}^{n} \left\lvert\, \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}\right., 1 \leq i \leq l\right\},  \tag{1.62}\\
\mathcal{F}^{n, \alpha} & :=\left\{\omega \in C_{\bar{\Omega}}^{n, \alpha}\left|\frac{\partial \omega}{\partial \tau}\right| \partial \Omega\right.  \tag{1.63}\\
\mathcal{U}^{n, \alpha} & :=\left\{\psi \in \mathcal{F}^{n, \alpha} \mid \psi_{\mid \Gamma_{0}}=0\right\},  \tag{1.64}\\
\mathcal{U}_{\gamma_{i}}^{n, \alpha} & :=\left\{\psi \in \mathcal{U}^{n, \alpha} \left\lvert\, \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}\right., 1 \leq i \leq l\right\},  \tag{1.65}\\
\mathcal{F} & :=\left\{\omega \in C_{\bar{\Omega}}^{\infty}\left|\frac{\partial \omega}{\partial \tau}\right| \partial \Omega\right.  \tag{1.66}\\
\mathcal{U} & :=0\},  \tag{1.67}\\
\mathcal{U}^{\prime} & \left\{\psi \in \mathcal{F} \mid \psi_{\mid \Gamma_{0}}=0\right\},  \tag{1.68}\\
\mathcal{U}_{\gamma_{i}} & :=\left\{\psi \in \mathcal{U} \left\lvert\, \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}\right., 1 \leq i \leq l\right\} .
\end{align*}
$$

When $l=1$, we will also work with the open subsets

$$
\begin{align*}
\mathcal{F}_{+}^{n} & :=\left\{\omega \in \mathcal{F}^{n} \mid \nabla \omega \neq 0, \omega_{\mid \Gamma_{i}}<\omega_{\mid \Gamma_{o}}\right\},  \tag{1.69}\\
\mathcal{F}_{+}^{n, \alpha} & :=\left\{\omega \in \mathcal{F}^{n, \alpha} \mid \nabla \omega \neq 0, \omega_{\mid \Gamma_{i}}<\omega_{\mid \Gamma_{o}}\right\},  \tag{1.70}\\
\mathcal{F}_{+} & :=\left\{\omega \in \mathcal{F} \mid \nabla \omega \neq 0, \omega_{\mid \Gamma_{i}}<\omega_{\mid \Gamma_{o}}\right\} . \tag{1.71}
\end{align*}
$$

For simplicity, we will sometimes simply write $C^{n}, C^{n, \alpha}$, or $C^{\infty}$ for these spaces.

## 2 The solution operator $\psi=S(F)$

A solution operator $F \mapsto \psi$ returning a uniquely defined steady-state solution can be constructed in a neighborhood of the reference steady-state provided it satisfies the nondegeneracy condition (ND1). For this part of the proof no restriction on the geometry of $\Omega$ is necessary, and thus we consider (in this Section only) a bounded domain $\Omega \subset \mathbb{R}^{2}$ with outer boundary component $\Gamma_{0}$ and an arbitrary number $l$ of inner boundary components, $\Gamma_{i}, 1 \leq i \leq l$.

We assume given a reference steady-state:

$$
\begin{equation*}
\Delta \bar{\psi}=\bar{F}(\bar{\psi}), \quad \bar{\psi}_{\mid \Gamma_{o}}=0, \quad \frac{\partial \bar{\psi}}{\partial \tau}{ }_{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \bar{\psi}}{\partial N}=\gamma_{i}, \quad 1 \leq i \leq l \tag{2.1}
\end{equation*}
$$

where the $\gamma_{i}$ 's are fixed. It is assumed to satisfy the non-degeneracy condition (ND1). The goal of this section is to construct a solution operator $\psi=S(F)$ for $F$ in some neighborhood
of $\bar{F}$ (along with the desired estimates) returning a uniquely defined solution to the steadystate equation (1.24),

$$
\begin{equation*}
\Delta \psi=F(\psi), \quad \psi_{\mid \Gamma_{o}}=0, \left.\quad \frac{\partial \psi}{\partial \tau} \right\rvert\, \Gamma_{i}=0, \quad \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}, \quad 1 \leq i \leq l . \tag{2.2}
\end{equation*}
$$

Note that we do not assume in this section that $F^{\prime}$ does not vanish. Recall also that the case $F^{\prime}>0$ is special in that the corresponding solution automatically satisfies (ND1).

The boundary conditions in (2.2) define the affine space

$$
\begin{equation*}
\mathcal{U}_{\gamma_{i}}=\left\{\psi \in C \frac{\infty}{\Omega}\left|\psi_{\mid \Gamma_{o}}=0, \frac{\partial \psi}{\partial \tau}\right| \Gamma_{i}=0, \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N} d l=\gamma_{i}, 1 \leq i \leq l\right\} \tag{2.3}
\end{equation*}
$$

The $\gamma_{i}$ 's being fixed, its tangent space is

$$
\begin{equation*}
\left.\mathcal{U}_{0}=\left\{\phi \in C_{\bar{\Omega}}^{\infty}\left|\phi_{\mid \Gamma_{o}}=0, \frac{\partial \phi}{\partial \tau}\right| \Gamma_{i}\right], \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N} d l=0,1 \leq i \leq l\right\} \tag{2.4}
\end{equation*}
$$

We will also consider the following linear equation in $\phi \in \mathcal{U}_{0}$ parametrized by $c \in C \frac{\infty}{\Omega}$ :

$$
\begin{equation*}
\Delta \phi+c \phi=k, \quad \phi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \phi}{\partial \tau}{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0, \quad 1 \leq i \leq l \tag{2.5}
\end{equation*}
$$

### 2.1 Linear estimates

## Lemma 4 (Estimates for linear elliptic equations)

1. Given $\omega \in C_{\bar{\Omega}}^{\infty}$, there exists a unique $\psi \in C_{\bar{\Omega}}^{\infty}$ solving

$$
\begin{equation*}
\Delta \psi=\omega, \quad \psi_{\mid \Gamma_{0}}=0, \quad \frac{\partial \psi}{\partial \tau}{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}, \quad 1 \leq i \leq l \tag{2.6}
\end{equation*}
$$

It satisfies the tame estimates

$$
\begin{equation*}
\|\psi\|_{n+2, \alpha} \leq C\left(\|\omega\|_{n, \alpha}+\sum_{i=1}^{l}\left|\gamma_{i}\right|\right), \quad n \geq 0 \tag{2.7}
\end{equation*}
$$

2. For $c \in C \frac{\infty}{\bar{\Omega}}$ in a $\|\cdot\|_{0, \alpha}$-neighborhood and any $\phi \in \mathcal{U}_{0}$ satisfying

$$
\begin{equation*}
\Delta \phi+c \phi=k, \quad \phi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \phi}{\partial \tau}{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0, \quad 1 \leq i \leq l \tag{2.8}
\end{equation*}
$$

we have for $n \geq 0$

$$
\begin{equation*}
\|\phi\|_{n+2, \alpha} \leq C\left(\|\Delta \phi+c \phi\|_{n, \alpha}+\|c\|_{n, \alpha}\|\phi\|_{0, \alpha}\right) \tag{2.9}
\end{equation*}
$$

where the constant depends on $n$, but not on $c$ nor $\phi$.

## Proof

1 Hölder estimates on $\Delta \psi=\omega$ The construction of $\psi$ from $\omega$ is standard, see 10. In order to handle the boundary conditions of (2.2) one defines

$$
\begin{equation*}
u=\psi-\sum_{i=1}^{l} \psi_{\mid \Gamma_{i}} g_{i} \tag{2.10}
\end{equation*}
$$

where $g_{i} \in C_{\bar{\Omega}}^{\infty}, 1 \leq i \leq l$, are fixed functions with $g_{i \mid \Gamma_{0}}=0, g_{i \mid \Gamma_{i}}=1$, and $g_{i \mid \Gamma_{j}}=0$ for $j \neq i$. Then $u$ satisfies

$$
\begin{equation*}
\Delta u=\omega-\psi_{\mid \Gamma_{i}} \Delta g_{i}, \quad u_{\mid \partial \Omega}=0 \tag{2.11}
\end{equation*}
$$

for which we have the Schauder estimates: for $n \geq 0$,

$$
\begin{equation*}
\|u\|_{n+2, \alpha} \leq C\left(\|\Delta u\|_{n, \alpha}+\|u\|_{0, \alpha}\right) \tag{2.12}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\|\psi\|_{n+2, \alpha} \leq C \cdot\left(\|\omega\|_{n, \alpha}+\|\psi\|_{0, \alpha}+\sum_{i=1}^{l}\left|\psi_{\mid \Gamma_{i}}\right|\right) . \tag{2.13}
\end{equation*}
$$

Similarly we have in Sobolev spaces

$$
\begin{equation*}
\|\psi\|_{H^{2}} \leq C\left(\|\omega\|_{L^{2}}+\sum_{i=1}^{l}\left|\psi_{\mid \Gamma_{i}}\right|\right) . \tag{2.14}
\end{equation*}
$$

Observe that by the trace theorem, for each $i=1, \ldots, l$ we have

$$
\begin{equation*}
\left|\psi_{\mid \Gamma_{i}}\right| \leq C\|\psi\|_{L^{2}(\partial \Omega)} \leq C\|\nabla \psi\|_{L^{2}(\Omega)} \tag{2.15}
\end{equation*}
$$

while an integration by parts gives

$$
\begin{align*}
\int_{\Omega}|\nabla \psi|^{2} & =-\int_{\Omega} \psi \Delta \psi+\sum_{i=1}^{l} \psi_{\mid \Gamma_{i}} \gamma_{i}  \tag{2.16}\\
& \leq \epsilon^{2}\|\psi\|_{L^{2}}^{2}+\frac{1}{4 \epsilon^{2}}\|\omega\|_{L^{2}}^{2}+\epsilon^{2} \sum_{i=1}^{l} \psi_{\mid \Gamma_{i}}^{2}+\frac{1}{4 \epsilon^{2}} \sum_{i=1}^{l} \gamma_{i}^{2} \tag{2.17}
\end{align*}
$$

and therefore

$$
\begin{equation*}
\|\nabla \psi\|_{L^{2}} \leq \epsilon\|\psi\|_{L^{2}}+\frac{C}{\epsilon}\|\omega\|_{L^{2}}+\epsilon \sum_{i=1}^{l}\left|\psi_{\mid \Gamma_{i}}\right|+\frac{C}{\epsilon} \sum_{i=1}\left|\gamma_{i}\right| . \tag{2.18}
\end{equation*}
$$

Taking $\epsilon$ sufficiently small, we can achieve simultaneously

$$
\begin{equation*}
\sum_{i=1}^{l}\left|\psi_{\mid \Gamma_{i}}\right| \leq C\left(\epsilon\|\psi\|_{L^{2}}+\|\omega\|_{L^{2}}+\sum_{i=1}^{l}\left|\gamma_{i}\right|\right) \tag{2.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\|\psi\|_{H^{2}} \leq C\left(\|\omega\|_{L^{2}}+\sum_{i=1}^{l}\left|\gamma_{i}\right|\right) . \tag{2.20}
\end{equation*}
$$

Finally, by Sobolev's embedding (dimension is 2),

$$
\begin{equation*}
\|\psi\|_{0, \alpha} \leq C\|\psi\|_{H^{2}} \leq C\left(\|\omega\|_{L^{2}}+\sum_{i=1}^{l}\left|\gamma_{i}\right|\right) \leq C\left(\|\omega\|_{0, \alpha}+\sum_{i=1}^{l}\left|\gamma_{i}\right|\right) . \tag{2.21}
\end{equation*}
$$

2 Estimates for $\Delta \phi+c \phi=k \quad$ For $c \in C \frac{\infty}{\Omega}$ and $\phi \in \mathcal{U}_{0}$ let $\Delta \phi+c \phi=k$. Writing $\Delta \phi=k-c \phi$, from (2.7), and paying attention that the boundary conditions for $\phi \in \mathcal{U}_{0}$ are those of (2.4), we deduce by (5.114) that

$$
\begin{equation*}
\|\phi\|_{n+2, \alpha} \leq C\left(\|k\|_{n, \alpha}+\|c\|_{0, \alpha}\|\phi\|_{n, \alpha}+\|c\|_{n, \alpha}\|\phi\|_{0, \alpha}\right) . \tag{2.22}
\end{equation*}
$$

Now restricting $c$ to a $\|\cdot\|_{0, \alpha}$-neighborhood and using the interpolation $\|\phi\|_{n, \alpha} \leq \epsilon\|\phi\|_{n+2, \alpha}+$ $C(\epsilon, n)\|\phi\|_{0, \alpha}$, we can choose $\epsilon$ sufficiently small (and independent of $c$ in the $\|\cdot\|_{0, \alpha^{-}}$ neighborhood) to get (2.9).

We say that $c \in C_{\bar{\Omega}}^{\infty}$ is non-degenerate if

$$
\begin{equation*}
E(c) \phi=\Delta \phi+c \phi=0, \quad \phi_{\mid \Gamma_{o}}=0, \left.\quad \frac{\partial \phi}{\partial \tau} \right\rvert\, \Gamma_{i}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0, \quad 1 \leq i \leq l \tag{2.23}
\end{equation*}
$$

has only the trivial solution $\phi \equiv 0$.
Proposition 5 Suppose that $\bar{c} \in C_{\bar{\Omega}}^{\infty}$ is non-degenerate. Then, there exists a $\|\cdot\|_{0, \alpha^{-}}$ neighborhood of $\bar{c}$,

$$
\begin{equation*}
\mathcal{V}_{E}(\bar{c})=\left\{c \in C_{\bar{\Omega}}^{\infty} \mid\|c-\bar{c}\|_{0, \alpha}<\epsilon_{E}\right\} \tag{2.24}
\end{equation*}
$$

such that $E: \mathcal{V}_{E}(\bar{c}) \times \mathcal{U}_{0} \rightarrow C \frac{\infty}{\Omega}$ has a smooth tame family of inverses

$$
V E:\left\{\begin{array}{ccccc}
\mathcal{V}_{E}(\bar{c}) & \times & C_{\bar{\Omega}}^{\infty} & & \mathcal{U}_{0}  \tag{2.25}\\
c & & k & & \phi
\end{array}\right\} .
$$

For $n \geq 0, c \in \mathcal{V}_{E}(\bar{c})$, and $k \in C_{\bar{\Omega}}^{\infty}$,

$$
\begin{equation*}
\|\phi\|_{n+2, \alpha} \leq C\left(\|k\|_{n, \alpha}+\|c\|_{n, \alpha}\|k\|_{0, \alpha}\right) . \tag{2.26}
\end{equation*}
$$

The first derivative with respect to $c$ is given by

$$
\begin{equation*}
D V E(c) \cdot(k, \chi)=V E(c) \cdot(-\chi \phi) . \tag{2.27}
\end{equation*}
$$

## Proof

1 Invertibility for $c$ in a $\|\cdot\|_{0, \alpha}$-neighborhood of $\bar{c}$ Assume $\bar{c}$ is non-degenerate: $\Delta+\bar{c}$, as an operator of Fréchet spaces with the boundary conditions (2.23), has trivial kernel. Let $\phi \in C_{\bar{\Omega}}^{2, \alpha}$ solve

$$
\begin{equation*}
\Delta \phi+\bar{c} \phi=0, \quad \phi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \phi}{\partial \tau}{ }_{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0, \quad 1 \leq i \leq l . \tag{2.28}
\end{equation*}
$$

Then, $\phi \in C^{4, \alpha}$ and, repeating, $\phi \in C_{\bar{\Omega}}^{\infty}$. Thus, $\phi$ is in the kernel of $\Delta+\bar{c}$ as an operator on Fréchet spaces, and by non-degeneracy assumption $\phi \equiv 0$. That is, $\Delta+\bar{c}$ has trivial kernel as an operator of Banach spaces $C^{2, \alpha} \rightarrow C^{0, \alpha}$. The Fredholm alternative then implies that it is in fact an isomorphism of Banach spaces. In particular, we have the estimate

$$
\begin{equation*}
\|\phi\|_{2, \alpha} \leq C\|\Delta \phi+\bar{c} \phi\|_{0, \alpha} \tag{2.29}
\end{equation*}
$$

for all $\phi \in C_{\bar{\Omega}}^{2, \alpha}$ satisfying the boundary conditions (2.23), and in particular when $\phi \in C_{\bar{\Omega}}^{\infty}$.
Let now $c \in C_{\bar{\Omega}}^{\infty}$ :

$$
\begin{align*}
\|\phi\|_{0, \alpha} \leq\|\phi\|_{2, \alpha} & \leq C\|\Delta \phi+\bar{c} \phi\|_{0, \alpha}  \tag{2.30}\\
& \leq C\left(\|\Delta \phi+c \phi\|_{0, \alpha}+\|(c-\bar{c}) \phi\|_{0, \alpha}\right)  \tag{2.31}\\
& \leq C\left(\|\Delta \phi+c \phi\|_{0, \alpha}+\|c-\bar{c}\|_{0, \alpha}\|\phi\|_{0, \alpha}\right) . \tag{2.32}
\end{align*}
$$

Take then $\epsilon_{E}$ sufficiently small in (2.24) so that the last term can be incorporated to the left-hand side for all $\phi \in \mathcal{U}_{0}$ and any $c \in \mathcal{V}_{E}(\bar{c})$ :

$$
\begin{equation*}
\|\phi\|_{0, \alpha} \leq C\|\Delta \phi+c \phi\|_{0, \alpha} \tag{2.33}
\end{equation*}
$$

But from (2.9) we have

$$
\begin{align*}
\|\phi\|_{n+2, \alpha} & \leq C\left(\|\Delta \phi+c \phi\|_{n, \alpha}+\|c\|_{n, \alpha}\|\phi\|_{0, \alpha}\right)  \tag{2.34}\\
& \leq C\left(\|\Delta \phi+c \phi\|_{n, \alpha}+\|c\|_{n, \alpha}\|\Delta \phi+c \phi\|_{0, \alpha}\right) \tag{2.35}
\end{align*}
$$

as desired.

2 Continuity in $c$ and $k \quad$ Let $c, \tilde{c} \in \mathcal{V}_{E}(\bar{c})$ and $k, \tilde{k} \in C_{\bar{\Omega}}^{\infty}$, and let $\Delta \phi+c \phi=k$, $\Delta \tilde{\phi}+\tilde{c} \tilde{\phi}=\tilde{k}$. Then, $\phi-\tilde{\phi}$ solves

$$
\begin{equation*}
(\Delta+\tilde{c}) \cdot(\phi-\tilde{\phi})=(k-\tilde{k})-(c-\tilde{c}) \phi \tag{2.36}
\end{equation*}
$$

and the estimates from the previous paragraph and (5.114) give

$$
\begin{align*}
& \|\phi-\tilde{\phi}\|_{n+2, \alpha}  \tag{2.37}\\
\leq & C \cdot\left(\|k-\tilde{k}\|_{n, \alpha}+\|(c-\tilde{c}) \phi\|_{n, \alpha}+\|\tilde{c}\|_{n, \alpha}\left(\|k-\tilde{k}\|_{0, \alpha}+\|(c-\tilde{c}) \phi\|_{0, \alpha}\right)\right.  \tag{2.38}\\
\leq & C\left(\|k-\tilde{k}\|_{n, \alpha}+\|c-\tilde{c}\|_{n, \alpha}\|\phi\|_{n, \alpha}\right) \tag{2.39}
\end{align*}
$$

where the constant depends on $\|\tilde{c}\|_{n, \alpha}$ only. With $\tilde{k}$ and $\tilde{c}$ fixed, $\|\phi\|_{n, \alpha}$ remains bounded if $\|k\|_{n, \alpha}$ and $\|c\|_{n, \alpha}$ remain bounded. Then, $\|\phi-\tilde{\phi}\|_{n+2, \alpha}$ can be made arbitrarily small provided $\|k-\tilde{k}\|_{n, \alpha}$ and $\|c-\tilde{c}\|_{n, \alpha}$ are taken sufficiently small. That is,

$$
\left\{\begin{array}{ccccc}
\mathcal{V}_{E}(\bar{c})^{n, \alpha} & \times & C_{\bar{\Omega}}^{n, \alpha} & & C_{\bar{\Omega}}^{n+2, \alpha}  \tag{2.40}\\
c & & k & & \phi
\end{array}\right\} \quad(n \geq 0)
$$

is continuous as a map of Banach spaces. Remark that this also implies that

$$
\left.\left\{\begin{array}{cccc}
\mathcal{V}_{E}(\bar{c})^{n} & \times & C \frac{n}{\Omega} & C_{\overline{\bar{\Omega}}}+1  \tag{2.41}\\
c & k
\end{array}\right] \quad \begin{array}{l}
\phi
\end{array}\right\} \quad(n \geq 0)
$$

is continuous as a map of Banach spaces.

3 First derivative of $V E$ Since $V E(c) \cdot k$ is linear in $k$, its derivative in $k$ exists and is simply $V E(c) \cdot k$. For the derivative in $c$, fix then $\chi \in C_{\bar{\Omega}}^{\infty}$ and let $c_{t}=c+t \chi$. Denote the solutions

$$
\begin{equation*}
\Delta \phi_{t}+c_{t} \phi_{t}=k, \quad \Delta \phi+c \phi=k . \tag{2.42}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\Delta \frac{\phi_{t}-\phi}{t}+c \frac{\phi_{t}-\phi}{t}=-\chi \phi_{t}, \quad \text { or } \quad \frac{\phi_{t}-\phi}{t}=V E(c) \cdot\left(-\chi \phi_{t}\right) . \tag{2.43}
\end{equation*}
$$

By continuity of $V E$, the limit of $\frac{\phi_{t}-\phi}{t}$ exists in the $C^{\infty}$-topology. Furthermore, it is given by

$$
\begin{equation*}
D V E(c) \cdot(k, \chi)=V E(c) \cdot(-\chi \phi) . \tag{2.44}
\end{equation*}
$$

which is a continuous function of $c, k, \chi$ as a map of Fréchet spaces. This shows that $V E$ is continuously differentiable as a map of Fréchet spaces. Finally, it is clearly a tame map of $c, k, \chi$ since $V E$ is tame.
$44 V$ is smooth tame $E(c) \cdot \phi=\Delta \phi+c \phi$ is a smooth tame map as the sum of a linear differential operator with constant coefficients and multiplication of functions. Since its inverse $V E(c) \cdot k$ is tame and continuously differentiable with tame first derivative, Theorem 5.3.1, p. 102, Part I, and Theorem 3.1.1, p. 150, Part II of [7] imply that $V E(c) \cdot k$ is a smooth tame map.

### 2.2 The solution operator

We recall that $\Omega$ is not assumed to be diffeomorphic to an annulus, and that the reference steady-state $\bar{F}$ is assumed to satisfy (ND1), i.e. $\bar{c}=-\bar{F}^{\prime}(\bar{\psi})$ is non-degenerate in the sense that (2.23) has only the trivial solution.

Let $I \subset \mathbb{R}$ be a closed, bounded interval strictly larger than range $(\bar{\psi})$.

Proposition 6 Let $\bar{\omega}=\bar{F}(\bar{\psi})$ satisfy (ND1). Then there exists a smooth tame solution operator to (2.2):

$$
S:\left\{\begin{array}{ccc}
\left(\mathcal{V}_{S}(\bar{F}) \subset C_{I}^{\infty}\right) & & C_{\bar{\Omega}}^{\infty}  \tag{2.45}\\
F & \psi
\end{array}\right\}
$$

where $\mathcal{V}_{S}(\bar{F})$ is a $\|\cdot\|_{2, \alpha}$-neighborhood. For $F \in \mathcal{V}_{S}(\bar{F})$,

$$
\begin{equation*}
\|\psi\|_{n+2, \alpha} \leq C\left(\|F\|_{n, \alpha}+1\right), \quad n \geq 0 \tag{2.46}
\end{equation*}
$$

Letting $\mathcal{V}_{S}^{n}(\bar{F})$ be the completion of $\mathcal{V}_{S}(\bar{F})$ in $\|\cdot\|_{n}$, it is continuous as a map

$$
S:\left\{\begin{array}{ccc}
\mathcal{V}_{S}^{n}(\bar{F}) & & C_{\bar{\Omega}}^{n+1}  \tag{2.47}\\
F & & \psi
\end{array}\right\} \quad(n \geq 1)
$$

The first derivative, given by

$$
\begin{equation*}
\phi=D S(F) \cdot f=V E\left(-F^{\prime}(\psi)\right) \cdot(f \circ \psi) \tag{2.48}
\end{equation*}
$$

is continuous as a map

$$
D S:\left\{\begin{array}{ccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & & C \begin{array}{c}
n \\
\Omega \\
F
\end{array}  \tag{2.49}\\
& f & & \phi
\end{array}\right\} \quad(n \geq 1)
$$

For $F \in \mathcal{V}_{S}(\bar{F})$, and any $f \in C_{I}^{\infty}$,

$$
\begin{align*}
\|\phi\|_{n+2, \alpha} & \leq C\left(\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}\right), \quad n \geq 1  \tag{2.50}\\
\|\phi\|_{2, \alpha} & \leq C\|f\|_{0, \alpha} \tag{2.51}
\end{align*}
$$

The second derivative $\phi_{12}=D^{2} S(F)\left(f_{1}, f_{2}\right)$ is continuous as a map

$$
D^{2} S:\left\{\begin{array}{cccccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & \times & C_{I}^{n-1} & & C_{\bar{\Omega}}^{n-1}  \tag{2.52}\\
F & & f_{1} & & & f_{2} & & \\
\phi_{12}
\end{array}\right\} \quad(n \geq 1)
$$

It satisfies the tame estimates, for $F \in \mathcal{V}_{S}(\bar{F})$ and any $f_{1}, f_{2} \in C_{I}^{\infty}$,

$$
\begin{equation*}
\left\|\phi_{12}\right\|_{n+2, \alpha} \leq C\left(\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{1, \alpha}+\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right) \tag{2.53}
\end{equation*}
$$

for $n \geq 1$ and

$$
\begin{equation*}
\left\|\phi_{12}\right\|_{2, \alpha} \leq C\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{1, \alpha} \tag{2.54}
\end{equation*}
$$

## Proof

1 Construction of the solution operator The solution operator will be constructed by first applying the Implicit Function Theorem to suitable Banach spaces. Since ultimately we want a solution operator in the smooth category, several choices for these spaces are possible. We make the following choices. Let $V_{0}(\bar{\psi}) \subset\left\{\psi \in H_{\Omega}^{2}\left|\psi_{\mid \Gamma_{0}}=0, \frac{\partial \psi}{\partial \tau}\right| \Gamma_{i}=\right.$ $\left.0, \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N} d l=\gamma_{i}, 1 \leq i \leq l\right\}$ a neighborhood of $\bar{\psi}$ in which range $(\psi) \subset I$. (Note that this makes sense since $H_{\Omega}^{2}$ continuously embeds into $C_{\bar{\Omega}}^{0}$, and that $\int_{\Gamma_{i}} \frac{\partial \psi}{\partial N} d l$ makes sense by the trace lemma.) We may then define

$$
H:\left\{\begin{array}{ccccc}
C_{I}^{1} & \times & V_{0}(\bar{\psi}) & & L_{\Omega}^{2}  \tag{2.55}\\
F & & \psi & & \Delta \psi-F(\psi)
\end{array}\right\}
$$

This operator is continuously Fréchet-differentiable (for emphasis we work in the Banachcategory). The assumption that $\bar{\omega}=\bar{F}(\bar{\psi})$ satisfies (ND1) implies that

$$
\begin{equation*}
D_{\psi} H(\bar{F}, \bar{\psi}) \phi=\Delta \phi-\bar{F}^{\prime}(\bar{\psi}) \phi \tag{2.56}
\end{equation*}
$$

has trivial kernel and hence is invertible as an operator

$$
\begin{equation*}
\left\{\psi \in H_{\Omega}^{2} \mid \psi_{\mid \Gamma_{0}}=0, \frac{\partial \psi}{\partial \tau}{ }_{\mid \Gamma_{i}}=0, \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N} d l=\gamma_{i}\right\} \longrightarrow L_{\Omega}^{2} \tag{2.57}
\end{equation*}
$$

The classical Implicit Function Theorem guarantees the existence of neighborhoods $V_{S}(\bar{F}) \subset$ $C_{I}^{1}$ of $\bar{F}$ and $V_{S}(\bar{\psi}) \subset V_{0}(\bar{\psi})$ of $\bar{\psi}$, and a continuously Fréchet-differentiable map of Banach spaces $F \in V_{S}(\bar{F}) \mapsto \psi \in V_{S}(\bar{\psi})$ such that $\Delta \psi=F(\psi)$. (We will just say that this is a solution operator $C_{I}^{1} \rightarrow H_{\Omega}^{2}$.) By Sobolev's embedding theorem, this is in fact a solution operator $C_{I}^{1} \rightarrow C_{\bar{\Omega}}^{0, \beta}$ for any $0<\beta<1$. Now composition $(F, \psi) \mapsto F \circ \psi$ is continuous as an operator $C^{1} \times C^{0, \beta} \rightarrow C^{0, \alpha}$ for any $0<\alpha<\beta$ (see Lemma 26). Thus, by elliptic regularity the solution operator is continuous $C_{I}^{1} \rightarrow C_{\bar{\Omega}}^{2, \alpha}$. This along with a simple induction implies that the solution operator is continuous

$$
\left\{\begin{array}{ccc}
C_{I}^{l} & & C_{\bar{\Omega}}^{l+1}  \tag{2.58}\\
F & & \psi
\end{array}\right\} \quad(l \geq 1)
$$

The above implies as well that the solution operator is continuous

$$
\left\{\begin{array}{ccc}
C_{I}^{l, \alpha} & & C_{\bar{\Omega}}^{l+1, \alpha}  \tag{2.59}\\
F & & \psi
\end{array}\right\} \quad(l \geq 1)
$$

A first requirement on $\mathcal{V}_{S}(\bar{F})$ is that it be a $\|\cdot\|_{1, \alpha}$-neighborhood such that

$$
\begin{equation*}
\mathcal{V}_{S}(\bar{F}) \quad \subset \quad\left(V_{S}(\bar{F}) \cap C_{\bar{\Omega}}^{\infty}\right) \tag{2.60}
\end{equation*}
$$

2 Tame estimates on $\psi=S(F)$ in the $C^{n, \alpha}$-grading From Lemma 4, $F \in \mathcal{V}_{S}(\bar{F})$, by (5.121) we have for $n \geq 1$

$$
\begin{equation*}
\|\psi\|_{n+2, \alpha} \leq C \cdot\left(\|F \circ \psi\|_{n, \alpha}+1\right) \leq C \cdot\left(\|F\|_{n, \alpha}+\|\psi\|_{n, \alpha}\|F\|_{1, \alpha}\right), \tag{2.61}
\end{equation*}
$$

and as long as $\psi$ remains in a $\|\cdot\|_{1, \alpha}$-neighborhood. This is the case, e.g. by continuity $F \in C^{1, \alpha} \mapsto \psi \in C^{2, \alpha}$ and since $F \in \mathcal{V}_{S}(\bar{F})$ remains in a $\|\cdot\|_{1, \alpha}$-neighborhood. Recall the inequality $\|\psi\|_{n, \alpha} \leq \epsilon_{1}\|\psi\|_{n+1, \alpha}+C\left(\epsilon_{1}, n\right)\|\psi\|_{0, \alpha}$. Choosing $\epsilon_{1}$ sufficiently small (depending on $n$ but independent of $F \in \mathcal{V}_{S}(\bar{F})$ ) we find

$$
\begin{equation*}
\|\psi\|_{n+2, \alpha} \leq C \cdot\left(\|F\|_{n, \alpha}+\|\psi\|_{0, \alpha}\right) . \tag{2.62}
\end{equation*}
$$

Since $\|\psi\|_{0, \alpha}$ is bounded (because $\|F\|_{1, \alpha}$ is) we obtain the desired tame estimates (2.46) for $n \geq 1$. Finally, since $\|F\|_{1, \alpha}$ remains bounded, so does $\|\psi\|_{0, \alpha}$ and the estimates in fact holds for $n \geq 0$ by increasing the constant if necessary.

3 First derivative Let $F \in \mathcal{V}_{S}(\bar{F})$ and $f \in C_{I}^{\infty}$, and denote $\psi_{t}=S(F+t f)$ and $\psi=S(F)$ the corresponding solutions. Then,

$$
\begin{equation*}
\Delta \frac{\psi_{t}-\psi}{t}-F^{\prime}(\psi) \frac{\psi_{t}-\psi}{t}=\left(\frac{F\left(\psi_{t}\right)-F(\psi)}{t}-F^{\prime}(\psi) \frac{\psi_{t}-\psi}{t}\right)+f\left(\psi_{t}\right) \tag{2.63}
\end{equation*}
$$

Now recall that the solution operator obtained from Step 1 is obtained via the classical Implicit Function Theorem, and therefore is a Fréchet-differentiable map of Banach spaces $F \in C_{I}^{1} \mapsto \psi \in H_{\Omega}^{2}$. Thus, the limit $\phi$ of $\frac{\psi_{t}-\psi}{t}$ exists in $H^{2}$, and it satisfies

$$
\begin{equation*}
\Delta \phi-F^{\prime}(\psi) \phi=f(\psi) \quad \text { in } \quad L^{2} \tag{2.64}
\end{equation*}
$$

By Sobolev's embedding, $\frac{\psi_{t}-\psi}{t} \rightarrow_{t} \phi$ also in $C^{0, \alpha}$, and in turn the right-hand side of (2.63) converges in $C^{0, \alpha}$ as well. This implies that $\frac{\psi_{t}-\psi}{t} \rightarrow_{t} \phi$ in $C^{2, \alpha}$ by continuity of $V E$. Repeating, one can show that $\frac{\psi_{t}-\psi}{t} \rightarrow_{t} \phi$ in $C^{k, \alpha}$ for each $k$, in other words that it converges in the $C^{\infty}$-topology. This proves that the derivative $\phi$ of $\psi=S(F)$ at $F$ in the direction $f$ exists. Furthermore, it is given by $\phi=V E\left(-F^{\prime} \circ \psi\right)(f \circ \psi)$ which is clearly continuous and tame. More precisely, from (2.41), we find that

$$
\left\{\begin{array}{ccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & & C  \tag{2.65}\\
C \frac{n}{\Omega} \\
F & f & & \phi=D S(F) f
\end{array}\right\} \quad(n \geq 1)
$$

is continuous as a map of Banach spaces. (Taking $f \in C_{I}^{n}$ does not improve $\phi$.)

4 Tame estimates on $\phi=D S(F) f$ in the $C^{n, \alpha}$-grading Since we will invoke tame estimates from Lemma 4 to derive tame estimates on $\phi=D S(F) f$, we will take $\epsilon_{S}$ sufficiently small so that

$$
\begin{equation*}
\mathcal{V}_{S}(\bar{F})=\left\{F:\|F-\bar{F}\|_{2, \alpha}<\epsilon_{S}\right\} \subset\left(V_{S}(\bar{F}) \cap C_{\bar{\Omega}}^{\infty}\right) \tag{2.66}
\end{equation*}
$$

and such that $F^{\prime}(\psi) \in \mathcal{V}_{E}(\bar{c})$ whenever $F \in V_{S}(\bar{F})$, where $\bar{c}=-\bar{F}^{\prime}(\bar{\psi})$ and $\mathcal{V}_{E}(\bar{c})$ is as in Lemma 5, By (2.26) and (5.121), we have for $F \in \mathcal{V}_{S}(\bar{F})$ and $n \geq 1$

$$
\begin{equation*}
\|\phi\|_{n+2, \alpha} \leq C\left(\|f(\psi)\|_{n, \alpha}+\|f(\psi)\|_{0, \alpha}\left\|F^{\prime}(\psi)\right\|_{n, \alpha}\right) . \tag{2.67}
\end{equation*}
$$

Now

$$
\begin{align*}
\|f(\psi)\|_{n, \alpha} & \leq C\left(\|f\|_{n, \alpha}+\|\psi\|_{n, \alpha}\|f\|_{1, \alpha}\right)  \tag{2.68}\\
& \leq C\left(\|f\|_{n, \alpha}+\|\psi\|_{n+3, \alpha}\|f\|_{1, \alpha}\right)  \tag{2.69}\\
& \leq C\left(\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}\right) \tag{2.70}
\end{align*}
$$

for $n \geq 1$ while

$$
\begin{equation*}
\|f(\psi)\|_{0, \alpha} \leq C\|f\|_{0, \alpha} \tag{2.71}
\end{equation*}
$$

since $\|\psi\|_{1, \alpha}$ remains bounded (because $\|F\|_{1, \alpha}$ does). Next, for $n \geq 1$

$$
\begin{align*}
\left\|F^{\prime}(\psi)\right\|_{n, \alpha} & \leq C\left(\left\|F^{\prime}\right\|_{n, \alpha}+\|\psi\|_{n, \alpha}\left\|F^{\prime}\right\|_{1, \alpha}\right)  \tag{2.72}\\
& \leq C\left(\|F\|_{n+1, \alpha}+\|\psi\|_{n+3, \alpha}\|F\|_{2, \alpha}\right)  \tag{2.73}\\
& \leq C\|F\|_{n+1, \alpha} \tag{2.74}
\end{align*}
$$

and this holds in fact for $n=0$ as well since $\left\|F^{\prime}(\psi)\right\|_{0, \alpha} \leq C\|F\|_{1, \alpha}$ because $\|\psi\|_{1, \alpha}$ remains bounded. With the above, we have for $n \geq 1$

$$
\begin{equation*}
\|\phi\|_{n+2, \alpha} \leq C\left(\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}\right) \tag{2.75}
\end{equation*}
$$

and for $n=0$

$$
\begin{align*}
\|\phi\|_{2, \alpha} & \leq C\left(\|f(\psi)\|_{0, \alpha}+\left\|F^{\prime}(\psi)\right\|_{0, \alpha}\|f(\psi)\|_{0, \alpha}\right)  \tag{2.76}\\
& \leq C\left(\|f\|_{0, \alpha}+\|F\|_{1, \alpha}\|f\|_{0, \alpha}\right)  \tag{2.77}\\
& \leq C\|f\|_{0, \alpha} \tag{2.78}
\end{align*}
$$

since $\|\psi\|_{1, \alpha}$ and $\|F\|_{1, \alpha}$ remain bounded for $F \in \mathcal{V}_{S}(\bar{F})$.

5 Second derivative Since $\phi$ is a differentiable tame map of $F$ and $f$, it is immediate that $S(F)=\psi$ is twice continuously differentiable and tame (and in fact smooth tame).

Implicit differentiation on $\Delta \phi_{1}=F^{\prime}(\psi) \phi_{1}+f_{1}(\psi)$ shows that the second derivative $\phi_{12}$ in the directions $f_{1}, f_{2}$ satisfies

$$
\begin{equation*}
\Delta \phi_{12}=F^{\prime}(\psi) \phi_{12}+F^{\prime \prime}(\psi) \phi_{1} \phi_{2}+f_{2}^{\prime}(\psi) \phi_{1}+f_{1}^{\prime}(\psi) \phi_{2} \tag{2.79}
\end{equation*}
$$

where $\Delta \phi_{2}=F^{\prime}(\psi) \phi_{2}+f_{2}(\psi)$. From (2.65), we conclude that

$$
D^{2} S:\left\{\begin{array}{ccccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & \times & C_{I}^{n-1} & & \begin{array}{c}
C_{\bar{\Omega}}^{n-1} \\
F
\end{array}  \tag{2.80}\\
& f_{1} & & f_{2} & & & \phi_{12}
\end{array}\right\} \quad(n \geq 1)
$$

is continuous as a map of Banach spaces.

6 Tame estimates on the second derivatives We need first tame estimates on $F^{\prime \prime}(\psi) \phi_{1} \phi_{2}+f_{2}^{\prime}(\psi) \phi_{1}+f_{1}^{\prime}(\psi) \phi_{2}$. Since $F \in \mathcal{V}_{S}(\bar{F})$, we have for $n \geq 1$

$$
\begin{align*}
\left\|F^{\prime \prime}(\psi)\right\|_{n, \alpha} & \leq C\left(\left\|F^{\prime \prime}\right\|_{n, \alpha}+\|\psi\|_{n, \alpha}\left\|F^{\prime \prime}\right\|_{1, \alpha}\right)  \tag{2.81}\\
& \leq C\left(\|F\|_{n+2, \alpha}+\|F\|_{3, \alpha}\right)  \tag{2.82}\\
& \leq C\|F\|_{n+2, \alpha} \tag{2.83}
\end{align*}
$$

and this holds in fact also for $n=0$ since $\|\psi\|_{1, \alpha}$ remains bounded. Next with $i=1,2$,

$$
\begin{equation*}
\left\|f_{i}^{\prime}(\psi)\right\|_{n, \alpha} \leq C\left(\left\|f_{i}\right\|_{n+1, \alpha}+\|F\|_{n, \alpha}\left\|f_{i}\right\|_{2, \alpha}\right) \tag{2.84}
\end{equation*}
$$

for $n \geq 1$ while $\left\|f_{i}^{\prime}(\psi)\right\|_{0, \alpha} \leq C\left\|f_{i}\right\|_{1, \alpha}$ since $\|\psi\|_{1, \alpha}$ remains bounded. Thus, recalling (2.50) and (2.51), we have for $n \geq 3$

$$
\begin{align*}
& \left\|F^{\prime \prime}(\psi) \phi_{1} \phi_{2}+f_{2}^{\prime}(\psi) \phi_{1}+f_{1}^{\prime}(\psi) \phi_{2}\right\|_{n, \alpha}  \tag{2.85}\\
\leq & C\left(\left\|F^{\prime \prime}(\psi)\right\|_{n, \alpha}\left\|\phi_{1}\right\|_{0, \alpha}\left\|\phi_{2}\right\|_{0, \alpha}+\left\|\phi_{1}\right\|_{n, \alpha}\left\|\phi_{2}\right\|_{0, \alpha}+\left\|\phi_{1}\right\|_{0, \alpha}\left\|\phi_{2}\right\|_{n, \alpha}\right.  \tag{2.86}\\
& +\left\|f_{1}^{\prime}(\psi)\right\|_{n, \alpha}\left\|\phi_{2}\right\|_{0, \alpha}+\left\|f_{1}^{\prime}(\psi)\right\|_{0, \alpha}\left\|\phi_{2}\right\|_{n, \alpha}  \tag{2.87}\\
& \left.+\left\|f_{2}^{\prime}(\psi)\right\|_{n, \alpha}\left\|\phi_{1}\right\|_{0, \alpha}+\left\|f_{2}^{\prime}(\psi)\right\|_{0, \alpha}\left\|\phi_{1}\right\|_{n, \alpha}\right)  \tag{2.88}\\
\leq & C\left(\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{0, \alpha}\left\|f_{2}\right\|_{0, \alpha}\right.  \tag{2.89}\\
& +\left(\left\|f_{1}\right\|_{n-2, \alpha}+\|F\|_{n-1, \alpha}\left\|f_{1}\right\|_{1, \alpha}\right)\left\|f_{2}\right\|_{0, \alpha}  \tag{2.90}\\
& +\left\|f_{2}\right\|_{0, \alpha}\left(\left\|f_{2}\right\|_{n-2, \alpha}+\|F\|_{n-1, \alpha}\left\|f_{2}\right\|_{1, \alpha}\right)  \tag{2.91}\\
& +\left(\left\|f_{1}\right\|_{n+1, \alpha}+\|F\|_{n, \alpha}\left\|f_{1}\right\|_{2, \alpha}\right)\left\|f_{2}\right\|_{0, \alpha}  \tag{2.92}\\
& +\left\|f_{1}\right\|_{1, \alpha}\left(\left\|f_{2}\right\|_{n-2, \alpha}+\|F\|_{n-1, \alpha}\left\|f_{2}\right\|_{1, \alpha}\right)  \tag{2.93}\\
& +\left(\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right)\left\|f_{1}\right\|_{0, \alpha}  \tag{2.94}\\
& \left.+\left\|f_{2}\right\|_{1, \alpha}\left(\left\|f_{1}\right\|_{n-2, \alpha}+\|F\|_{n-1, \alpha}\left\|f_{1}\right\|_{1, \alpha}\right)\right)  \tag{2.95}\\
\leq & C\left(\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{1, \alpha}+\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right) \tag{2.96}
\end{align*}
$$

and one can verify that this also holds for $n=1,2$ since $\|\psi\|_{4, \alpha}$ remains bounded for $F \in \mathcal{V}_{S}(\bar{F})$. Finally, for $n=0$, we have

$$
\begin{equation*}
\left\|F^{\prime \prime}(\psi) \phi_{1} \phi_{2}+f_{2}^{\prime}(\psi) \phi_{1}+f_{1}^{\prime}(\psi) \phi_{2}\right\|_{n, \alpha} \leq C\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{1, \alpha} \tag{2.97}
\end{equation*}
$$

The tame estimates on $\phi_{12}$ are thus, for $n \geq 1$,

$$
\begin{align*}
& \left\|\phi_{12}\right\|_{n+2, \alpha}  \tag{2.98}\\
\leq & C\left(\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{1, \alpha}+\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right.  \tag{2.99}\\
& \left.+\|F\|_{n+1, \alpha}\left(\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}+\|F\|_{3, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right)\right)  \tag{2.100}\\
\leq & C\left(\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{1, \alpha}+\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right) \tag{2.101}
\end{align*}
$$

where we have used interpolation to get $\|F\|_{n+1, \alpha}\|F\|_{3, \alpha} \leq C\|F\|_{n+2, \alpha}\|F\|_{2, \alpha} \leq C\|F\|_{n+2, \alpha}$. Finally for $n=0$ we have

$$
\begin{equation*}
\left\|\phi_{12}\right\|_{2, \alpha} \leq C\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{1, \alpha} \tag{2.102}
\end{equation*}
$$

Remarks Gårding's inequality fails in the $C^{n}$-grading. Yet, $\omega=F(\psi)$ is $C^{n}$ if $F$ is $C^{n}$, and furthermore (see Lemma (26) the map

$$
\left\{\begin{array}{ccc}
\mathcal{V}_{S}^{n}(\bar{F}) & & \begin{array}{c}
C_{\bar{\Omega}}^{n} \\
F
\end{array}  \tag{2.103}\\
\omega=F(\psi)
\end{array}\right\} \quad(n \geq 1)
$$

is continuous. Likewise, from $\nu=\Delta \phi$ it would appear that $\nu$ loses two derivatives from $F$. However, writing $\nu=F^{\prime}(\psi) \phi+f(\psi)$, one finds that in fact

$$
\left\{\begin{array}{ccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & & C_{\bar{\Omega}}^{n-1}  \tag{2.104}\\
F & & f & & \nu
\end{array}\right\} \quad(n \geq 1)
$$

is continuous as a map of Banach spaces. Finally, the second derivative $\nu_{12}$ of $F \mapsto \omega$ in the directions $\left(f_{1}, f_{2}\right)$ is continuous

$$
\left\{\begin{array}{cccccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & \times & C_{I}^{n-1} & & C & C_{\bar{\Omega}}^{n-2}  \tag{2.105}\\
F & & f_{1} & & f_{2} & & & \nu_{12}
\end{array}\right\} \quad(n \geq 1)
$$

as a map of Banach spaces.

## 3 Distribution functions and co-adjoint orbits

We are now assuming that $\Omega$ is diffeomorphic to an annulus:

$$
\begin{equation*}
\partial \Omega=\Gamma_{o} \cup \Gamma_{i} . \tag{3.1}
\end{equation*}
$$

This Section is concerned with establishing properties and estimates on distribution functions $A_{\omega}$ for functions $\omega$, which are locally constant on the boundary and have no critical points. These results will apply equally to stream functions, which further satisfy $\psi_{\mid \Gamma_{o}}=0$, and to steady-state vorticity functions, which are locally constant on $\partial \Omega$ as observed in the Introduction. Without loss of generality we will assume that $\omega_{\mid \Gamma_{i}}<\omega_{\mid \Gamma_{o}}$, and therefore we will work with spaces $\mathcal{F}_{+}^{n}$ and $\mathcal{F}_{+}$introduced at the end of Section 17

We begin with the following useful result.
Lemma 7 (Global coordinates on $\bar{\Omega}$ induced by $\omega$ ) For each $\omega \in \mathcal{F}_{+}$there exists a global coordinate system for $\bar{\Omega}, z:[0,1] \times \mathbb{S}^{1} \rightarrow \bar{\Omega}$, such that $\{t\} \times \mathbb{S}^{1}$ is mapped onto the level set $\{x \in \bar{\Omega}: \omega(x)=\min \omega+t(\max \omega-\min \omega)\}$ :

$$
\begin{equation*}
\omega(z(t, s))=\min \omega+t(\max \omega-\min \omega) . \tag{3.2}
\end{equation*}
$$

The map $\omega \mapsto z$ is continuous $\mathcal{F}_{+}^{n} \rightarrow C_{[0,1] \times \mathbb{S}^{1}}^{n-1}, n \geq 2$.
Remark The proof of Lemma 7 can easily be adapted to achieve a continuous map $\mathcal{F}_{+}^{n} \rightarrow C_{[0,1] \times \mathbb{S}^{1}}^{n} \stackrel{[3]}{[3]}$ However, this will make no difference in the rest of the paper as $z$ will always be used in connection with a factor $\frac{1}{|\nabla \omega|}$ which is of class $C^{n-1}$ when $\omega \in \mathcal{F}_{+}^{n}$.

Proof Let $c_{i}(s), 0 \leq s \leq 1$, be a smooth parametrization of $\Gamma_{i}$. For each $s$, let $z=z(t, s)$ be the solution to

$$
\begin{equation*}
\dot{z}=(\max \omega-\min \omega) \frac{\nabla \omega(z)}{|\nabla \omega(z)|^{2}}, \quad z(0)=c_{i}(s) . \tag{3.3}
\end{equation*}
$$

Since $\frac{d}{d t} \omega(z(t))=\max \omega-\min \omega$, we have $\omega(z(t, s))=\min \omega+t(\max \omega-\min \omega)$ for $t \in[0,1]$ and $s \in \mathbb{S}^{1}$. It is standard that if $\omega \in C^{n}, n \geq 2$, then $z$ is $C^{n}$ in $t$ and $C^{n-1}$ in $s$.

For $\omega, \omega_{1}$, denote $f=(\max \omega-\min \omega) \frac{\nabla \omega}{|\nabla \omega|^{2}}, f_{1}=\left(\max \omega_{1}-\min \omega_{1}\right) \frac{\nabla \omega_{1}}{\left|\nabla \omega_{1}\right|^{2}}$, and $z, z_{1}$ the

[^3]corresponding coordinate systems. For $t \in[0,1], s \in \mathbb{S}^{1}$,
\[

$$
\begin{align*}
& z(t, s)-z_{1}(t, s)  \tag{3.4}\\
= & \int_{0}^{t}\left(f(z(\theta, s))-f\left(z_{1}(\theta, s)\right)\right) d \theta+\int_{0}^{t}\left(f-f_{1}\right)\left(z_{1}(\theta, s)\right) d \theta  \tag{3.5}\\
= & \int_{0}^{t} \int_{0}^{1} \nabla f\left(z_{1}(\theta, s)+\tau\left(z(\theta, s)-z_{1}(\theta, s)\right) \cdot\left(z(\theta, s)-z_{1}(\theta, s) d \tau d \theta\right.\right.  \tag{3.6}\\
& +\int_{0}^{t}\left(f-f_{1}\right)\left(z_{1}(\theta, s)\right) d \theta \tag{3.7}
\end{align*}
$$
\]

so that

$$
\begin{equation*}
\left|z(t, s)-z_{1}(t, s)\right| \leq \sup |\nabla f| \int_{0}^{t}\left|z(\theta, s)-z_{1}(\theta, s)\right| d \theta+\sup \left|f-f_{1}\right| . \tag{3.8}
\end{equation*}
$$

This is the same as

$$
\begin{equation*}
\frac{d}{d t}\left(e^{-\sup |\nabla f| t} \int_{0}^{t}\left|z-z_{1}\right|(\theta, s) d \theta\right) \leq e^{-\sup |\nabla f| t} \sup \left|f-f_{1}\right| \tag{3.9}
\end{equation*}
$$

hence integrating

$$
\begin{equation*}
\int_{0}^{t}\left|z-z_{1}\right|(\theta, s) d \theta \leq e^{\sup |\nabla f| t} \int_{0}^{t} e^{-\sup |\nabla f| \theta} \sup \left|f-f_{1}\right| d \theta \leq \sup \left|f-f_{1}\right| \frac{e^{|\sup \nabla f|}}{|\sup \nabla f|} \tag{3.10}
\end{equation*}
$$

In turn,

$$
\begin{equation*}
\left|z(t, s)-z_{1}(t, s)\right| \leq \sup \left|f-f_{1}\right|\left(e^{\sup |\nabla f|}+1\right) . \tag{3.11}
\end{equation*}
$$

Given $f_{1} \in C^{1}$, for any $f \in C^{1}$ in a neighborhood of $f_{1}$ such that $|\nabla f| \leq M$, then sup $\left|z-z_{1}\right|$ can be made arbitrarily small provided $\sup \left|f-f_{1}\right|$ is chosen sufficiently small. In other words, $f \mapsto z$ is continuous as a map $C^{1} \rightarrow C^{0}$.

Estimate now

$$
\begin{align*}
& \left|\dot{z}(t, s)-\dot{z}_{1}(t, s)\right|  \tag{3.12}\\
\leq & \left|f(z(t, s))-f\left(z_{1}(t, s)\right)\right|+\left|\left(f-f_{1}\right)(z(t, s))\right|  \tag{3.13}\\
\leq & \int_{0}^{1}\left|\nabla f\left(z_{1}(t, s)+\tau\left(z(t, s)-z_{1}(t, s)\right)\right)\right|\left|z(t, s)-z_{1}(t, s)\right| d \tau+\sup \left|f-f_{1}\right|  \tag{3.14}\\
\leq & \sup |\nabla f|\left|z-z_{1}\right|+\sup \left|f-f_{1}\right| \tag{3.15}
\end{align*}
$$

which shows that $f \mapsto \dot{z}$ is continuous $C^{1} \rightarrow C^{0}$.
To estimate the derivative $z^{\prime}$ in $s$, differentiate (3.3):

$$
\begin{equation*}
\frac{d}{d t} z^{\prime}(t, s)=\nabla f(z(t, s)) \cdot z^{\prime}(t, s), \quad \frac{d}{d t} z_{1}^{\prime}(t, s)=\nabla f_{1}\left(z_{1}(t, s)\right) \cdot z_{1}^{\prime}(t, s) \tag{3.16}
\end{equation*}
$$

and observe that $z^{\prime}(0, s)=z_{1}^{\prime}(0, s), s \in \mathbb{S}^{1}$. Integrating,

$$
\begin{align*}
z^{\prime}(t, s)-z_{1}^{\prime}(t, s)= & \int_{0}^{t} \nabla f\left(z_{1}(\theta, s)\right)\left(z^{\prime}(\theta, s)-z_{1}^{\prime}(\theta, s)\right) d \theta  \tag{3.17}\\
& +\int_{0}^{t}\left(\nabla f(z(\theta, s))-\nabla f_{1}\left(z_{1}(\theta, s)\right)\right) z_{1}^{\prime}(\theta, s) d \theta \tag{3.18}
\end{align*}
$$

and thus

$$
\begin{equation*}
\left|z^{\prime}(t, s)-z_{1}^{\prime}(t, s)\right| \leq \sup |\nabla f| \int_{0}^{t}\left|z^{\prime}(\theta, s)-z_{1}^{\prime}(\theta, s)\right| d \theta+\sup \left|z_{1}^{\prime}\right| \sup \left|\nabla\left(f-f_{1}\right)\right| . \tag{3.19}
\end{equation*}
$$

As above, we find

$$
\begin{equation*}
\left|z^{\prime}(t, s)-z_{1}^{\prime}(t, s)\right| \leq \sup \left|z_{1}^{\prime}\right| \sup \left|\nabla\left(f-f_{1}\right)\right|\left(\frac{e^{\sup |\nabla f|}}{\sup |\nabla f|}+1\right) \tag{3.20}
\end{equation*}
$$

showing that $f \mapsto z^{\prime}$ is continuous $C^{1} \rightarrow C^{0}$.
Therefore, $f \mapsto z$ is continuous $C^{1} \rightarrow C^{1}$. Differentiating further in $t$ and $s$, one establishes easily by induction that $f \mapsto z$ is continuous $C^{k} \rightarrow C^{k}$ for each $k \geq 1$.

### 3.1 Properties of the distribution function $A_{\omega}$

Linearizing $\omega \circ \eta$ at $\eta=\mathrm{id}$, a tangent $\nu$ to $\mathcal{O}_{\omega}$ at $\omega$ is of the form

$$
\begin{equation*}
\nu=\{\omega, \alpha\} \quad \text { for some stream function } \alpha \in \mathcal{U} . \tag{3.21}
\end{equation*}
$$

This is a first order PDE in $\alpha$ hence can be locally integrated along the characteristics, which are here closed curves. The compatibility conditions (to be able to integrate round these closed curves) are

$$
\begin{equation*}
\int_{\{x: \omega(x)=\lambda\}} \frac{\nu(x)}{|\nabla \omega(x)|} d l(x)=0, \quad \min \omega \leq \lambda \leq \max \omega . \tag{3.22}
\end{equation*}
$$

On the other hand, if $\omega \in \mathcal{F}_{+}$, then the distribution function $A_{\omega}(\lambda)=|\{x: \omega(x)<\lambda\}|$ can be expressed as

$$
\begin{equation*}
A_{\omega}(\lambda)=\int_{\min \omega}^{\lambda} \int_{x: \omega(x)=\lambda} \frac{1}{|\nabla \omega(x)|} d l(x) d \lambda, \quad \lambda \in[\min \omega, \max \omega] . \tag{3.23}
\end{equation*}
$$

This is easily seen using the coarea formula (see § 3.2 in [6])

$$
\begin{equation*}
\int_{\Omega} u(x)|\nabla \omega(x)| \zeta(\omega(x)) d x=\int_{\min \omega}^{\max \omega} \zeta\left(\lambda^{\prime}\right) \int_{x: \omega(x)=\lambda^{\prime}} u(x) d l(x) d \lambda^{\prime} \tag{3.24}
\end{equation*}
$$

with $u(x)=\frac{1}{|\nabla \omega(x)|}$ and $\zeta\left(\lambda^{\prime}\right)$ the characteristic function over the interval $[\min \omega, \lambda]$. Much of the work will be devoted to the detailed study of

$$
\begin{equation*}
J_{\omega} u(\lambda)=\int_{x: \omega(x)=\lambda} u(x) d l(x), \quad \lambda \in[\min \omega, \max \omega] . \tag{3.25}
\end{equation*}
$$

With this notation, the coarea formula reads

$$
\begin{equation*}
\int_{\Omega} u(x)|\nabla \omega(x)| \zeta(\omega(x)) d x=\int_{\min \omega}^{\max \omega} \zeta(\lambda) J_{\omega} u(\lambda) d \lambda . \tag{3.26}
\end{equation*}
$$

The derivative of $A_{\omega}$ can then be written as

$$
\begin{equation*}
A_{\omega}^{\prime}=J_{\omega} \frac{1}{|\nabla \omega|} \tag{3.27}
\end{equation*}
$$

and, differentiating $A_{\omega}^{-1}\left(A_{\omega}(\lambda)\right)=\lambda$, that of $A_{\omega}^{-1}$ is

$$
\begin{equation*}
\left(A_{\omega}^{-1}\right)^{\prime}=\frac{1}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}} . \tag{3.28}
\end{equation*}
$$

## Proposition 8 ( $A_{\omega}$ characterizes $\mathcal{O}(\omega)$ locally)

1. For $\omega, \omega_{1} \in \mathcal{F}_{+}$, let $z, z_{1}$ be the global coordinate transformations from Lemma 7 . If $\omega, \omega_{1} \in \mathcal{F}_{+}$and $A_{\omega_{1}}=A_{\omega}$, then there exists $\eta \in \mathcal{D}_{\text {vol }}$ such that $\omega=\omega_{1} \circ \eta$.
2. Let $\omega \in \mathcal{F}_{+}$and let $\nu \in \mathcal{F}$ such that $\nu_{\mid \partial \Omega}=0$. Then, there exists a stream function $\alpha \in \mathcal{U}$ such that

$$
\begin{equation*}
\nu=\{\omega, \alpha\} \tag{3.29}
\end{equation*}
$$

if and only if

$$
\begin{equation*}
\int_{\omega=\lambda} \frac{\nu}{|\nabla \omega|} d l=0, \quad \lambda \in \operatorname{range}(\omega) . \tag{3.30}
\end{equation*}
$$

3. Let $\omega_{\epsilon} \in \mathcal{F}_{+}, \omega_{0}=\omega$. Then, $\frac{d A_{\omega_{\epsilon}}}{d \epsilon}{ }_{\mid \epsilon=0}=0$, if and only if $\nu=\left.\frac{d \omega_{\epsilon}}{d \epsilon}\right|_{\epsilon \in 0}$ is tangent to the orbit $\mathcal{O}(\omega)$ at $\omega$, i.e. $\nu=\{\omega, \alpha\}$ for some $\alpha \in \mathcal{U}$.

## Proof

$1 A_{\omega}$ characterizes $\mathcal{O}(\omega)$ locally If $A_{\omega_{1}}=A_{\omega}$, then setting $\phi=z_{1} \circ z^{-1}$, we have $\omega=\omega_{1} \circ \phi$. Next we construct a diffeomorphism $\psi$ which moves points along the level sets of $\omega_{1}$ in such a way that $\eta=\psi \circ \phi \in \mathcal{D}_{\text {vol }}$. We write it in the form

$$
\begin{equation*}
\psi(x)=z_{1}(t, a(t, s)), \quad \text { where } \quad x=z_{1}(t, s) \tag{3.31}
\end{equation*}
$$

or $\psi=z_{1} \circ \alpha$, where $\alpha(t, s)=(t, a(t, s))$. The condition $\operatorname{det}\left(\frac{\partial(\psi \circ \phi)}{\partial(x, y)}\right)=1$ can be written in the form $(Z \circ \alpha) \partial_{s} a=F$ where $Z=\operatorname{det}\left(\frac{\partial z_{1}}{\partial(t, s)}\right)$ and $F=\frac{1}{\operatorname{det}\left(\frac{\partial \phi}{\partial(x, y)}\right) \circ \phi^{-1}}$. This is a collection
of ODEs in $s$ parametrized by $t$. Let then $a(t, s),(t, s) \in[0,1] \times[0,1]$, denote the solution such that $a(t, 0)=0,0 \leq t \leq 1$. Setting $\eta=\psi \circ \phi$ defines a local diffeomorphism $\bar{\Omega} \rightarrow \bar{\Omega}$ such that $\operatorname{det}\left(\frac{\partial \eta}{\partial(x, y)}\right)=1$. One needs to check that $\eta$ is a global diffeomorphism of $\bar{\Omega}$.

Fix $\lambda \in \operatorname{range}(\omega)$ so that $z\left(\frac{\lambda-\min \omega}{\max \omega-\min \omega}, 0\right)$ is a point on the level set $\{\omega=\lambda\}$, and denote $q=q(t)$ the solution to

$$
\begin{equation*}
\dot{q}=\nabla^{\perp} \omega(q), \quad q(0)=z\left(\frac{\lambda-\min \omega}{\max \omega-\min \omega}, 0\right) \tag{3.32}
\end{equation*}
$$

Clearly, $q(t)$ travels around the level set $\{\omega=\lambda\}$. This is a Hamiltonian flow with respect to the symplectic form dvol $=d x \wedge d y$. Now $\eta$ is a local symplectomorphism, so that $q_{1}=\eta \circ q$ solves

$$
\begin{equation*}
\dot{q}_{1}=\nabla^{\perp} \omega_{1}\left(q_{1}\right), \quad q_{1}(0)=\eta\left(z\left(\frac{\lambda-\min \omega}{\max \omega-\min \omega}, 0\right)\right) . \tag{3.33}
\end{equation*}
$$

Clearly, $q_{1}(t)$ travels around the level set $\left\{\omega_{1}=\lambda\right\}$. But $A_{\omega}=A_{\omega_{1}}$, so that the travel time of $q(t)$ and $q_{1}(t)$ around $\{\omega=\lambda\}$ and $\left\{\omega_{1}=\lambda\right\}$ respectively are the same:

$$
\begin{equation*}
\int_{\omega_{1}=\lambda} \frac{d l}{\left|\nabla \omega_{1}\right|}=\frac{d}{d \lambda} A_{\omega_{1}}(\lambda)=\frac{d}{d \lambda} A_{\omega}(\lambda)=\int_{\omega=\lambda} \frac{d l}{|\nabla \omega|} . \tag{3.34}
\end{equation*}
$$

Finally, $q_{1}=\eta \circ q$ so $\eta$ takes $\{\omega=\lambda\}$ onto $\left\{\omega_{1}=\lambda\right\}$ for each $\lambda$ and the claim is proved.

2 Characterizing $\nu \| \mathcal{O}(\omega)$ Equation (3.29) is a first order PDE in $\alpha$, and the characteristics are the level sets of $\omega$. Thus, it is locally solvable and (3.30) are precisely the compatibility conditions that ensure that $\alpha$ is globally defined (the characteristics are simple closed curves).

3 Characterizing $\nu \| \mathcal{O}(\omega)$ in terms of distribution functions We will show that (3.30) holds. This is immediate once the identity

$$
\begin{equation*}
\frac{\partial}{\partial \epsilon}{ }_{\mid \epsilon=0} A_{\omega_{\epsilon}}(\lambda)=-J_{\omega} \frac{\nu}{|\nabla \omega|}(\lambda)=-\int_{\omega=\lambda} \frac{\nu}{|\nabla \omega|} d l \tag{3.35}
\end{equation*}
$$

is proved. Let $\lambda$ be in the interior of range $(\omega)$ and let $f$ be a $C^{1}$ function of $\mathbb{R}$ such that $f^{\prime}$ has support contained in the interior of range $(\omega)$. Then, by the coarea formula (3.26)

$$
\begin{equation*}
\frac{\partial}{\partial}_{\mid \epsilon=0} \int_{\Omega} f\left(\omega_{\epsilon}(x)\right) d x=\int_{\Omega} f^{\prime}(\omega(x)) \nu(x) d x=\int_{\mathbb{R}} f^{\prime}(\lambda) J_{\omega} \frac{\nu}{|\nabla \omega|}(\lambda) d \lambda . \tag{3.36}
\end{equation*}
$$

By approximation, this holds for $f$ a continuous, piecewise linear function. Fix then $\lambda$ and for $\delta>0$ let $f^{\delta}\left(\lambda^{\prime}\right)$ have value 1 for $\lambda^{\prime}<\lambda, 0$ for $\lambda^{\prime}>\lambda+\delta$, and linear in between. Then $\left.\frac{\partial}{\partial \epsilon}\right|_{\epsilon=0} \int_{\Omega} f^{\delta}\left(\omega_{\epsilon}(x)\right) d x \rightarrow_{\delta} \frac{\partial}{\partial \partial \epsilon}{ }_{\mid \epsilon=0} A_{\omega_{\epsilon}}(\lambda)$ while $\int_{\mathbb{R}}\left(f^{\delta}\right)^{\prime}\left(\lambda^{\prime}\right) J_{\omega} \frac{\nu}{|\nabla \omega|}\left(\lambda^{\prime}\right) d \lambda^{\prime} \rightarrow_{\delta}-\left(J_{\omega} u\right)(\lambda)$.

Lemma 9 Setting $N=\frac{\nabla \omega}{|\nabla \omega|}$,

$$
\begin{equation*}
\frac{d}{d \lambda}\left(J_{\omega} u\right)(\lambda)=J_{\omega}\left(\frac{\operatorname{div}(u N)}{|\nabla \omega|}\right)(\lambda), \quad \lambda \in[\min \omega, \max \omega] \tag{3.37}
\end{equation*}
$$

Proof Let $\zeta$ be an arbitrary smooth function with compact support in $(\min \omega, \max \omega)$. Integrating by parts the coarea formula (3.26), using the identities $\nabla(\zeta \circ \omega)=\zeta^{\prime}(\omega) \nabla \omega$ and $\operatorname{div}((\zeta \circ \omega) u N)=(\zeta \circ \omega) \operatorname{div}(u N)+\langle\nabla(\zeta \circ \omega), u N\rangle$, and the coarea formula again, we have

$$
\begin{align*}
\int_{\min \omega}^{\max \omega} \zeta(\lambda)\left(J_{\omega} u\right)^{\prime}(\lambda) d \lambda & =-\int_{\Omega} u(x)|\nabla \omega(x)| \zeta^{\prime}(\omega(x)) d x  \tag{3.38}\\
& =\int_{\Omega} \operatorname{div}(u N)(\zeta \circ \omega) d x  \tag{3.39}\\
& =\int_{\min \omega}^{\max \omega} \zeta(\lambda) J_{\omega}\left(\frac{\operatorname{div}(u N)}{|\nabla \omega|}\right)(\lambda) d \lambda . \tag{3.40}
\end{align*}
$$

Lemma 10 Let $\omega_{\epsilon} \in \mathcal{F}$ and $u \in C \frac{\infty}{\Omega}$. Then we have the pointwise derivative

$$
\begin{equation*}
\frac{\partial}{\partial \epsilon}{ }_{\mid \epsilon=0} J_{\omega_{\epsilon}} u(\lambda)=-J_{\omega}\left\{\frac{\nu \operatorname{div}(u N)}{|\nabla \omega|}\right\}(\lambda), \quad \text { where } \quad \nu:=\frac{\partial}{\partial \epsilon}{ }_{\mid \epsilon=0} \omega_{\epsilon} . \tag{3.41}
\end{equation*}
$$

Proof For an arbitrary function $\zeta$ with compact support in $(\min \omega, \max \omega)$, differentiating the co-area formula (3.26)

$$
\begin{equation*}
\int_{\Omega} u(x)\left|\nabla \omega_{\epsilon}(x)\right| \zeta\left(\omega_{\epsilon}(x)\right) d x=\int_{\min \omega}^{\max \omega} \zeta(\lambda) J_{\omega_{\epsilon}} u(\lambda) d \lambda \tag{3.42}
\end{equation*}
$$

at $\epsilon=0$ we obtain

$$
\begin{equation*}
\int_{\Omega} u(x)\left[\frac{\langle\nabla \omega, \nabla \nu\rangle}{|\nabla \omega|} \zeta(\omega(x))+|\nabla \omega| \zeta^{\prime}(\omega(x)) \nu(x)\right] d x=\left[\int_{\min \omega}^{\max \omega} \zeta(\lambda)\left(\frac{\partial}{\partial \epsilon}{ }_{\mid \epsilon=0} J_{\omega_{\epsilon}} u\right)(\lambda) d \lambda\right] . \tag{3.43}
\end{equation*}
$$

On the one hand, the coarea formula (3.26) and an integration by parts give

$$
\begin{equation*}
\int_{\Omega} u(x) \nu(x)|\nabla \omega| \zeta^{\prime}(\omega(x)) d x=\int_{\min \omega}^{\max \omega} \zeta^{\prime}(\lambda)\left(J_{\omega}(u \nu)\right)(\lambda) d \lambda=-\int_{\min \omega}^{\max \omega} \zeta(\lambda)\left(J_{\omega}(u \nu)\right)^{\prime}(\lambda) d \lambda \tag{3.44}
\end{equation*}
$$

and on the other,

$$
\begin{equation*}
\int_{\Omega} u(x) \frac{\langle\nabla \omega, \nabla \nu\rangle}{|\nabla \omega|} \zeta(\omega(x)) d x=\int_{\min \omega}^{\max \omega} \zeta(\lambda) J_{\omega}\left(u \frac{\langle\nabla \omega, \nabla \nu\rangle}{|\nabla \omega|^{2}}\right)(\lambda) d \lambda \tag{3.45}
\end{equation*}
$$

Since $\zeta$ is arbitrary, and using (3.37), we conclude with

$$
\begin{equation*}
\frac{\partial}{\partial \epsilon}_{\mid \epsilon=0}\left(J_{\omega_{\epsilon}} u(\lambda)\right)=-\left(J_{\omega}(u \nu)\right)^{\prime}(\lambda)+J_{\omega}\left(u \frac{\langle\nabla \omega, \nabla \nu\rangle}{|\nabla \omega|^{2}}\right)(\lambda)=-J_{\omega}\left\{\frac{\nu \operatorname{div}(u N)}{|\nabla \omega|}\right\}(\lambda) . \tag{3.46}
\end{equation*}
$$

## $3.2 \quad A_{\omega}^{-1}$ is a smooth tame map of $\omega$

We will heavily rely on continuity and tameness of operators introduced in the Appendix. Also, we recall the Faà di Bruno formula: if $f, g$ are two functions of one variable, then the $n$-th derivative $(n \geq 0)$ of $f \circ g$ is

$$
\begin{equation*}
(f \circ g)^{(n)}=\sum_{k=1}^{n} \sum_{\substack{j_{1}+\ldots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{1}, \ldots, j_{k}}\left(f^{(k)} \circ g\right) g^{\left(j_{1}\right)} \cdots g^{\left(j_{k}\right)}, \quad n \geq 0 \tag{3.47}
\end{equation*}
$$

where $c_{k ; j_{1}, \ldots, j_{k}}$ are constants. If $f=g^{-1}$, then $f^{\prime}=\frac{1}{g^{\circ} \circ f}$ and $(f \circ g)^{(n)}=0$ for $n \geq 2$, hence

$$
\begin{equation*}
f^{(n)}=-\left(f^{\prime}\right)^{n} \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{1}, \ldots, j_{k}} f^{(k)}\left(g^{\left(j_{1}\right)} \circ f\right) \cdots\left(g^{\left(j_{k}\right)} \circ f\right), \quad n \geq 2 . \tag{3.48}
\end{equation*}
$$

Proposition $11\left(Q(\omega)=A_{\omega}^{-1}\right.$ is smooth tame)

1. The operator

$$
Q:\left\{\begin{array}{c}
\mathcal{F}_{+}  \tag{3.49}\\
\omega
\end{array} \longrightarrow \begin{array}{c}
C_{[0,|\Omega|]}^{\infty} \\
A_{\omega}^{-1}
\end{array}\right\}
$$

is a smooth tame map of Fréchet spaces with first derivative

$$
\begin{equation*}
D Q(\omega) \cdot \nu=\frac{J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}} . \tag{3.50}
\end{equation*}
$$

For $n \geq 0, \omega \in \mathcal{F}_{+}$in $a\|\cdot\|_{2, \alpha}$-neighborhood and any $\nu, \nu_{1}, \nu_{2} \in \mathcal{F}$,

$$
\begin{align*}
\left\|A_{\omega}^{-1}\right\|_{n, \alpha} \leq & C\left(\|\omega\|_{n, \alpha}+1\right)  \tag{3.51}\\
\|D Q(\omega) \nu\|_{n, \alpha} \leq & C\left(\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right)  \tag{3.52}\\
\left\|D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)\right\|_{n, \alpha} \leq & C\left\{\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}\right.  \tag{3.53}\\
& \left.+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{2, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\right\} . \tag{3.54}
\end{align*}
$$

2. The operator

$$
\left\{\begin{array}{ccccc}
\mathcal{F}_{+} & \times & C_{\bar{\Omega}}^{\infty} & & \begin{array}{c}
C_{[0,|\Omega|]}^{\infty} \\
\omega
\end{array}  \tag{3.55}\\
u & & J_{\omega} u \circ A_{\omega}^{-1}
\end{array}\right\}
$$

is smooth tame. For $n \geq 0, \omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and any $u$ (without restriction),

$$
\begin{equation*}
\left\|J_{\omega} \frac{u}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha} \leq C \cdot\left(\|u\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|u\|_{1, \alpha}\right) . \tag{3.56}
\end{equation*}
$$

What prevents one from working directly with $A_{\omega}$ is that the range of $\omega$ is not fixed and that $A_{\omega}^{\prime}$ is necessarily discontinuous at the endpoints of range $(\omega)$. This problem is resolved by working instead with the inverse distribution function $A_{\omega}^{-1}$, at the cost of a fair amount of technical complications.

Proposition 11 is split into Lemmas 12, 13, and 14, For continuity, the $C^{n}$-grading will be more convenient. However, tame estimates will still be derived in the $C^{n, \alpha}$-grading.

Lemma $12\left(Q(\omega)=A_{\omega}^{-1}\right.$ is continuous) $Q(\omega)=A_{\omega}^{-1}$ is continuous as a map of Fréchet spaces. More precisely, the following are continuous as maps of Banach spaces:

$$
\begin{gather*}
Q:\left\{\begin{array}{clc}
\mathcal{F}_{+}^{m} & \longrightarrow \begin{array}{c}
C_{[0,|\Omega|]}^{m} \\
\omega
\end{array} & A_{\omega}^{-1}
\end{array}\right\} \quad(m \geq 2)  \tag{3.57}\\
\left\{\begin{array}{ccc}
\mathcal{F}_{+}^{m} & \times C_{\bar{\Omega}}^{m-1} & \longrightarrow \\
\omega & u & \\
\left.\left(J_{\omega} u\right) \circ,|\Omega|\right] \\
\omega & A_{\omega}^{-1}
\end{array}\right\} \quad(m \geq 2) \tag{3.58}
\end{gather*}
$$

Proof
$1 \omega \mapsto A_{\omega}^{-1}$ is continuous $C^{m} \rightarrow C^{m}, m \geq 2 \quad$ In order to alleviate some of the complications of working with $A_{\omega}^{-1}$ (rather than $A_{\omega}$ directly), we will use the following device. Fix $\omega_{1} \in \mathcal{F}_{+}$. For $\omega \in \mathcal{F}_{+}$, let $\omega_{2} \in \mathcal{F}_{+}$and $a$ a monotone increasing function such that range $\left(\omega_{2}\right)=\operatorname{range}\left(\omega_{1}\right)$ and $\omega=a \circ \omega_{2}$. We will take $a$ to be affine:

$$
\begin{equation*}
\omega=a \circ \omega_{2}=\min \omega+\frac{\max \omega-\min \omega}{\max \omega_{2}-\min \omega_{2}}\left(\omega_{2}-\min \omega_{2}\right) . \tag{3.59}
\end{equation*}
$$

Fix $k \geq 0$. Then

$$
\begin{align*}
\left\|A_{\omega}^{-1}-A_{\omega_{1}}^{-1}\right\|_{k} & \leq\left\|A_{a \circ \omega_{2}}^{-1}-A_{\omega_{2}}^{-1}\right\|_{k}+\left\|A_{\omega_{2}}^{-1}-A_{\omega_{1}}^{-1}\right\|_{k}  \tag{3.60}\\
& =\left\|(a-\mathrm{id}) \circ A_{\omega_{2}}^{-1}\right\|_{k}+\left\|A_{\omega_{2}}^{-1}-A_{\omega_{1}}^{-1}\right\|_{k}  \tag{3.61}\\
& =I_{k}+I I_{k} . \tag{3.62}
\end{align*}
$$

2 Estimating $I_{k}$ In view of (3.47), $I_{k}$ is arbitrarily small provided $\|a-\mathrm{id}\|_{k}$ is sufficiently small while $\left\|A_{\omega_{2}}^{-1}\right\|_{k}$ remains bounded. Since $a$ is affine, the former is small provided $\left\|\omega-\omega_{1}\right\|_{0}$ is small. Next, we show that $\left\|A_{\omega_{2}}^{-1}\right\|_{k}$ remains bounded provided $\left\|\omega_{2}\right\|_{k}$ is bounded and $\left\|\omega_{2}-\omega_{1}\right\|_{2}$ is sufficiently small.

Observe first that $\left\|A_{\omega_{2}}^{-1}\right\|_{0}=\max \left(\left|\max \omega_{2}\right|,\left|\min \omega_{2}\right|\right)=\max (|\max \omega|,|\min \omega|)$ which is fixed. For $\lambda \in \operatorname{range}\left(\omega_{2}\right)$,

$$
\begin{equation*}
A_{\omega_{2}}^{\prime}(\lambda)=\int_{\omega_{2}=\lambda} \frac{d l}{\left|\nabla \omega_{2}\right|}=\int_{s \in \mathbb{S}^{1}} \frac{\left|\frac{\partial z_{2}}{\partial s}\right|\left(\frac{\lambda-\min \omega_{2}}{\max \omega_{2}-\min \omega_{2}}, s\right)}{\left|\nabla \omega_{2}\left(z_{2}\left(\frac{\lambda-\min \omega_{2}}{\max \omega_{2}-\min \omega_{2}}, s\right)\right)\right|} d s . \tag{3.63}
\end{equation*}
$$

Fix then $L_{1}$ and a neighborhood

$$
\begin{equation*}
\left\|\omega_{2}-\omega_{1}\right\|_{2}<\epsilon_{0} \tag{3.64}
\end{equation*}
$$

where

$$
\begin{equation*}
\frac{1}{L_{1}}<\int_{\omega_{2}=\lambda} d l<L_{1} \quad \text { and } \quad \frac{1}{L_{1}}<A_{\omega_{2}}^{\prime}(\lambda)<L_{1} \tag{3.65}
\end{equation*}
$$

for $\lambda \in \operatorname{range}\left(\omega_{2}\right)$. Now

$$
\begin{equation*}
\left(A_{\omega_{2}}^{-1}\right)^{\prime}(\mu)=\frac{1}{A_{\omega_{2}}^{\prime} \circ A_{\omega_{2}}^{-1}(\mu)}, \quad \mu \in[0,|\Omega|] \tag{3.66}
\end{equation*}
$$

so that $\left\|A_{\omega_{2}}^{-1}\right\|_{1}$ remains bounded when $\left\|\omega_{2}-\omega_{1}\right\|_{2}<\epsilon_{0}$.
For the higher order derivatives, we use (3.48): $\frac{d^{n} A_{\omega_{2}}^{-1}}{d \mu^{n}}, n \geq 2$, is given by

$$
\begin{equation*}
-\left(\frac{d A_{\omega_{2}}^{-1}}{d \mu}\right)^{n} \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{i}}\left(\frac{d^{k} A_{\omega_{2}}^{-1}}{d \mu^{k}}\right)\left(\frac{d^{j_{1}} A_{\omega_{2}}}{d \lambda^{j_{1}}} \circ A_{\omega_{2}}^{-1}\right) \cdots\left(\frac{d^{j_{k}} A_{\omega_{2}}}{d \lambda^{j_{k}}} \circ A_{\omega_{2}}^{-1}\right) . \tag{3.67}
\end{equation*}
$$

An induction shows that $\frac{d^{n} A_{\omega_{2}}^{-1}}{d \mu^{n}}$ is bounded provided $\frac{d A_{\omega_{2}}}{d \mu}, \ldots, \frac{d^{n} A_{\omega_{2}}}{d \mu^{n}}$. We show that this holds when $\left\|\omega_{2}\right\|_{n}$ is bounded. Setting

$$
\begin{equation*}
v_{0}=\frac{1}{\left|\nabla \omega_{2}\right|}, \quad v_{m}=\frac{\operatorname{div}\left(v_{m-1} N_{2}\right)}{\left|\nabla \omega_{2}\right|}, \quad m \geq 1 \tag{3.68}
\end{equation*}
$$

where $N_{2}=\frac{\nabla \omega_{2}}{\left|\nabla \omega_{2}\right|}$, Lemma 9 yields

$$
\begin{align*}
& \frac{d^{j} A_{\omega_{2}}(\lambda)}{d \lambda^{j}}  \tag{3.69}\\
= & \left(J_{\omega_{2}} v_{j-1}\right)(\lambda)=\int_{\omega_{2}=\lambda} v_{j-1} d l  \tag{3.70}\\
= & \int_{s \in \mathbb{S}^{1}} v_{j-1}\left(z\left(\left(\frac{\lambda-\min \omega_{2}}{\max \omega_{2}-\min \omega_{2}}, s\right)\right)\left|\frac{\partial z}{\partial s}\right|\left(\frac{\lambda-\min \omega_{2}}{\max \omega_{2}-\min \omega_{2}}, s\right) d s\right. \tag{3.71}
\end{align*}
$$

But $v_{j-1}$ is a smooth expression of the derivatives of $\omega_{2}$ up to order $j$. Therefore, with $\left\|\omega_{2}-\omega_{1}\right\|_{2}<\epsilon_{0}$ and $\left\|\omega_{2}\right\|_{j}$ bounded, $\left\|A_{\omega_{2}}\right\|_{C_{\text {range }\left(\omega_{2}\right)}^{j}}$ is bounded.

5 Estimating $I I_{k} \quad$ Since $\omega_{2}$ and $\omega_{1}$ have same range, we may invoke general results on the inversion operator, see Lemma 27 in the Appendix: for $k \geq 1, I I_{k}$ is arbitrarily small provided $\left\|A_{\omega_{2}}-A_{\omega_{1}}\right\|_{C_{\left.\text {range( } \omega_{1}\right)}^{k}}$ is taken sufficiently small. Restricting to functions such that range $\left(\omega_{2}\right)=\operatorname{range}\left(\omega_{1}\right)$, it remains to show that $\omega_{2} \mapsto A_{\omega_{2}}$ is continuous $C^{k} \rightarrow C^{k}, k \geq 2$. But this is immediate in view of (3.71) and the fact that $\omega \mapsto \frac{\partial z}{\partial s}$ is continuous $C^{j} \rightarrow C^{0}$ for $j \geq 2$.
$4(\omega, u) \mapsto\left(J_{\omega} u\right) \circ A_{\omega}^{-1}$ is continuous $C^{m} \times C^{m-1} \rightarrow C^{m-1}, m \geq 2 \quad$ Let $u$, $u_{1}$, and $\omega, \omega_{1}$ with corresponding coordinate systems $z, z_{1}$. For $\mu \in[0,|\Omega|]$, write $\lambda=A_{\omega}^{-1}(\mu)$, $\lambda_{1}=A_{\omega_{1}}^{-1}(\mu)$, and $t=\frac{\lambda-\min \omega}{\max \omega-\min \omega}, t_{1}=\frac{\lambda_{1}-\min \omega_{1}}{\max \omega_{1}-\min \omega_{1}}$. Then

$$
\begin{align*}
& J_{\omega} u \circ A_{\omega}^{-1}(\mu)-J_{\omega_{1}} u_{1} \circ A_{\omega_{1}}^{-1}(\mu)  \tag{3.72}\\
= & \int_{\omega=\lambda} u d l-\int_{\omega_{1}=\lambda_{1}} u_{1} d l  \tag{3.73}\\
= & \int_{s \in \mathbb{S}^{1}}\left(u(z(t, s))\left|\frac{\partial z}{\partial s}\right|(t, s)-u_{1}\left(z_{1}\left(t_{1}, s\right)\right)\left|\frac{\partial z_{1}}{\partial s}\right|\left(t_{1}, s\right)\right) d s \tag{3.74}
\end{align*}
$$

and the integrand is small (uniformly in $s$ and $\mu$ ) provided $\left\|u-u_{1}\right\|_{0}$ and $\left\|\omega-\omega_{1}\right\|_{2}$ are taken sufficiently small. This shows that $(\omega, u) \mapsto J_{\omega} u \circ A_{\omega}^{-1}$ is continuous $C^{2} \times C^{0} \rightarrow C^{0}$.

Write $\frac{d\left(J_{\omega} u \circ A_{\omega}^{-1}\right)}{d \mu}=\left(J_{\omega} \frac{\operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right) \times\left(\frac{d A_{\omega}^{-1}}{d \mu}\right)$ where $N=\frac{\nabla \omega}{|\nabla \omega|}$. Since $(\omega, u) \mapsto \frac{\operatorname{div}(u N)}{|\nabla \omega|}$ is continuous $C^{2} \times C^{1} \rightarrow C^{0}$ and $\omega \mapsto A_{\omega}^{-1}$ is continuous $C^{2} \rightarrow C^{1}$, we conclude thanks to the previous paragraph that $(\omega, u) \mapsto J_{\omega} u \circ A_{\omega}^{-1}$ is continuous $C^{2} \times C^{1} \rightarrow C^{1}$.

As for higher order derivatives, set $v_{0}=u, v_{m}=\frac{\operatorname{div}\left(v_{m-1} N\right)}{|\nabla \omega|}, m \geq 1$. A simple induction shows that $(\omega, u) \mapsto v_{m}$ is continuous $\mathcal{F}_{+}^{m+1} \times C_{\bar{\Omega}}^{m} \rightarrow C_{\bar{\Omega}}^{0}$. Now from (3.47) we find

$$
\begin{align*}
& \frac{d^{n}\left(J_{\omega} u \circ A_{\omega}^{-1}\right)}{d \mu^{n}}  \tag{3.75}\\
= & \sum_{k=1}^{n} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\
j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{i}}\left(\frac{d^{k} J_{\omega} u}{d c^{k}} \circ A_{\omega}^{-1}\right)\left(\frac{d^{j_{1}} A_{\omega}^{-1}}{d \mu^{j_{1}}}\right) \ldots\left(\frac{d^{j_{k}} A_{\omega}^{-1}}{d \mu^{j_{k}}}\right)  \tag{3.76}\\
= & \sum_{k=1}^{n} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\
j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{i}}\left(J_{\omega} u_{k} \circ A_{\omega}^{-1}\right)\left(\frac{d^{j_{1}} A_{\omega}^{-1}}{d \mu^{j_{1}}}\right) \ldots\left(\frac{d^{j_{k}} A_{\omega}^{-1}}{d \mu^{j_{k}}}\right) . \tag{3.77}
\end{align*}
$$

Again by induction $(\omega, u) \mapsto J_{\omega} u \circ A_{\omega}^{-1}$ is continuous $C^{n} \times C^{n-1} \rightarrow C^{n-1}, n \geq 2$.

Lemma $13\left(Q(\omega)=A_{\omega}^{-1}\right.$ is tame)

1. For $n \geq 0$ and $\omega$ in $\|\cdot\|_{2, \alpha}$-neighborhood,

$$
\begin{equation*}
\left\|A_{\omega}^{-1}\right\|_{n, \alpha} \leq C \cdot\left(\|\omega\|_{n, \alpha}+1\right) \tag{3.78}
\end{equation*}
$$

2. For $n \geq 0$, $\omega$ in $a\|\cdot\|_{2, \alpha}$-neighborhood and any $u$ (without restriction),

$$
\begin{equation*}
\left\|J_{\omega} u \circ A_{\omega}^{-1}\right\|_{n, \alpha} \leq C \cdot\left(\|u\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|u\|_{1, \alpha}\right) . \tag{3.79}
\end{equation*}
$$

## Proof

1 Estimate on $\left\|J_{\omega} u \circ A_{\omega}^{-1}\right\|_{0, \alpha} \quad$ Recall that $\epsilon_{0}$ is defined in (3.64). Clearly,

$$
\begin{equation*}
\left\|J_{\omega} u \circ A_{\omega}^{-1}\right\|_{0} \leq L_{1}\|u\|_{0} \tag{3.80}
\end{equation*}
$$

To estimate the Hölder-constant of $J_{\omega} u \circ A_{\omega}^{-1}$, fix $\mu, \mu^{\prime} \in[0,|\Omega|]$ and set $\lambda=A_{\omega}^{-1}(\mu)$, $\lambda^{\prime}=A_{\omega}^{-1}\left(\mu^{\prime}\right), t=\frac{\lambda-\min \omega}{\max \omega-\min \omega}, t^{\prime}=\frac{\lambda^{\prime}-\min \omega}{\max \omega-\min \omega}$. Then,

$$
\begin{align*}
& J_{\omega} u\left(\lambda^{\prime}\right)-J_{\omega} u(\lambda)  \tag{3.81}\\
= & \int_{s \in \mathbb{S}^{1}}\left(u\left(z\left(t^{\prime}, s\right)\right)\left|\frac{\partial z}{\partial s}\right|\left(t^{\prime}, s\right)-u(z(t, s))\left|\frac{\partial z}{\partial s}\right|(t, s)\right) d s  \tag{3.82}\\
= & \int_{s \in \mathbb{S}^{1}} u\left(z\left(t^{\prime}, s\right)\right)\left(\left|\frac{\partial z}{\partial s}\right|\left(t^{\prime}, s\right)-\left|\frac{\partial z}{\partial s}\right|(t, s)\right)  \tag{3.83}\\
& +\left(u\left(z\left(t^{\prime}, s\right)\right)-u(z(t, s))\right)\left|\frac{\partial z}{\partial s}\right|(t, s) d s . \tag{3.84}
\end{align*}
$$

Since $\omega$ is restricted to a $\|\cdot\|_{2, \alpha}$-neighborhood, Lemma 7 gives that the first term is bounded by $C\|u\|_{0}\left|t^{\prime}-t\right|$ while the second is bounded by $C[u]_{\alpha}\left|z\left(t^{\prime}, s\right)-z(t, s)\right|^{\alpha} \leq C[u]_{\alpha}\left|t^{\prime}-t\right|^{\alpha}$. Now from the proof of Lemma 12, $\left(A_{\omega}^{-1}\right)^{\prime}$ is bounded, so that $\left|t^{\prime}-t\right|=C\left|\lambda^{\prime}-\lambda\right| \leq C\left|\mu^{\prime}-\mu\right|$. In conclusion, for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and any $u$,

$$
\begin{equation*}
\left\|J_{\omega} u \circ A_{\omega}^{-1}\right\|_{0, \alpha} \leq C\|u\|_{0, \alpha} \tag{3.85}
\end{equation*}
$$

for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and any $u$ (without restriction). This implies (3.79) for $m=0$.

52 Estimate on $\left\|A_{\omega}^{-1}\right\|_{1, \alpha} \quad$ With $\left(A_{\omega}^{-1}\right)^{\prime}=\frac{1}{J_{\omega} \frac{1}{\mid \nabla \omega}{ }^{\circ} A_{\omega}^{-1}}$, with $\|\omega-\bar{\omega}\|_{1}$ sufficiently small $\frac{1}{\mid \nabla \omega}$ remains bounded, and thus making $\epsilon_{0}$ smaller if necessary, $\left(A_{\omega}^{-1}\right)^{\prime}$ remains bounded in $\|\cdot\|_{0, \alpha}$ for $\|\omega-\bar{\omega}\|_{2, \alpha}<\epsilon_{0}$.

3 Estimate on $\left\|u_{m}\right\|_{0, \alpha} \quad$ Suppose first (for simplicity) that $u, a$ are smooth functions of one variable and set $u_{0}:=u, u_{m}:=\left(u_{m-1} a\right)^{\prime} a, m \geq 1$. By induction, one verifies that $u_{m}$ is then of the form

$$
\begin{equation*}
u_{m}=\sum_{\substack{j_{0}+\ldots+j_{m}=m \\ j_{0}, \ldots, j_{m} \geq 1}} c_{m ; j_{1}, \ldots, j_{m}} u^{\left(j_{0}\right)} a^{\left(j_{1}\right)} \cdots a^{\left(j_{m}\right)} a^{m} . \tag{3.86}
\end{equation*}
$$

Using (5.114), and interpolation inequalities (1.39) on all factors (between their $\|\cdot\|_{0, \alpha}$ and $\|\cdot\|_{m-j_{0}, \alpha^{-}}$-norms),

$$
\begin{align*}
\left\|u_{m}\right\|_{0, \alpha} & \leq C \sum_{\substack{j_{0}+\ldots+j_{m}=m \\
j_{0}, \ldots, j_{m} \geq 1}}\|u\|_{j_{0}, \alpha}\|a\|_{j_{1}, \alpha} \cdots\|a\|_{j_{m}, \alpha}\left\|a^{m}\right\|_{0, \alpha}  \tag{3.87}\\
& \leq C\|a\|_{0, \alpha}^{l_{m}} \sum_{\substack{j_{0}+\ldots+j_{m}=m \\
j_{0}, \ldots, j_{m} \geq 0}}\|u\|_{j_{0}, \alpha}\|a\|_{m-j_{0}, \alpha}^{\frac{j_{1}}{m-j_{0}}} \cdots\|a\|_{m-j_{0}, \alpha}^{\frac{j_{m}}{m-j_{0}}}  \tag{3.88}\\
& \leq C\|a\|_{0, \alpha}^{l_{m}}\|u\|_{j_{0}, \alpha}\|a\|_{m-j_{0}, \alpha}  \tag{3.89}\\
& \left.\leq C\|a\|_{0, \alpha}^{l_{m}}\|u\|_{0, \alpha}\|a\|_{m, \alpha}+\|u\|_{m, \alpha}\|a\|_{0, \alpha}\right) \tag{3.90}
\end{align*}
$$

where $l_{m}$ is some positive integer depending on $m$. In turn,

$$
\begin{equation*}
\left\|u_{m}\right\|_{0, \alpha} \leq C\left(\|u\|_{m, \alpha}+\|a\|_{m, \alpha}\|u\|_{0, \alpha}\right), \quad m \geq 1 \tag{3.91}
\end{equation*}
$$

for all $u$ without restriction, and $a$ in a $\|\cdot\|_{0, \alpha}$-neighborhood.
The situation with $u_{m}$ as defined by $u_{m}=\frac{\operatorname{div}\left(u_{m-1} N\right)}{|\nabla \omega|}$ can be dealt with in a similar fashion, only the details are more tedious. Here, $a$ plays the rôle of $\frac{1}{|\nabla \omega|}$ or $N=\frac{\nabla \omega}{|\nabla \omega|}$. In conclusion,

$$
\begin{equation*}
\left\|u_{m}\right\|_{0, \alpha} \leq C\left(\|u\|_{m, \alpha}+\|\omega\|_{m+1, \alpha}\|u\|_{0, \alpha}\right), \quad m \geq 1 \tag{3.92}
\end{equation*}
$$

for all $\omega$ in a $\|\cdot\|_{1, \alpha}$-neighborhood and all $u$ without restriction.

44 Estimates on $\left\|J_{\omega} u_{m} \circ A_{\omega}^{-1}\right\|_{0, \alpha} \quad$ We easily conclude from (3.85) and (3.92) that

$$
\begin{equation*}
\left\|\frac{d^{m} J_{\omega} u}{d \lambda^{m}} \circ A_{\omega}^{-1}\right\|_{0, \alpha}=\left\|J_{\omega} u_{m} \circ A_{\omega}^{-1}\right\|_{0, \alpha} \leq C \cdot\left(\|u\|_{m, \alpha}+\|\omega\|_{m+1, \alpha}\|u\|_{0, \alpha}\right), \quad m \geq 0 \tag{3.93}
\end{equation*}
$$

for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood.

5 Estimate (3.78) on $\left\|A_{\omega}^{-1}\right\|_{n, \alpha}$ In this paragraph only, we write

$$
\begin{equation*}
f=A_{\omega}^{-1}, \quad g=A_{\omega}, \quad \text { and } \quad J=J_{\omega} \frac{1}{|\nabla \omega|} \tag{3.94}
\end{equation*}
$$

with derivatives $f^{\prime}, g^{\prime}, J^{\prime}$. Since $g^{\prime}=J$,

$$
\begin{equation*}
f^{(n)}=-\left(f^{\prime}\right)^{n} \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\ldots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{1}, \ldots, j_{k}} f^{(k)}\left(J^{\left(j_{1}-1\right)} \circ f\right) \cdots\left(J^{\left(j_{k}-1\right)} \circ f\right) \tag{3.95}
\end{equation*}
$$

hence by (5.114)

$$
\begin{equation*}
\left\|f^{(n)}\right\|_{0, \alpha} \leq C\left\|f^{\prime}\right\|_{0, \alpha}^{n} \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\ldots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}}\left\|f^{(k)}\right\|_{0, \alpha}\left\|J^{\left(j_{1}-1\right)} \circ f\right\|_{0, \alpha} \cdots\left\|J^{\left(j_{k}-1\right)} \circ f\right\|_{0, \alpha} . \tag{3.96}
\end{equation*}
$$

We prove by induction the estimate (3.78). We have already seen that $\|f\|_{1, \alpha}$ is bounded for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood. Suppose the estimate (3.78) proved up to some $n-1 \geq 1$. From (3.93) $\left\|J^{(j-1)} \circ f\right\|_{0, \alpha} \leq C \cdot\left(\|\omega\|_{j, \alpha}+1\right), j \geq 1$, hence by the induction hypothesis

$$
\begin{equation*}
\left\|f^{(n)}\right\|_{0, \alpha} \leq C\left\|f^{\prime}\right\|_{0, \alpha}^{n} \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}}\left(\|\omega\|_{k, \alpha}+1\right)\left(\|\omega\|_{j_{1}, \alpha}+1\right) \cdots\left(\|\omega\|_{j_{k}, \alpha}+1\right) . \tag{3.97}
\end{equation*}
$$

The double sum is the sum of $1,\|\omega\|_{k, \alpha}\|\omega\|_{j_{1}, \alpha} \cdots\|\omega\|_{j_{k}, \alpha}$, and products of fewer factors. For simplicity, consider only the term $\|\omega\|_{k, \alpha}\|\omega\|_{j_{1}, \alpha} \cdots\|\omega\|_{j_{k}, \alpha}$ (the other terms are in fact easier). We interpolate each factor between its $\|\cdot\|_{1, \alpha^{-}}$and $\|\cdot\|_{n, \alpha^{-}}$norms using (1.39)

$$
\begin{align*}
& \|\omega\|_{k, \alpha}\|\omega\|_{j_{1}, \alpha} \cdots\|\omega\|_{j_{k}, \alpha}  \tag{3.98}\\
\leq & C \cdot\|\omega\|_{1, \alpha}^{\frac{n-k}{n-1}}\|\omega\|_{n, \alpha}^{\frac{k-1}{n-1}} \cdot\|\omega\|_{1, \alpha}^{\frac{n-j_{1}}{n-1}}\|\omega\|_{n, \alpha}^{\frac{j_{1}-1}{n-1}} \cdots\|\omega\|_{1, \alpha}^{\frac{n-j_{k}}{n-1}}\|\omega\|_{n, \alpha}^{\frac{j_{k}-1}{n-1}}  \tag{3.99}\\
\leq & C \cdot\|\omega\|_{n, \alpha} \tag{3.100}
\end{align*}
$$

since $\|\omega\|_{1, \alpha}$ remains bounded and $j_{1}+\cdots+j_{k}=n$. This establishes the tame estimate (3.78) on $A_{\omega}^{-1}$ for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood.
$6 J_{\omega} u \circ A_{\omega}^{-1}$ is tame Setting $f=J_{\omega} u$ and $g=A_{\omega}^{-1}$ in (3.47) and with

$$
\begin{equation*}
u_{0}=u, \quad u_{m}=\frac{\operatorname{div}\left(u_{m-1} N\right)}{|\nabla \omega|}, \quad m \geq 1 \tag{3.101}
\end{equation*}
$$

we have

$$
\begin{equation*}
\frac{d^{n}\left(J_{\omega} u \circ A_{\omega}^{-1}\right)}{d \mu^{n}}=\sum_{k=1}^{n} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\ j_{1}, \ldots, j_{k} \geq 1}} c_{k ; j_{i}}\left(J_{\omega} u_{k} \circ A_{\omega}^{-1}\right)\left(\frac{d^{j_{1}} A_{\omega}^{-1}}{d \mu^{j_{1}}}\right) \cdots\left(\frac{d^{j_{k}} A_{\omega}^{-1}}{d \mu^{j_{k}}}\right) . \tag{3.102}
\end{equation*}
$$

From (3.93), (3.78), and (3.92), and by (5.114), we have for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and $u$ in a $\|\cdot\|_{0, \alpha}$-neighborhood,

$$
\begin{align*}
& \left\|\left(J_{\omega} u_{k} \circ A_{\omega}^{-1}\right)\left(\frac{d^{j_{1}} A_{\omega}^{-1}}{d \mu^{j_{1}}}\right) \cdots\left(\frac{d^{j_{k}} A_{\omega}^{-1}}{d \mu^{j_{k}}}\right)\right\|_{0, \alpha}  \tag{3.103}\\
\leq & C \cdot\left(\|u\|_{k, \alpha}+\|\omega\|_{k+1, \alpha}\right)\left(\|\omega\|_{j_{1}, \alpha}+1\right) \cdots\left(\|\omega\|_{j_{k}, \alpha}+1\right) . \tag{3.104}
\end{align*}
$$

Products of $\|\omega\|_{j, \alpha}$ 's can be estimated by $\|\omega\|_{n, \alpha}$ in a way similar to that which led to (3.100). Consider now $\|u\|_{k, \alpha}\|\omega\|_{j_{1}, \alpha} \cdots\|\omega\|_{j_{k}, \alpha}$ (the remaining terms are easier and their
estimation will be omitted). Interpolating each factor between its $\|\cdot\|_{1, \alpha}$ and $\|\cdot\|_{n, \alpha}$-norms with (1.39) and using the inequality $x^{\delta} y^{1-\delta} \leq x+y(0 \leq \delta \leq 1)$, we find

$$
\begin{align*}
\|u\|_{k, \alpha}\|\omega\|_{j_{1}, \alpha} \cdots\|\omega\|_{j_{k}, \alpha} & \leq C\|u\|_{1, \alpha}^{\frac{n-k}{n-1}}\|u\|_{n, \alpha}^{\frac{k-1}{n-1}}\|\omega\|_{1, \alpha}^{\frac{(k-1) n}{n-1}}\|\omega\|_{n, \alpha}^{\frac{n-k}{n-1}}  \tag{3.105}\\
& \leq C \cdot\|\omega\|_{1, \alpha}^{k-1}\left[\|u\|_{1, \alpha}\|\omega\|_{n, \alpha}+\|u\|_{n, \alpha}\|\omega\|_{1, \alpha}\right]  \tag{3.106}\\
& \leq C \cdot\left(\|u\|_{n, \alpha}+\|\omega\|_{n, \alpha}\right) \tag{3.107}
\end{align*}
$$

for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and $u$ in a $\|\cdot\|_{1, \alpha}$-neighborhood. Replacing $u$ by $\nabla \omega$ one immediately obtains

$$
\begin{equation*}
\|\omega\|_{k+1, \alpha}\|\omega\|_{j_{1}, \alpha} \cdots\|\omega\|_{j_{k}, \alpha} \leq C \cdot\|\omega\|_{n+1, \alpha}, \quad n \geq 1 \tag{3.108}
\end{equation*}
$$

for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood. Putting these together, we arrive at

$$
\begin{align*}
& \left\|\left(J_{\omega} u_{k} \circ A_{\omega}^{-1}\right)\left(\frac{d^{j_{1}} A_{\omega}^{-1}}{d \mu^{j_{1}}}\right) \cdots\left(\frac{d^{j_{k}} A_{\omega}^{-1}}{d \mu^{j_{k}}}\right)\right\|_{0, \alpha}  \tag{3.109}\\
\leq & C \cdot\left(\|\omega\|_{n+1, \alpha}+\|u\|_{n, \alpha}\right) \tag{3.110}
\end{align*}
$$

and summing,

$$
\begin{equation*}
\left\|J_{\omega} u \circ A_{\omega}^{-1}\right\|_{n, \alpha} \leq C \cdot\left(\|\omega\|_{n+1, \alpha}+\|u\|_{n, \alpha}+1\right), \quad n \geq 1 \tag{3.111}
\end{equation*}
$$

for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and all $u$ in a $\|\cdot\|_{1, \alpha}$-neighborhood. Since $J_{\omega} u \circ A_{\omega}^{-1}$ is linear in $u$, see Proposition 2.

$$
\begin{equation*}
\left\|J_{\omega} u \circ A_{\omega}^{-1}\right\|_{n, \alpha} \leq C \cdot\left(\|u\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|u\|_{1, \alpha}\right), \quad n \geq 1 \tag{3.112}
\end{equation*}
$$

for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and all $u$ without restriction. This is (3.79) for $m \geq 1$.

Even though $Q$ is in fact smooth, two derivatives are sufficient for the Moser iteration.
Lemma $14\left(Q(\omega)=A_{\omega}^{-1}\right.$ is smooth tame) $Q: \mathcal{F}_{+} \rightarrow C_{[0,|\Omega|]}^{\infty}$ is twice continuously differentiable as a map of Fréchet spaces. Its first derivative is given by

$$
\begin{equation*}
D Q(\omega) \cdot \nu=\frac{J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}} . \tag{3.113}
\end{equation*}
$$

For $n \geq 0, \omega$ in $a\|\cdot\|_{2, \alpha}$-neighborhood and any $\nu$ without restriction, we have

$$
\begin{equation*}
\left\|J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha} \quad \text { and } \quad\|D Q(\omega) \nu\|_{n, \alpha} \leq C\left\{\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right\} . \tag{3.114}
\end{equation*}
$$

The first derivative

$$
\left\{\begin{array}{ccccc}
\mathcal{F}_{+}^{m} & \times & C_{\bar{\Omega}}^{m-1} & & \begin{array}{c}
C_{\bar{\Omega}}^{m-1} \\
\omega
\end{array}  \tag{3.115}\\
\nu & & D Q(\omega) \nu
\end{array}\right\} \quad(m \geq 2)
$$

is continuous as a map of Banach spaces.
More generally, $(\omega, u) \in \mathcal{F}_{+} \times C_{0,|\Omega|]}^{\infty} \mapsto J_{\omega} u \circ A_{\omega}^{-1} \in C_{[0,|\Omega|]}^{\infty}$ is smooth tame. For $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood, $u \in C_{\bar{\Omega}}^{\infty}, \nu \in \mathcal{F}$, the first derivative in $\omega$ satisfies

$$
\begin{equation*}
\left\|D_{\omega}\left(J_{\omega} u \circ A_{\omega}^{-1}\right) \nu\right\|_{n, \alpha} \leq C\left\{\|u\|_{n+1, \alpha}\|\nu\|_{1, \alpha}+\|u\|_{2, \alpha}\|\nu\|_{n, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\|\nu\|_{1, \alpha}\right\} . \tag{3.116}
\end{equation*}
$$

For $\omega \in \mathcal{F}_{+}$in a $\|\cdot\|_{2, \alpha}$-neighborhood, $\nu_{1}, \nu_{2} \in \mathcal{F}$, and $n \geq 0$,
$\left\|D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)\right\|_{n, \alpha} \leq C\left\{\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{2, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\right\}$.

## Proof

1 Preliminary remark Let $f(\mu, \epsilon)$ be a smooth function on $[0,|\Omega|] \times\left[0, \epsilon_{0}\right]$. Then,

$$
\begin{equation*}
\frac{f(\mu, \epsilon)-f(\mu, 0)}{\epsilon} \rightarrow_{\epsilon}\left(\frac{\partial f}{\partial \epsilon}\right)(\mu, 0) \tag{3.118}
\end{equation*}
$$

uniformly in $\mu \in[0,|\Omega|]$, and thus

$$
\begin{equation*}
\left\|\frac{f(\cdot, \epsilon)-f(\cdot, 0)}{\epsilon}-\left(\frac{\partial f}{\partial \epsilon}\right)(\cdot, 0)\right\|_{0} \rightarrow_{\epsilon} 0 . \tag{3.119}
\end{equation*}
$$

Since $f$ is smooth, the same holds for all derivatives $\frac{\partial^{n} f}{\partial \mu^{n}}$, i.e.

$$
\begin{equation*}
\left\|\frac{f(\cdot, \epsilon)-f(\cdot, 0)}{\epsilon}-\left(\frac{\partial f}{\partial \epsilon}\right)(\cdot, 0)\right\|_{n} \rightarrow_{\epsilon} 0 . \tag{3.120}
\end{equation*}
$$

In other words, $\frac{f(\cdot, \epsilon)-f(\cdot, 0)}{\epsilon} \rightarrow_{\epsilon}\left(\frac{\partial f}{\partial \epsilon}\right)(\cdot, 0)$ in the $C^{\infty}$-topology.

2 Differentiability of $Q: \mathcal{F}_{+} \rightarrow C_{[0,|\Omega|]}^{\infty} \quad$ Let $\omega \in \mathcal{F}_{+}, \nu \in \mathcal{F}$, and set $\omega_{\epsilon}=\omega+\epsilon \nu$. From the previous paragraph, it is enough to show that $A_{\omega_{\epsilon}}^{-1}(\mu)$ is a smooth function of $\mu$ and $\epsilon$. But by definition, we have $A_{\omega_{\epsilon}}^{-1}\left(A_{\omega_{\epsilon}}(\lambda)\right)=\lambda$. The classical Implicit Function Theorem with parameter $\epsilon$ shows that $A_{\omega_{\epsilon}}^{-1}(\mu)$ is smooth in $\mu$ and $\epsilon$ provided $A_{\omega_{\epsilon}}(\lambda)$ is smooth in $\lambda$ and $\epsilon$. But observe that, using the change of coordinates $z^{\epsilon}$ corresponding to $\omega_{\epsilon}$, see Lemma 7 ,

$$
\begin{align*}
\frac{d}{d \lambda} A_{\omega_{\epsilon}}(\lambda) & =J_{\omega_{\epsilon}} \frac{1}{\left|\nabla \omega_{\epsilon}\right|}(\lambda)  \tag{3.121}\\
& =\int_{\omega_{\epsilon}=\lambda} \frac{1}{\left|\nabla \omega_{\epsilon}\right|} d l  \tag{3.122}\\
& =\int_{s \in \mathbb{S}^{1}} \frac{1}{\left|\nabla \omega_{\epsilon}\left(z^{\epsilon}\left(t^{\epsilon}, s\right)\right)\right|}\left|\frac{\partial z^{\epsilon}}{\partial s}\left(t^{\epsilon}, s\right)\right| d s \tag{3.123}
\end{align*}
$$

where $t^{\epsilon}=\frac{\lambda-\min \omega_{\epsilon}}{\max \omega_{\epsilon}-\min \omega_{\epsilon}}$ is obviously a smooth function of $\lambda$ and $\epsilon$. This shows that $\frac{d}{d \lambda} A_{\omega_{\epsilon}}(\lambda)$ is smooth in $\lambda$ and $\epsilon$, hence that $A_{\omega_{\epsilon}}(\lambda)$ is, as desired.

3 First derivative $D Q(\omega) \cdot \nu \quad$ Differentiating $A_{\omega_{\epsilon}}^{-1}\left(A_{\omega_{\epsilon}}(\lambda)\right)=\lambda$ at $\epsilon=0$, we find

$$
\begin{equation*}
\frac{\partial}{\partial \epsilon}_{\mid \epsilon=0} A_{\omega_{\epsilon}}^{-1}(\mu)+\frac{d A_{\omega}^{-1}}{d \mu}(\mu) \frac{\partial}{\partial \epsilon}{ }_{\mid \epsilon=0} A_{\omega_{\epsilon}}(\lambda)=0, \quad \text { where } \quad \lambda=A_{\omega}^{-1}(\mu) \tag{3.124}
\end{equation*}
$$

so that, thanks to (3.35),

$$
\begin{equation*}
\frac{\partial}{\partial \epsilon}{ }_{\mid \epsilon=0} A_{\omega_{\epsilon}}^{-1}(\mu)=\frac{d A_{\omega}^{-1}}{d \mu}(\mu)\left(J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)(\mu) \tag{3.125}
\end{equation*}
$$

i.e.

$$
\begin{equation*}
D Q(\omega) \cdot \nu=\frac{J_{\omega} \frac{\nu}{\mid \nabla \omega} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}} \tag{3.126}
\end{equation*}
$$

which is a rational function of continuous tame maps of $\omega$ and $\nu$. In particular, $D Q(\omega) \cdot \nu$ is continuous as a map of Fréchet spaces and tame in $\omega$ and $\nu$ by Lemma 12. More precisely, Lemma 12 implies that

$$
\left\{\begin{array}{ccccc}
\mathcal{F}_{+}^{m} & \times & C_{\bar{\Omega}}^{m-1} & & \begin{array}{c}
C_{[0,|\Omega|]}^{m-1} \\
\omega
\end{array}  \tag{3.127}\\
\nu & & & D Q(\omega) \nu
\end{array}\right\}
$$

is a continuous map of Banach spaces, and Lemma 13 that it is tame.

4 Tame estimates on $J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1} \quad$ For $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and any $\nu$, we have for $n \geq 0$

$$
\begin{align*}
& \left\|J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha}  \tag{3.128}\\
\leq & C\left(\left\|\frac{\nu}{|\nabla \omega|}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\frac{\nu}{|\nabla \omega|}\right\|_{1, \alpha}\right)  \tag{3.129}\\
\leq & C\left\{\|\nu\|_{n, \alpha}+\|\nu\|_{0, \alpha}\left(\|\omega\|_{n+1, \alpha}+1\right)+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right\}  \tag{3.130}\\
\leq & C\left(\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right) \tag{3.131}
\end{align*}
$$

by (3.79), (5.114), and (5.116).

5 Tame estimates on $D Q(\omega) \nu \quad$ For $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood and any $\nu$, we have for $n \geq 0$

$$
\begin{align*}
\|D Q(\omega) \nu\|_{n, \alpha} \leq & \left\|J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha}\left\|\frac{1}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}}\right\|_{0, \alpha}  \tag{3.132}\\
& +\left\|J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{0, \alpha}\left\|\frac{1}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}}\right\|_{n, \alpha}  \tag{3.133}\\
\leq & C\left\{\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}+\|\nu\|_{1, \alpha}\left(\|\omega\|_{n+1, \alpha}+1\right)\right\}  \tag{3.134}\\
\leq & C\left\{\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right\} \tag{3.135}
\end{align*}
$$

by (3.131) and (5.114).

6 First derivative of $J_{\omega} u \circ A_{\omega}^{-1}$ The operator $J_{\omega} u \circ A_{\omega}^{-1}$ is linear in $u$ so we only need to worry about differentiability in $\omega . J_{\omega_{\epsilon}} u \circ A_{\omega_{\epsilon}}^{-1}(\mu)$ is a smooth function of $(\mu, \epsilon)$, and a similar argument as for $A_{\omega}^{-1}$ shows that $\left.\frac{d}{d \epsilon}\right|_{\epsilon=0}\left(J_{\omega_{\epsilon}} u \circ A_{\omega_{\epsilon}}^{-1}\right)$ exists in the $C^{\infty}$-topology. For $\mu \in(0,|\Omega|)$, setting $\lambda=A_{\omega}^{-1}(\mu)$, and using (3.41), (3.37), and (3.50),

$$
\begin{align*}
& \frac{\partial}{\partial \epsilon}_{\left.\right|_{\epsilon=0}}\left(J_{\omega_{\epsilon}} u\left(A_{\omega_{\epsilon}}^{-1}(\mu)\right)\right)=D_{\omega}\left(J_{\omega} u \circ A_{\omega}^{-1}\right) \nu  \tag{3.136}\\
& =-J_{\omega}\left(\frac{\nu \operatorname{div}(u N)}{|\nabla \omega|}\right) \circ A_{\omega}^{-1}(\mu)  \tag{3.137}\\
& +\left(J_{\omega}\left(\frac{\operatorname{div}(u N)}{|\nabla \omega|}\right) \circ A_{\omega}^{-1}\right)(\mu)\left(\frac{J_{\omega} \frac{\nu}{|\nabla \omega|} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}}\right)(\mu) . \tag{3.138}
\end{align*}
$$

Lemma 12 implies that $(\omega, u) \mapsto J_{\omega} u \circ A_{\omega}^{-1}$ is continuously differentiable, and Lemma 13 that the derivative is tame. (From this expression, it is not too difficult to see, by an induction argument, that in fact $J_{\omega} u \circ A_{\omega}^{-1}$ is infinitely differentiable and all derivatives are
tame.)

7 Tame estimates on $D_{\omega}\left(J_{\omega} u \circ A_{\omega}^{-1}\right) \nu \quad$ We estimate $\left\|D_{\omega}\left(J_{\omega} u \circ A_{\omega}^{-1}\right) \nu\right\|_{n, \alpha}$ by the sum of three terms: setting $N=\frac{\nabla \omega}{|\nabla \omega|}$, (5.114) gives

$$
\begin{align*}
I+I I+I I I & =\left\|J_{\omega} \frac{\nu \operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha}  \tag{3.139}\\
& +\left\|J_{\omega} \frac{\operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha}\|D Q(\omega) \nu\|_{0, \alpha}  \tag{3.140}\\
& +\left\|J_{\omega} \frac{\operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{0, \alpha}\|D Q(\omega) \nu\|_{n, \alpha} . \tag{3.141}
\end{align*}
$$

Estimates on $I \quad$ By (3.131), (5.114), and (5.116), we have for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood, any $\nu, u$, and $n \geq 0$,

$$
\begin{align*}
& \left\|J_{\omega} \frac{\nu \operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha}  \tag{3.142}\\
\leq & C\left\{\|\nu \operatorname{div}(u N)\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu \operatorname{div}(u N)\|_{1, \alpha}\right\}  \tag{3.143}\\
\leq & C\left\{\|\nu\|_{n, \alpha}\|u\|_{1, \alpha}+\|u N\|_{n+1, \alpha}\|\nu\|_{0, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\|u\|_{2, \alpha}\right\}  \tag{3.144}\\
\leq & C\left\{\|\nu\|_{n, \alpha}\|u\|_{1, \alpha}+\|u\|_{n+1, \alpha}\|\nu\|_{0, \alpha}+\|N\|_{n+1, \alpha}\|u\|_{0, \alpha}\|\nu\|_{1, \alpha}\right.  \tag{3.145}\\
& \left.+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\|u\|_{2, \alpha}\right\}  \tag{3.146}\\
\leq & C\left\{\|\nu\|_{n, \alpha}\|u\|_{1, \alpha}+\|u\|_{n+1, \alpha}\|\nu\|_{0, \alpha}+\left(\|\omega\|_{n+2, \alpha}+1\right)\|u\|_{2, \alpha}\|\nu\|_{1, \alpha}\right\}  \tag{3.147}\\
\leq & C\left\{\|\nu\|_{n, \alpha}\|u\|_{1, \alpha}+\|u\|_{n+1, \alpha}\|\nu\|_{0, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\|\nu\|_{1, \alpha}\right\}  \tag{3.148}\\
\leq & C\left\{\|u\|_{1, \alpha}\|\nu\|_{n, \alpha}+\|u\|_{n+1, \alpha}\|\nu\|_{0, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\|\nu\|_{1, \alpha}\right\} \tag{3.149}
\end{align*}
$$

Estimates on $J_{\omega} \frac{\operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1} \quad$ By (3.131), (5.114), and (5.116), we have for $n \geq 0$

$$
\begin{align*}
& \left\|J_{\omega} \frac{\operatorname{div}(u N)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right\|_{n, \alpha}  \tag{3.150}\\
\leq & C\left\{\|u N\|_{n+1, \alpha}+\|\omega\|_{n+1, \alpha}\|u N\|_{2, \alpha}\right\}  \tag{3.151}\\
\leq & C\left\{\|u\|_{n+1, \alpha}+\|u\|_{0, \alpha}\|N\|_{n+1, \alpha}+\|\omega\|_{n+1, \alpha}\|u\|_{2, \alpha}\left(\|\omega\|_{3, \alpha}+1\right)\right\}  \tag{3.152}\\
\leq & C\left\{\|u\|_{n+1, \alpha}+\|u\|_{0, \alpha}\left(\|\omega\|_{n+2, \alpha}+1\right)+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\right\}  \tag{3.153}\\
\leq & C\left\{\|u\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\right\} \tag{3.154}
\end{align*}
$$

provided $\omega$ remains in a $\|\cdot\|_{2, \alpha}$-neighborhood. (We have used that, by interpolation inequalities, $\|\omega\|_{n+1, \alpha}\|\omega\|_{3, \alpha} \leq C\|\omega\|_{n+2, \alpha}\|\omega\|_{2, \alpha} \leq C\|\omega\|_{n+2, \alpha}$ for $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood.)

Estimates on II+III For $\omega$ in a $\|\cdot\|_{2, \alpha}$-neighborhood, and any $u, \nu$, we conclude from the above and (5.114) that for $n \geq 0$

$$
\begin{align*}
& I I+I I I  \tag{3.155}\\
\leq & C\left\{\left(\|u\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\right)\|\nu\|_{1, \alpha}\right.  \tag{3.156}\\
& \left.+\|u\|_{2, \alpha}\left(\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right)\right\}  \tag{3.157}\\
\leq & C\left\{\|u\|_{n+1, \alpha}\|\nu\|_{1, \alpha}+\|u\|_{2, \alpha}\|\nu\|_{n, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\|\nu\|_{1, \alpha}\right\} \tag{3.158}
\end{align*}
$$

Conclusion Putting the above together,

$$
\begin{align*}
& \left\|D_{\omega}\left(J_{\omega} u \circ A_{\omega}^{-1}\right) \nu\right\|_{n, \alpha}  \tag{3.159}\\
\leq & C\left\{\|u\|_{n+1, \alpha}\|\nu\|_{1, \alpha}+\|u\|_{2, \alpha}\|\nu\|_{n, \alpha}+\|\omega\|_{n+2, \alpha}\|u\|_{2, \alpha}\|\nu\|_{1, \alpha}\right\} . \tag{3.160}
\end{align*}
$$

8 Second derivatives of $Q$ From (3.113) it is clear that $D Q(\omega) \nu$ is continuously differentiable in $\omega$. Using (3.136) and after some simplification, the second derivative $D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)$, being the partial derivative of $D Q(\omega) \nu_{1}$ with respect to $\omega$ in the direc-
tion $\nu_{2}$, is then given by

$$
\begin{align*}
& D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)  \tag{3.161}\\
= & \frac{\left(J_{\omega} \left\lvert\, \frac{\nu_{1}}{\nabla \omega \mid} \circ A_{\omega}^{-1}\right.\right)\left(J_{\omega} \frac{\operatorname{div}\left(\frac{\nu_{2} N}{\nabla \omega \mid}\right)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)}{\left(J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)^{2}}  \tag{3.162}\\
+ & \frac{\left(J_{\omega} \frac{\nu_{2}}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)\left(J_{\omega} \frac{\operatorname{div}\left(\frac{\nu_{1} N}{|\nabla \omega|}\right)}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)}{\left(J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)^{2}}  \tag{3.163}\\
- & \frac{J_{\omega} \frac{\operatorname{div}\left(\frac{N}{\mid \nabla \omega}\right)}{|\nabla \omega|} \circ A_{\omega}^{-1}}{\left(J_{\omega} \frac{1}{\mid \omega \omega} \circ A_{\omega}^{-1}\right)^{3}}\left(J_{\omega} \frac{\nu_{1}}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)\left(J_{\omega} \frac{\nu_{2}}{|\nabla \omega|} \circ A_{\omega}^{-1}\right)  \tag{3.164}\\
- & \frac{J_{\omega} \frac{\left.\operatorname{div} \frac{\nu_{1} \nu_{2} N}{|\nabla \omega|}\right)}{|\nabla \omega|} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}} . \tag{3.165}
\end{align*}
$$

(One verifies that this expression is symmetric in $\nu_{1}$ and $\nu_{2}$.) By Lemma 12 and Lemma 13, a moment's concentration shows that this is a continuous map

$$
D^{2} Q:\left\{\begin{array}{ccccccc}
\mathcal{F}_{+}^{m} & \times & C_{\bar{\Omega}}^{m-1} & \times & C_{\bar{\Omega}}^{m-1} & & C_{[0, \mid \Omega]}^{m-2}  \tag{3.166}\\
\omega & & \nu_{1} & & \nu_{2} & & D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)
\end{array}\right\}
$$

and that it is tame.

9 Tame estimates on $D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right) \quad$ We write the above expression as $D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)=$ $I_{2}+I I_{2}+I I I_{2}+I V_{2}$. Note that all the factors not depending on $\nu_{1}$ nor $\nu_{2}$ have their $\|\cdot\|_{n, \alpha^{-}}$ norms bounded by $C\left(\|\omega\|_{n+1, \alpha}+1\right)$.

## Estimates on $I_{2}$

Using (3.85) for $\|\cdot\|_{0, \alpha}$-estimates,

$$
\begin{align*}
\left\|I_{2}\right\|_{n, \alpha} \leq & C\left\{\left(\|\omega\|_{n+1, \alpha}+1\right)\left\|\frac{\nu_{1}}{|\nabla \omega|}\right\|_{0, \alpha}\left\|\frac{\operatorname{div}}{|\nabla \omega|}\right\|_{0, \alpha}\right.  \tag{3.167}\\
& +\left(\left\|\nu_{1}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\nu_{1}\right\|_{1, \alpha}\right)\left\|\nu_{2}\right\|_{1, \alpha}  \tag{3.168}\\
& \left.+\left\|\nu_{1}\right\|_{0, \alpha}\left(\left\|\operatorname{div}\left(\frac{\nu_{2} N}{|\nabla \omega|}\right)\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\operatorname{div}\left(\frac{\nu_{2} N}{|\nabla \omega|}\right)\right\|_{1, \alpha}\right)\right\}  \tag{3.169}\\
\leq & C\left\{\left(\|\omega\|_{n+1, \alpha}+1\right)\left\|\nu_{1}\right\|_{0, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right.  \tag{3.170}\\
& +\left(\left\|\nu_{1}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\nu_{1}\right\|_{1, \alpha}\right)\left\|\nu_{2}\right\|_{1, \alpha}  \tag{3.171}\\
& \left.\left\|\nu_{1}\right\|_{0, \alpha}\left(\left\|\frac{\nu_{2} N}{|\nabla \omega|}\right\|_{n+1, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\frac{\nu_{2} N}{|\nabla \omega|}\right\|_{2, \alpha}\right)\right\}  \tag{3.172}\\
\leq & C\left\{\|\omega\|_{n+1, \alpha}\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{n, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right.  \tag{3.173}\\
& +\left\|\nu_{1}\right\|_{0, \alpha}\left(\left\|\nu_{2}\right\|_{n+1, \alpha}+\left(\|\omega\|_{n+2, \alpha}+1\right)\left\|\nu_{2}\right\|_{0, \alpha}\right.  \tag{3.174}\\
& \left.\left.+\|\omega\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\left(\|\omega\|_{3, \alpha}+1\right)\right)\right\}  \tag{3.175}\\
\leq & C\left\{\left\|\nu_{1}\right\|_{n, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{0, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\right\}(3 \tag{3.176}
\end{align*}
$$

where we have used interpolations to get $\|\omega\|_{n+1, \alpha}\|\omega\|_{3, \alpha} \leq C\|\omega\|_{n+2, \alpha}\|\omega\|_{2, \alpha} \leq C\|\omega\|_{n+2, \alpha}$.

Estimates on $I_{2}+I I_{2}$
Since $I I_{2}$ is obtained by interchanging $\nu_{1}$ and $\nu_{2}$, we immediately have

$$
\begin{equation*}
\left\|I_{2}+I I_{2}\right\|_{n, \alpha} \leq C\left\{\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{2, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\right\} . \tag{3.177}
\end{equation*}
$$

## Estimates on $\mathrm{III}_{2}$

We have

$$
\begin{align*}
& \left\|I I I_{2}\right\|_{n, \alpha}  \tag{3.178}\\
\leq & C\left\{\left(\|\omega\|_{n+1, \alpha}+1\right)\left\|\nu_{1}\right\|_{0, \alpha}\left\|\nu_{2}\right\|_{0, \alpha}\right.  \tag{3.179}\\
& +\left(\left\|\nu_{1}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\nu_{1}\right\|_{1, \alpha}\right)\left\|\nu_{2}\right\|_{0, \alpha}  \tag{3.180}\\
& \left.+\left\|\nu_{1}\right\|_{0, \alpha}\left(\left\|\nu_{2}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right)\right\}  \tag{3.181}\\
\leq & C\left\{\left\|\nu_{1}\right\|_{n, \alpha}\left\|\nu_{2}\right\|_{0, \alpha}+\left\|\nu_{1}\right\|_{0, \alpha}\left\|\nu_{2}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right\} \tag{3.182}
\end{align*}
$$

## Estimates on $I V_{2}$

We have

$$
\begin{align*}
& \|I V\|_{n, \alpha}  \tag{3.183}\\
\leq & C\left\{\left(\|\omega\|_{n+1, \alpha}+1\right)\left\|\frac{\nu_{1} \nu_{2} N}{|\nabla \omega|}\right\|_{1, \alpha}+\left\|\frac{\nu_{1} \nu_{2} N}{|\nabla \omega|}\right\|_{n+1, \alpha}\right\}  \tag{3.184}\\
\leq & C\left\{\left(\|\omega\|_{n+1, \alpha}+1\right)\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right.  \tag{3.185}\\
+ & \left.\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{0, \alpha}+\left\|\nu_{1}\right\|_{0, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\left(\|\omega\|_{n+2, \alpha}+1\right)\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right\}  \tag{3.186}\\
\leq & C\left\{\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}\right\} \tag{3.187}
\end{align*}
$$

## Conclusion

Putting the estimates on $I_{2}, I I_{2}, I I I_{2}, I V_{2}$ together,
$\left\|D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)\right\|_{n, \alpha} \leq C\left\{\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{2, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\right\}$.

## 4 Proof of Theorem 1

## Outline of proof

That $T$ is smooth tame is immediate since $T(F)=(Q \circ \Delta \circ S)(F)$ is a composition of smooth tame maps. The crucial part in the surjective part of the Nash-Moser theorem is to establish that $D T(F) f$ has a tame family of right-inverses. We emphasize that the nondegeneracy condition (ND2) is only made at the reference steady-state, and not in an entire neighborhood of the reference steady-state. The problem of finding such right-inverse for $D T(F) f$, given in (4.14), is equivalent to inverting a map of the form $g+K(F) g=h$ where $K(F) g$ can be thought of as a "compact perturbation" of the first term $g$, see (4.16). This is precisely what was done in Lemma 5 with the elliptic operator $\Delta \phi+c \phi=k$ (augmented with suitable boundary conditions): $c \phi$ is a "compact perturbation" of $\Delta \phi$. There, the estimates on the (bilinear) term $c \phi$ were standard. Here, the term $K(F) g$ is more complicated and requires considerably more work.

The injective part of Theorem 1 is proved in Section 4.4. At the conceptual level, the proof is an adjustment of the injective part of the Nash-Moser Inverse Function Theorem as presented in Section 1.3, Part III [7]. (One cannot use this theorem directly because of complications created by the lack of injectivity of the map $F \mapsto \psi$.)

## Assumptions

We recall the main assumptions.
The domain $\Omega$ is assumed diffeomorphic to the annulus so that

$$
\begin{equation*}
\partial \Omega=\Gamma_{o} \cup \Gamma_{i} . \tag{4.1}
\end{equation*}
$$

We assume that the reference steady-state $\bar{\omega}=\bar{F}(\bar{\psi})$ is such that $\bar{F}^{\prime} \neq 0$, $\bar{\omega}$ has no critical points, and satisfies the non-degeneracy conditions (ND1) and (ND2). By continuity of $F \in C^{1} \mapsto \psi \in C^{2}$, make $\epsilon_{S}$ in $\mathcal{V}_{S}(\bar{F})$ smaller if necessary (see Proposition 6), so that the corresponding $\psi$ has no critical points either and that $F^{\prime} \neq 0$. We will then assume without loss of generality that

$$
\begin{equation*}
\psi \leq 0 \tag{4.2}
\end{equation*}
$$

The interval $I$ introduced in Section 2.2 can now be taken of the form

$$
\begin{equation*}
I=[\underline{c}, 0] \quad \text { where } \quad \underline{c}<\min \bar{\psi} \quad \text { is fixed. } \tag{4.3}
\end{equation*}
$$

For simplicity, the calculations will be performed assuming that

$$
\begin{equation*}
F^{\prime}>0 \tag{4.4}
\end{equation*}
$$

As observed in the Introduction, the case $F^{\prime}>0$ is special in that the corresponding solution automatically satisfies both non-degeneracy conditions (ND1) and (ND2). But this property will never be used in the following.

## The first derivative $D T(F) f$

The map

$$
T:\left\{\begin{array}{ccc}
\mathcal{V}_{S}^{n}(\bar{F}) & C_{[0,|\Omega|]}^{n}  \tag{4.5}\\
F & & A_{\omega}^{-1}
\end{array}\right\} \quad(n \geq 2)
$$

is continuous as a map of Banach spaces (Lemma 12 and (2.103) after Proposition 6). Write the first derivative as $D T(F) f=D Q(\omega) \nu$ where $\omega=\Delta \psi$ and $\psi=S(F)$ solves the steady-state equation

$$
\begin{equation*}
\Delta \psi=F(\psi), \quad \psi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \psi}{\partial \tau}{ }_{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \psi}{\partial N}=\gamma_{i}, \tag{4.6}
\end{equation*}
$$

and $\nu=\Delta \phi$ where $\phi=D S(F) f$ solves the linearized steady-state equation

$$
\begin{equation*}
\Delta \phi=F^{\prime}(\psi) \phi+f(\psi), \quad \phi_{\mid \Gamma_{o}}=0, \quad \frac{\partial \phi}{\partial \tau}{\mid \Gamma_{i}}=0, \quad \int_{\Gamma_{i}} \frac{\partial \phi}{\partial N}=0 . \tag{4.7}
\end{equation*}
$$

Thus,

$$
D T:\left\{\begin{array}{ccccc}
\mathcal{V}_{S}^{n}(\bar{F}) & \times & C_{I}^{n-1} & & \begin{array}{c}
C_{[0,|\Omega|]}^{n-1} \\
F
\end{array}  \tag{4.8}\\
& f & & D T(F) f
\end{array}\right\} \quad(n \geq 2)
$$

is continuous as a map of Banach spaces (Lemma 14 and (2.104) after Proposition 6). From Proposition 11 ,

$$
\begin{equation*}
D T(F) f=\frac{J_{\omega} \frac{\nu}{\mid \nabla \omega} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}}=\frac{J_{\omega} \frac{f(\psi)}{|\nabla \omega|} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}}+\frac{J_{\omega} \frac{F^{\prime}(\psi) \phi}{|\nabla \omega|} \circ A_{\omega}^{-1}}{J_{\omega} \frac{1}{|\nabla \omega|} \circ A_{\omega}^{-1}} \tag{4.9}
\end{equation*}
$$

To simplify this, use the identity (when $F^{\prime}>0$ )

$$
\begin{equation*}
T(F)=F \circ A_{\psi}^{-1} \tag{4.10}
\end{equation*}
$$

which follows from $A_{\psi}(\lambda)=|\{\psi<\lambda\}|=|\{\omega<F(\lambda)\}|=A_{\omega}(F(\lambda))$. (A similar expression holds in the case $F^{\prime}<0$.) Then, $\omega(x)=A_{\omega}^{-1}(\lambda)$ if and only if $\psi(x)=A_{\psi}^{-1}(\lambda)$ and

$$
\begin{equation*}
\int_{\omega=A_{\omega}^{-1}(\lambda)} \frac{f(\psi)}{|\nabla \omega|} d l=f\left(A_{\psi}^{-1}(\lambda)\right) \int_{\omega=A_{\omega}^{-1}(\lambda)} \frac{1}{|\nabla \omega|} d l . \tag{4.11}
\end{equation*}
$$

Also, $\nabla \omega=F^{\prime}(\psi) \nabla \psi$ so that after some elementary calculations we obtain

$$
\begin{align*}
D T(F) \cdot f & =f \circ A_{\psi}^{-1}+\left(F^{\prime} \circ A_{\psi}^{-1}\right)\left(\frac{J_{\psi \frac{\phi}{|\nabla \psi|}} \circ A_{\psi}^{-1}}{J_{\psi} \frac{1}{|\nabla \psi|} \circ A_{\psi}^{-1}}\right)  \tag{4.12}\\
& =f \circ A_{\psi}^{-1}+\left(\frac{d A_{\omega}^{-1}}{d \mu}\right)\left(J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right)  \tag{4.13}\\
& =: B(F) f+\tilde{K}(F) f . \tag{4.14}
\end{align*}
$$

Construction of a tame right-inverse $f=L(F) h$ to $h=D T(F) \cdot f$
In order to construct a right-inverse $f=L(F) h$ to $h=D T(F) f$ we will construct first the inverse to a modification $h=M(F) g$ of $h=D T(F) f$. Set

$$
\begin{equation*}
M(F):=D T(F) \cdot V B(F)=\operatorname{Id}_{[[0,|\Omega|]}^{\infty}+K(F) \tag{4.15}
\end{equation*}
$$

where

$$
\begin{equation*}
K(F) \cdot g:=\tilde{K}(F) \cdot V B(F) \cdot g=\left(\frac{d A_{\omega}^{-1}}{d \mu}\right)\left(J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right) \tag{4.16}
\end{equation*}
$$

and $\Delta \phi=F^{\prime}(\psi) \phi+f(\psi), f=V B(F) \cdot g$, and $f=V B(F) g$ is a tame right-inverse to $g=B(F) f$ (intuitively, it is " $f=g \circ A_{\psi}$ "). The latter is constructed in Lemma 15, We will show that $h=M(F) g$ has a tame inverse $g=V M(F) h$, and thus that $h=D T(F) f$ has a tame right-inverse by setting

$$
\begin{equation*}
L(F) \cdot h:=V B(F) \cdot V M(F) \cdot h . \tag{4.17}
\end{equation*}
$$

This is possible because $K(F) g$ can be viewed as a "compact perturbation" of $g$. This is due to the fact that $\phi$ gains sufficient regularity from $f$. In fact, the proof is completely analogous to the proof that $\Delta \phi+c \phi=k$ has a family of tame inverses: compare Lemma 4 with Lemma 16, and Proposition 5 with Proposition 18 ,

### 4.1 Right-inverse to $B(F) f=f \circ A_{\psi}^{-1}$

First, as an auxiliary step, we need to construct a right-inverse $f=V B(F) g$ to $g=B(F) f$. Since $B(F)$ is surjective for each $F$, we know that a right-inverse exists for each $F$. However, we need the inverse $f=V B(F) g$ to $g=B(F) f$ to be continuous in both $F$ and $h$, and to satisfy tame estimates. Naively, the inverse of $g=f \circ A_{\psi}^{-1}$ should be " $f=g \circ A_{\psi}$ ", but $f$ is defined on an interval larger than the domain of $A_{\psi}$.

Lemma 15 The map $B(F) \cdot f$ has a smooth tame family of right-inverses $f=V B(F) \cdot g$ defined on a sufficiently small $\|\cdot\|_{1, \alpha}$-neighborhood $\mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F})$ of $\bar{F}$ :

$$
\begin{equation*}
V B(F): C_{[0,|\Omega|]}^{\infty} \rightarrow C_{[\underline{c}, 0]}^{\infty}, \quad B(F) \cdot V B(F)=I d_{C_{[0,|\Omega|]}^{\infty}} \tag{4.18}
\end{equation*}
$$

For $n \geq 2, F \in \mathcal{V}_{B}(\bar{F})$ and any $g$ (without restriction),

$$
\begin{equation*}
\|f\|_{n, \alpha} \leq C \cdot\left\{\|g\|_{n, \alpha}+\|F\|_{n-2, \alpha}\|g\|_{1, \alpha}\right\} \tag{4.19}
\end{equation*}
$$

while $\|f\|_{0, \alpha} \leq C\|g\|_{0, \alpha},\|f\|_{1, \alpha} \leq C\|g\|_{1, \alpha}$.

Proof Fix $D>0$ and let $\bar{B} \in C_{[-D,|\Omega|+D]}^{\infty}$ be a monotone increasing extension of $A_{\bar{\psi}}^{-1} \in$ $C_{[0,|\Omega|]}^{\infty}$. It can be arranged so that range $(\bar{B})=[\underline{c},|\underline{c}|]$ (see (4.3)). See proof of Corollary 1.3.7, p. 138, Part II of [7]. Let

$$
\begin{equation*}
\mathcal{E}_{0}: C_{[0,|\Omega|]}^{\infty} \quad \longrightarrow \quad\left\{b \in C_{[-D,|\Omega|+D]}^{\infty} \mid b(-D)=b(|\Omega|+D)=0\right\} \tag{4.20}
\end{equation*}
$$

be an extension operator taking functions on $[0,|\Omega|]$ to functions on $[-D,|\Omega|+D]$ vanishing at the endpoints. (The target space is easily seen to be a tame Fréchet space). It can be made tame linear of degree 0: again from the proof of Corollary II.1.3.7, p. 138, [Hamilton]), extend $b=b(\lambda) \in C_{[0,|\Omega|]}^{\infty}$ to $\lambda \leq 0$, then to $\lambda \geq|\Omega|$, and finally multiply by a smooth cut-off function with support in $(-D,|\Omega|+D)$ and equal to 1 on $[0,|\Omega|]$. Then we have the tame estimates for $n \geq 0$

$$
\begin{equation*}
\left\|\mathcal{E}_{0} b\right\|_{n, \alpha} \leq C\|b\|_{n, \alpha} \tag{4.21}
\end{equation*}
$$

Define now

$$
\mathcal{E}:\left\{\begin{array}{ccc}
C_{[0,|\Omega|]}^{\infty} & \longrightarrow & \begin{array}{c}
C_{[-D,|\Omega|+D]}^{\infty} \\
B
\end{array}  \tag{4.22}\\
B+\mathcal{E}_{0}\left(B-A_{\bar{\psi}}^{-1}\right)
\end{array}\right\}
$$

which extends maps defined on $[0,|\Omega|]$ to maps defined on $[-D,|\Omega|]$ with fixed endpoint values $\underline{c}$ and $|\underline{c}|$ at $-D$ and $|\Omega|+D$ respectively. $\mathcal{E}$ is smooth tame since it is affine with tame linear part $\mathcal{E}_{0}$. For a sufficiently small $\|\cdot\|_{1}$-neighborhood $\mathcal{V}\left(A_{\bar{\psi}}^{-1}\right)$ of $A_{\bar{\psi}}^{-1}$, it also defines a map

$$
\begin{equation*}
\mathcal{E}:\left(\mathcal{V}\left(A_{\bar{\psi}}^{-1}\right) \subset C_{[0,|\Omega|]}^{\infty}\right) \quad \longrightarrow \quad \mathcal{D}_{I_{1}, I_{2}}^{\infty} \tag{4.23}
\end{equation*}
$$

where $\mathcal{D}_{I_{1}, I}^{\infty}$ is the set of smooth diffeomorphisms from $I_{1}=[-D,|\Omega|]$ to $I_{2}=[\underline{c},|\underline{c}|]$.
From $g=f \circ A_{\psi}^{-1}$ we find $\mathcal{E}(g)=\mathcal{E}\left(f \circ A_{\psi}^{-1}\right)=f \circ \mathcal{E}\left(A_{\psi}^{-1}\right)$. Let then

$$
V:\left\{\begin{array}{ccc}
\mathcal{D}_{I_{1}, I_{2}}^{\infty} & \longrightarrow & \mathcal{D}_{I_{2}, I_{1}}^{\infty}  \tag{4.24}\\
B & & B^{-1}
\end{array}\right\}
$$

denote the operator which takes inverses. Choose now

$$
\begin{equation*}
\mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F}) \tag{4.25}
\end{equation*}
$$

a sufficiently small $\|\cdot\|_{1, \alpha}$-neighborhood of $\mathcal{V}_{S}(\bar{F})$ from Proposition 6 so that the corresponding $A_{\psi}^{-1}$ remains in $\mathcal{V}\left(A_{\bar{\psi}}^{-1}\right)$ (use also Lemma 12). We have constructed a right-inverse

$$
\begin{equation*}
V B(F) \cdot g:=\left.\left(\mathcal{E} g \circ V\left(\mathcal{E}\left(A_{\psi}^{-1}\right)\right)\right)\right|_{[c, 0]} \tag{4.26}
\end{equation*}
$$

defined for any $F \in \mathcal{V}_{B}(\bar{F})$ and any $g \in C_{[0,|\Omega|]}^{\infty}$.
Clearly, $f=V B(F) g$ is smooth tame. Using the above estimates on $\mathcal{E}_{0}$, on $A_{\psi}^{-1}$ from Proposition 11, and on $\psi=S(F)$ from Proposition 6, we have for $n \geq 2$

$$
\begin{equation*}
\left\|\mathcal{E}\left(A_{\psi}^{-1}\right)\right\|_{n, \alpha} \leq C\left(\left\|A_{\psi}^{-1}\right\|_{n, \alpha}+1\right) \leq C\left(\|\psi\|_{n, \alpha}+1\right) \leq C\left(\|F\|_{n-2, \alpha}+1\right) . \tag{4.27}
\end{equation*}
$$

From tame estimates on composition of functions from Lemma 26 and on the inversion operator $V$ from Lemma 27$]$ of the Appendix, for $F \in \mathcal{V}_{B}(\bar{F})$ and any $g \in C_{[0,|\Omega|}^{\infty}$,

$$
\begin{align*}
\|f\|_{n, \alpha} & \leq C\left(\|\mathcal{E}(g)\|_{n, \alpha}+\left\|\mathcal{E}\left(A_{\psi}^{-1}\right)\right\|_{n, \alpha}\|\mathcal{E}(g)\|_{1, \alpha}\right)  \tag{4.28}\\
& \leq C\left(\|g\|_{n, \alpha}+\left(\|F\|_{n-2, \alpha}+1\right)\| \| g \|_{1, \alpha}\right)  \tag{4.29}\\
& \leq C\left(\|g\|_{n, \alpha}+\|F\|_{n-2, \alpha}\| \| g \|_{1, \alpha}\right) . \tag{4.30}
\end{align*}
$$

Note that $\left\|\mathcal{E}\left(A_{\psi}^{-1}\right)\right\|_{1, \alpha}$ remains bounded for $F \in \mathcal{V}_{B}(\bar{F})$. Thus, we deduce easily the desired estimates on $\|f\|_{0, \alpha}$ and $\|f\|_{1, \alpha}$.

### 4.2 Summary of tame estimates

Here we collect tame estimates which will be used abundantly in the next Sections. Some estimates will be given in two equivalent forms, the second being particularly useful for the estimates on the difference $K(F) g-K(\bar{F}) g$.

From Proposition 6; If $\omega=\Delta \psi=F(\psi)$ with $F \in \mathcal{V}_{S}(\bar{F})$, then for $n \geq 0$

$$
\begin{equation*}
\|\omega\|_{n, \alpha} \quad \text { and } \quad\|\psi\|_{n+2, \alpha} \leq C\left\{\|F\|_{n, \alpha}+1\right\} \leq C\left\{\|F-\bar{F}\|_{n, \alpha}+1\right\} \tag{4.31}
\end{equation*}
$$

from the triangle inequality $\|F\|_{n, \alpha} \leq\|\bar{F}\|_{n, \alpha}+\|F-\bar{F}\|_{n, \alpha} \leq C\left(\|F-\bar{F}\|_{n, \alpha}+1\right)$. If $\Delta \phi=F^{\prime}(\psi) \phi+f(\psi)$ with $F \in \mathcal{V}_{S}(\bar{F})$ and $f \in C_{\bar{\Omega}}^{\infty}$, then for $n \geq 0$

$$
\begin{align*}
\|\phi\|_{n+2, \alpha} & \leq C\left\{\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}\right\}  \tag{4.32}\\
\|\phi\|_{n+2, \alpha} & \leq C\left\{\|f\|_{n, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|f\|_{1, \alpha}\right\} \tag{4.33}
\end{align*}
$$

again by the triangle inequality and using $\|f\|_{1, \alpha} \leq\|f\|_{n, \alpha}$ for $n \geq 1$ (for $n=0$, the last term in (4.32) and (4.33) is actually not needed, see Proposition (6).
From Lemma 15 If $f=V B(F) g$ where $F \in \mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F})$ and $g \in C_{[0,|\Omega|]}^{\infty}$, then for $n \geq 2$,

$$
\begin{align*}
\|f\|_{n, \alpha} & \leq C\left\{\|g\|_{n, \alpha}+\|F\|_{n-2, \alpha}\|g\|_{1, \alpha}\right\}  \tag{4.34}\\
\|f\|_{n, \alpha} & \leq C\left\{\|g\|_{n, \alpha}+\|F-\bar{F}\|_{n-2, \alpha}\|g\|_{1, \alpha}\right\} \tag{4.35}
\end{align*}
$$

while $\|f\|_{0, \alpha} \leq C\|g\|_{0, \alpha},\|f\|_{1, \alpha} \leq C\|g\|_{1, \alpha}$.
Combining the above, we have for $n \geq 0$,

$$
\begin{align*}
\|\phi\|_{n+2, \alpha} & \leq C\left\{\|g\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\}  \tag{4.36}\\
\|\phi\|_{n+2, \alpha} & \leq C\left\{\|g\|_{n, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\} . \tag{4.37}
\end{align*}
$$

From Lemma 14; For $\psi \in \mathcal{F}_{+}$in a $\|\cdot\|_{2, \alpha}$-neighborhod, $\phi \in C_{\bar{\Omega}}^{\infty}$, and $n \geq 0$,

$$
\begin{align*}
& \left\|J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right\|_{n, \alpha} \text { and }\|D Q(\psi) \phi\|_{n, \alpha} \leq C\left(\|\phi\|_{n, \alpha}+\|\psi\|_{n+1, \alpha}\|\phi\|_{1, \alpha}\right),  \tag{4.38}\\
& \left\|J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right\|_{n, \alpha} \text { and }\|D Q(\psi) \phi\|_{n, \alpha} \leq C\left(\|\phi\|_{n, \alpha}+\|\psi-\bar{\psi}\|_{n+1, \alpha}\|\phi\|_{1, \alpha}\right) . \tag{4.39}
\end{align*}
$$

In turn, for $n \geq 0$ and $\psi \in \mathcal{F}_{+}$in a $\|\cdot\|_{2, \alpha}$-neighborhood (see proof below),

$$
\begin{equation*}
\left\|A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right\|_{n, \alpha} \leq C\|\psi-\bar{\psi}\|_{n+1, \alpha} . \tag{4.40}
\end{equation*}
$$

For $n \geq 0, F$ in the $\|\cdot\|_{2, \alpha}$-neighborhood $\mathcal{V}_{S}(\bar{F})$ (see proof below),

$$
\begin{equation*}
\|\omega-\bar{\omega}\|_{n, \alpha} \leq C\|\psi-\bar{\psi}\|_{n+2, \alpha} \leq C\|F-\bar{F}\|_{n+1, \alpha} . \tag{4.41}
\end{equation*}
$$

Combining these two estimates we obtain

$$
\begin{equation*}
\left\|A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right\|_{n, \alpha} \leq C\|F-\bar{F}\|_{n, \alpha} \tag{4.42}
\end{equation*}
$$

One would expect from the above that $\left\|A_{\omega}^{-1}-A_{\bar{\omega}}^{-1}\right\|_{n, \alpha} \leq C\|F-\bar{F}\|_{n+2, \alpha}$. However, the regularizing effect of $\omega=\Delta \psi=F(\psi)$ gives the better estimates (see proof below):

$$
\begin{equation*}
\left\|A_{\omega}^{-1}-A_{\bar{\omega}}^{-1}\right\|_{n, \alpha} \leq C\|F-\bar{F}\|_{n, \alpha} \tag{4.43}
\end{equation*}
$$

For $n \geq 0, h, \bar{h} \in C_{I}^{\infty}$, and $\psi, \bar{\psi} \in C_{\bar{\Omega}}^{\infty}$ in a $\|\cdot\|_{1, \alpha}$-neighborhood (see proof below),

$$
\begin{equation*}
\|h(\psi)-\bar{h}(\bar{\psi})\|_{n, \alpha} \leq C\left\{\|h\|_{n+1, \alpha}\|\psi-\bar{\psi}\|_{0, \alpha}+\|h\|_{2, \alpha}\|\psi-\bar{\psi}\|_{n, \alpha}+\|h-\bar{h}\|_{n, \alpha}\right\} \tag{4.44}
\end{equation*}
$$

Observe then that for $F$ in the $\|\cdot\|_{2, \alpha}$-neighborhood $\mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F})$,

$$
\begin{equation*}
\|F\|_{2, \alpha}, \quad\|\omega\|_{2, \alpha}, \quad\|\psi\|_{4, \alpha}, \quad\left\|A_{\omega}^{-1}\right\|_{2, \alpha}, \quad\left\|A_{\psi}^{-1}\right\|_{4, \alpha} \quad \leq \quad C \tag{4.45}
\end{equation*}
$$

This allows to incorporate a number of terms of lower order into a constant.
We also recall the standard interpolation inequalities

$$
\begin{align*}
\|u\|_{n+r, \alpha}\|v\|_{s+m, \alpha} & \leq C\left\{\|u\|_{n+r+m, \alpha}\|v\|_{s, \alpha}+\|u\|_{s, \alpha}\|v\|_{n+r+m, \alpha}\right\},  \tag{4.46}\\
\|u\|_{n+r, \alpha}\|u\|_{s+m, \alpha} & \leq C\|u\|_{n+r+m, \alpha}\|u\|_{s, \alpha}
\end{align*}
$$

which are consequences of (1.39).

Proof of (4.40) Set $\omega_{t}=\bar{\omega}+t(\omega-\bar{\omega})$ and note that $Q(\omega)-Q(\bar{\omega})=\int_{0}^{1} D Q\left(\omega_{t}\right)(\omega-\bar{\omega}) d t$. Then, $\omega_{t}$ remains in a $\|\cdot\|_{2, \alpha^{-}}$-neighborhood and by (4.39) we have

$$
\begin{align*}
\left\|A_{\omega}^{-1}-A_{\omega_{0}}^{-1}\right\|_{m, \alpha} & \leq \int_{0}^{1}\left\|D Q\left(\omega_{t}\right)(\omega-\bar{\omega})\right\|_{m, \alpha} d t  \tag{4.47}\\
& \leq C\left\{\|\omega-\bar{\omega}\|_{m, \alpha}+\|\omega-\bar{\omega}\|_{m+1, \alpha}\|\omega-\bar{\omega}\|_{1, \alpha}\right\}  \tag{4.48}\\
& \leq C\|\omega-\bar{\omega}\|_{m+1, \alpha} \tag{4.49}
\end{align*}
$$

since $\left\|\omega_{t}-\bar{\omega}\right\|_{j, \alpha} \leq\|\omega-\bar{\omega}\|_{j, \alpha}$ for $0 \leq t \leq 1$.

Proof of (4.41) Write $\omega-\bar{\omega}=\Delta \psi-\Delta \bar{\psi}$, and set $F_{t}=\bar{F}+t(F-\bar{F})$ and $\psi_{t}=S\left(F_{t}\right)$. Then, $\psi-\bar{\psi}=\int_{0}^{1} \phi_{t} d t$ where $\Delta \phi_{t}=F_{t}^{\prime}\left(\psi_{t}\right) \phi_{t}+(F-\bar{F})\left(\psi_{t}\right)$. Then by (4.36),

$$
\begin{align*}
\left\|\phi_{t}\right\|_{n+2, \alpha} & \leq C\left(\|F-\bar{F}\|_{n, \alpha}+\left\|F_{t}\right\|_{n+1, \alpha}\|F-\bar{F}\|_{1, \alpha}\right)  \tag{4.50}\\
& \leq C\left(\|F-\bar{F}\|_{n, \alpha}+\left(\|F-\bar{F}\|_{n+1, \alpha}+1\right)\|F-\bar{F}\|_{1, \alpha}\right)  \tag{4.51}\\
& \leq C\|F-\bar{F}\|_{n+1, \alpha} \tag{4.52}
\end{align*}
$$

Integrating, this gives the desired estimate.
Proof of (4.43) Write

$$
\begin{align*}
A_{\omega}^{-1}-A_{\bar{\omega}}^{-1} & =(F-\bar{F})\left(A_{\psi}^{-1}\right)+\bar{F}\left(A_{\psi}^{-1}\right)-\bar{F}\left(A_{\bar{\psi}}^{-1}\right)  \tag{4.53}\\
& =(F-\bar{F})\left(A_{\psi}^{-1}\right)+\int_{0}^{1} \bar{F}^{\prime}\left(A_{\bar{\psi}}^{-1}+t\left(A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right)\right)\left(A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right) d t . \tag{4.54}
\end{align*}
$$

The first term is estimated by

$$
\begin{align*}
\left\|(F-\bar{F})\left(A_{\psi}^{-1}\right)\right\|_{n, \alpha} & \leq\|F-\bar{F}\|_{n, \alpha}+\left\|A_{\psi}^{-1}\right\|_{n, \alpha}\|F-\bar{F}\|_{1, \alpha}  \tag{4.55}\\
& \leq\|F-\bar{F}\|_{n, \alpha}+\left\|A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right\|_{n, \alpha}\|F-\bar{F}\|_{1, \alpha}  \tag{4.56}\\
& \leq C\|F-\bar{F}\|_{n, \alpha} \tag{4.57}
\end{align*}
$$

while we have

$$
\begin{align*}
\left\|\bar{F}^{\prime}\left(A_{\bar{\psi}}^{-1}+t\left(A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right)\right)\right\|_{n, \alpha} & \leq C\left(1+\left\|A_{\bar{\psi}}^{-1}+t\left(A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right)\right\|_{n, \alpha}\right)  \tag{4.58}\\
& \leq C\left(1+\left\|A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right\|_{n, \alpha}\right)  \tag{4.59}\\
& \leq C\left(1+\|F-\bar{F}\|_{n, \alpha}\right) \tag{4.60}
\end{align*}
$$

so that the second term is also bounded by $\|F-\bar{F}\|_{n, \alpha}$.

Proof of (4.44) Write $\psi_{t}=\bar{\psi}+t(\psi-\bar{\psi})$. Then

$$
\begin{equation*}
h(\psi)-\bar{h}(\bar{\psi})=h(\psi)-h(\bar{\psi})+(h-\bar{h})(\bar{\psi})=\int_{0}^{1} h^{\prime}\left(\psi_{t}\right)(\psi-\bar{\psi}) d t+(h-\bar{h})(\bar{\psi}) \tag{4.61}
\end{equation*}
$$

so that by (5.114) and (5.121),

$$
\begin{align*}
& \|h(\psi)-\bar{h}(\bar{\psi})\|_{n, \alpha}  \tag{4.62}\\
\leq & C\left\{\left(\|h\|_{n+1, \alpha}+\|\psi-\bar{\psi}\|_{n, \alpha}\|h\|_{2, \alpha}\right)\|\psi-\bar{\psi}\|_{0, \alpha}\right.  \tag{4.63}\\
& \left.+\|h\|_{2, \alpha}\|\psi-\bar{\psi}\|_{n, \alpha}+\|h-\bar{h}\|_{n, \alpha}\right\}  \tag{4.64}\\
\leq & C\left\{\|h\|_{n+1, \alpha}\|\psi-\bar{\psi}\|_{0, \alpha}+\|h\|_{2, \alpha}\|\psi-\bar{\psi}\|_{n, \alpha}+\|h-\bar{h}\|_{n, \alpha}\right\} . \tag{4.65}
\end{align*}
$$

### 4.3 Id $+K(F) g$ has a tame family of inverses

Recall that $K(F) g$ is defined in (4.16). The procedure to construct a tame family of inverses to Id $+K(F) g$ completely parallels the proof that $\Delta \phi+c \phi=k$ (augmented with suitable boundary conditions) has a tame family of inverses, see Lemma 4 and Proposition 5. In this case though, the estimates on $K(F) g-K(\bar{F}) g$ require significantly more work and they are derived in the separate Lemma 17 ,

Lemma 16 (Estimates for $h=g+K(F) g)$ Let $h=g+K(F) g$ where $F \in \mathcal{V}_{B}(\bar{F}) \subset$ $\mathcal{V}_{S}(\bar{F})$ and $g \in C_{[0,|\Omega|]}^{\infty}$. Then, for any $n \geq 0$,

$$
\begin{equation*}
\|g\|_{n, \alpha} \leq C\left\{\|h\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\} \tag{4.66}
\end{equation*}
$$

where $C$ may depend on $n$ and $\mathcal{V}_{B}(\bar{F})$.
Proof The estimates on $g$ are derived by writing $g=h-K(F) g$. We first estimate $K(F) g$ : by (5.114), (4.38), (4.31), and (4.36),

$$
\begin{align*}
\|K(F) g\|_{n, \alpha} & \leq C\left\|A_{\omega}^{-1}\right\|_{n+1, \alpha}\left\|J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right\|_{0, \alpha}+C\left\|A_{\omega}^{-1}\right\|_{1, \alpha}\left\|J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right\|_{n, \alpha}(4 .  \tag{array}\\
& \leq C\left(\|\omega\|_{n+1, \alpha}+1\right)\|\phi\|_{1, \alpha}+C\left(\|\phi\|_{n, \alpha}+\|\psi\|_{n+1, \alpha}\|\phi\|_{1, \alpha}\right)  \tag{4.68}\\
& \leq C\left\{\|\phi\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|\phi\|_{1, \alpha}\right\}  \tag{4.69}\\
& \leq C\left\{\|g\|_{n-2, \alpha}+\|F\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\} . \tag{4.70}
\end{align*}
$$

for $F \in \mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F})$ and any $g$. In turn,

$$
\begin{align*}
\|g\|_{n, \alpha} & \leq\|h\|_{n, \alpha}+\|K(F) g\|_{n, \alpha}  \tag{4.71}\\
& \leq\|h\|_{n, \alpha}+C\left\{\|g\|_{n-2, \alpha}+\|F\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\} \tag{4.72}
\end{align*}
$$

and the term $\|g\|_{n-2, \alpha}$ can be incorporated into the left-hand side as usual and we find

$$
\begin{equation*}
\|g\|_{n, \alpha} \leq C\left\{\|h\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\} . \tag{4.73}
\end{equation*}
$$

We will need estimates on the difference $K(F) g-K(\bar{F}) g$ for $F \in \mathcal{V}_{B}(\bar{F})$. Recall that $K(F) g$ is defined in (4.16).

Lemma 17 (Estimates on $K(F) g-K(\bar{F}) g$ ) For $n \geq 0, F \in \mathcal{V}_{S}(\bar{F})$, and $g \in C_{[0,|\Omega|]}^{\infty}$,

$$
\begin{equation*}
\|K(F) g-K(\bar{F}) g\|_{n, \alpha} \leq C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{2, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|g\|_{n-1, \alpha}\right\} . \tag{4.74}
\end{equation*}
$$

Proof $\quad F$ and $g$ being given, let $f=V B(F) g, \bar{f}=V B(\bar{F}) g$ and $\phi, \bar{\phi}$ such that

$$
\begin{equation*}
\Delta \phi=F^{\prime}(\psi) \phi+f(\psi), \quad \Delta \bar{\phi}=\bar{F}^{\prime}(\bar{\psi}) \bar{\phi}+\bar{f}(\bar{\psi}) . \tag{4.75}
\end{equation*}
$$

Write then

$$
\begin{align*}
K(F) g-K(\bar{F}) g= & \left(A_{\omega}^{-1}-A_{\bar{\omega}}^{-1}\right)^{\prime}\left(J_{\psi} \circ \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right)  \tag{4.76}\\
& +\left(A_{\bar{\omega}}^{-1}\right)^{\prime}\left(J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}-J_{\bar{\psi}} \frac{\phi}{|\nabla \psi|} \circ A_{\bar{\psi}}^{-1}\right)  \tag{4.77}\\
& +\left(A_{\bar{\omega}}^{-1}\right)^{\prime} J_{\bar{\psi}}\left(\frac{\phi}{|\nabla \psi|}-\frac{\bar{\phi}}{|\nabla \bar{\psi}|}\right) \circ A_{\bar{\psi}}^{-1}  \tag{4.78}\\
= & I+I I+I I I . \tag{4.79}
\end{align*}
$$

1 Estimates on I in terms of $\|\phi\|_{j, \alpha} \quad$ By (15.114), (4.43), (4.39), and (4.41),

$$
\begin{align*}
& \|I\|_{n, \alpha}  \tag{4.80}\\
\leq & C\left\{\left\|A_{\omega}^{-1}-A_{\bar{\omega}}^{-1}\right\|_{n+1, \alpha}\left\|J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right\|_{0, \alpha}+\left\|A_{\omega}^{-1}-A_{\bar{\omega}}^{-1}\right\|_{1, \alpha}\left\|J_{\psi} \frac{\phi}{|\nabla \psi|} \circ A_{\psi}^{-1}\right\|_{n, \alpha}^{(4.8\})}\right. \\
\leq & C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{1, \alpha}+\|F-\bar{F}\|_{2, \alpha}\left[\|\phi\|_{n, \alpha}+\|\psi-\bar{\psi}\|_{n+1, \alpha}\|\phi\|_{1, \alpha}\right]\right\}  \tag{4.82}\\
\leq & C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{1, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|\phi\|_{n, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|F-\bar{F}\|_{n, \alpha}\|\phi\|_{1, \alpha}\right\}(4.83)  \tag{4.83}\\
\leq & C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{1, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|\phi\|_{n, \alpha}\right\} . \tag{4.84}
\end{align*}
$$

2 Estimates on II in terms of $\|\phi\|_{j, \alpha} \quad$ As a preliminary we derive the following estimates for $\psi \in \mathcal{F}_{+}$in a $\|\cdot\|_{1, \alpha}$-neighborhood and any $u \in C_{\bar{\Omega}}^{\infty}$. Write $\psi_{t}=\bar{\psi}+t(\psi-\bar{\psi})$, $N_{t}=\frac{\nabla \psi_{t}}{\left|\nabla \psi_{t}\right|}$ and

$$
\begin{equation*}
\left\|J_{\psi} u \circ A_{\psi}^{-1}-J_{\bar{\psi}} u \circ A_{\bar{\psi}}^{-1}\right\|_{n, \alpha} \leq \int_{0}^{1}\left\|D_{\psi}\left(J_{\psi_{t}} u \circ A_{\psi_{t}}^{-1}\right)(\psi-\bar{\psi})\right\|_{n, \alpha} d t \tag{4.85}
\end{equation*}
$$

Using tame estimates (3.116) of the derivative of $J_{\psi} u \circ A_{\psi}^{-1}$ in $\psi$, and (4.41),

$$
\begin{align*}
& \left\|D_{\psi}\left(J_{\psi_{t}} u \circ A_{\psi_{t}}^{-1}\right)(\psi-\bar{\psi})\right\|_{n, \alpha}  \tag{4.86}\\
\leq & C\left\{\|u\|_{n+1, \alpha}\|\psi-\bar{\psi}\|_{1, \alpha}+\|u\|_{2, \alpha}\|\psi-\bar{\psi}\|_{n, \alpha}+\left\|\psi_{t}\right\|_{n+2, \alpha}\|u\|_{2, \alpha}\|\psi-\bar{\psi}\|_{1, \alpha}\right\}  \tag{4.87}\\
\leq & C\left\{\|u\|_{n+1, \alpha}\|\psi-\bar{\psi}\|_{1, \alpha}+\|u\|_{2, \alpha}\|\psi-\bar{\psi}\|_{n, \alpha}+\|\psi-\bar{\psi}\|_{n+2, \alpha}\|u\|_{2, \alpha}\right\}  \tag{4.88}\\
\leq & C\left\{\|u\|_{n+1, \alpha}\|\psi-\bar{\psi}\|_{1, \alpha}+\|\psi-\bar{\psi}\|_{n+2, \alpha}\|u\|_{2, \alpha}\right\}  \tag{4.89}\\
\leq & C\left\{\|u\|_{n+1, \alpha}\|F-\bar{F}\|_{1, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|u\|_{2, \alpha}\right\} \tag{4.90}
\end{align*}
$$

using $\left\|\psi_{t}\right\|_{n+2, \alpha} \leq\|\bar{\psi}\|_{n+2, \alpha}+\|\psi-\bar{\psi}\|_{n+2, \alpha} \leq C\left(1+\|\psi-\bar{\psi}\|_{n+2, \alpha}\right)$. Now with $u=\frac{\phi}{|\nabla \psi|}$, we have by (4.41),

$$
\begin{align*}
\|u\|_{m, \alpha} & \leq C\left\{\|\phi\|_{m, \alpha}+\|\psi-\bar{\psi}\|_{m+1, \alpha}\|\phi\|_{0, \alpha}\right\}  \tag{4.91}\\
& \leq C\left\{\|\phi\|_{m, \alpha}+\|F-\bar{F}\|_{m, \alpha}\|\phi\|_{0, \alpha}\right\} \tag{4.92}
\end{align*}
$$

so that

$$
\begin{align*}
& \|I I\|_{n, \alpha}  \tag{4.93}\\
\leq & C\left\{\|\phi\|_{n+1, \alpha}\|F-\bar{F}\|_{1, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{0, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{2, \alpha}\right\}  \tag{4.94}\\
\leq & C\left\{\|\phi\|_{n+1, \alpha}\|F-\bar{F}\|_{1, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{2, \alpha}\right\} \tag{4.95}
\end{align*}
$$

3 Estimates on I+II in terms of $\|g\|_{m, \alpha} \quad$ We may now estimate $\|I\|_{n, \alpha}+\|I I\|_{n, \alpha}$ in terms of $g$ instead of $\phi$ (the third term $I I I$ must be dealt with differently). Putting (4.84) and (4.95) together, and using interpolation inequalities (4.46), and (4.37),

$$
\begin{align*}
& \|I\|_{n, \alpha}+\|I I\|_{n, \alpha}  \tag{4.96}\\
\leq & C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{2, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|\phi\|_{n+1, \alpha}\right\}  \tag{4.97}\\
\leq & C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{1, \alpha}+\|F-\bar{F}\|_{2, \alpha}\left(\|g\|_{n-1, \alpha}+\|F-\bar{F}\|_{n, \alpha}\|g\|_{1, \alpha}\right)\right\}  \tag{4.98}\\
\leq & C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{1, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|g\|_{n-1, \alpha}\right\} \tag{4.99}
\end{align*}
$$

4 Estimates on III Since $\bar{\omega}$ is fixed, we have by (5.114)

$$
\begin{align*}
\|I I I\|_{n, \alpha} & \leq C\left\|J_{\bar{\psi}}\left\{\left(\frac{\phi}{|\nabla \psi|}-\frac{\bar{\phi}}{|\nabla \bar{\psi}|}\right)\right\} \circ A_{\bar{\psi}}^{-1}\right\|_{n, \alpha}  \tag{4.100}\\
& \leq C\left\|\left(\frac{1}{|\nabla \psi|}-\frac{1}{|\nabla \bar{\psi}|}\right) \phi\right\|_{n, \alpha}+C\left\|\frac{1}{|\nabla \bar{\psi}|}(\phi-\bar{\phi})\right\|_{n, \alpha}  \tag{4.101}\\
& =I I I_{1}+I I I_{2} . \tag{4.102}
\end{align*}
$$

Using (5.114), (4.41) and (4.37), the first term is estimated by

$$
\begin{align*}
I I I_{1} & \leq C\left\{\left\|\frac{1}{|\nabla \psi|}-\frac{1}{|\nabla \bar{\psi}|}\right\|_{0, \alpha}\|\phi\|_{n, \alpha}+\left\|\frac{1}{|\nabla \psi|}-\frac{1}{|\nabla \bar{\psi}|}\right\|_{n, \alpha}\|\phi\|_{0, \alpha}\right\}  \tag{4.103}\\
& \leq C\left\{\|\psi-\bar{\psi}\|_{2, \alpha}\|\phi\|_{n, \alpha}+\|\psi-\bar{\psi}\|_{n+2, \alpha}\|\phi\|_{0, \alpha}\right\}  \tag{4.104}\\
& \leq C\left\{\|F-\bar{F}\|_{1, \alpha}\|\phi\|_{n, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|\phi\|_{0, \alpha}\right\}  \tag{4.105}\\
& \left.\left.\leq C\left\{\|F-\bar{F}\|_{1, \alpha}\left[\|g\|_{n-2, \alpha}+\|F-\bar{F}\|_{n-1, \alpha}\|g\|_{1, \alpha}\right]+\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{1, \alpha \alpha}\right\}\right\} 106\right) \\
& \leq C\left\{\|F-\bar{F}\|_{1, \alpha}\|g\|_{n-2, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{1, \alpha}\right\} . \tag{4.107}
\end{align*}
$$

For $I I I_{2}$ we "just" need to estimate $\|\phi-\bar{\phi}\|_{n, \alpha}$. Note that

$$
\begin{equation*}
\Delta(\phi-\bar{\phi})-\bar{F}^{\prime}(\bar{\psi})(\phi-\bar{\phi})=\left(F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right) \phi+f(\psi)-\bar{f}(\bar{\psi}) . \tag{4.108}
\end{equation*}
$$

From tame estimates (2.26), (observe as well that $\bar{F}$ and $\bar{\psi}$ are fixed), we have for $j \geq 0$

$$
\begin{equation*}
\|\phi-\bar{\phi}\|_{j+2, \alpha} \leq C\left\{\left\|\left(F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right) \phi\right\|_{j, \alpha}+\|f(\psi)-\bar{f}(\bar{\psi})\|_{j, \alpha}\right\} . \tag{4.109}
\end{equation*}
$$

5 Estimates on $\|\phi-\bar{\phi}\|_{j+2, \alpha}:\left\|\left(F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right) \phi\right\|_{j, \alpha} \quad$ With the above, using (4.44), (5.114) and (5.121), the first term is estimated as follows:

$$
\begin{align*}
& \left\|\left(F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right) \phi\right\|_{j, \alpha}  \tag{4.110}\\
\leq & C\left\|F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right\|_{j, \alpha}\|\phi\|_{0, \alpha}+C\left\|F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right\|_{0, \alpha}\|\phi\|_{j, \alpha}  \tag{4.111}\\
\leq & C\left(\|F\|_{j+2, \alpha}\|\psi-\bar{\psi}\|_{0, \alpha}+\|F\|_{3, \alpha}\|\psi-\bar{\psi}\|_{j, \alpha}+\|F-\bar{F}\|_{j+1, \alpha}\right)\|\phi\|_{0, \alpha}  \tag{4.112}\\
& +C\left(\|F\|_{3, \alpha}\|\psi-\bar{\psi}\|_{0, \alpha}+\|F-\bar{F}\|_{1, \alpha}\right)\|\phi\|_{j, \alpha}  \tag{4.113}\\
\leq & C\left(\left(\|F-\bar{F}\|_{j+2, \alpha}+1\right)\|\psi-\bar{\psi}\|_{0, \alpha}+\left(\|F-\bar{F}\|_{3, \alpha}+1\right)\|\psi-\bar{\psi}\|_{j, \alpha}\right.  \tag{4.114}\\
& \left.+\|F-\bar{F}\|_{j+1, \alpha}\right)\|\phi\|_{0, \alpha}  \tag{4.115}\\
& +\left(\left(\|F-\bar{F}\|_{3, \alpha}+1\right)\|\psi-\bar{\psi}\|_{0, \alpha}+\|F-\bar{F}\|_{1, \alpha}\right)\|\phi\|_{j, \alpha}  \tag{4.116}\\
\leq & C\left\{\|F-\bar{F}\|_{j+2, \alpha}\|\phi\|_{0, \alpha}+\|F-\bar{F}\|_{3, \alpha}\|\phi\|_{j, \alpha}\right\} \tag{4.117}
\end{align*}
$$

where we have used that

$$
\begin{equation*}
\|F-\bar{F}\|_{3, \alpha}\|\psi-\bar{\psi}\|_{j, \alpha} \leq C\|F-\bar{F}\|_{3, \alpha}\|F-\bar{F}\|_{j-1, \alpha} \leq C\|F-\bar{F}\|_{0, \alpha}\|F-\bar{F}\|_{j+2, \alpha} \tag{4.118}
\end{equation*}
$$

In turn, using (4.37),

$$
\begin{align*}
& \left\|\left(F^{\prime}(\psi)-\bar{F}^{\prime}(\bar{\psi})\right) \phi\right\|_{j, \alpha}  \tag{4.119}\\
\leq & C\left\{\|F-\bar{F}\|_{j+2, \alpha}\|g\|_{1, \alpha}+\|F-\bar{F}\|_{3, \alpha}\left(\|g\|_{j-2, \alpha}+\|F-\bar{F}\|_{j-1, \alpha}\|g\|_{1, \alpha}\right)\right\}(4  \tag{4.120}\\
\leq & C\left\{\|F-\bar{F}\|_{j+2, \alpha}\|g\|_{1, \alpha}+\|F-\bar{F}\|_{3, \alpha}\|g\|_{j-2, \alpha}\right\}  \tag{4.121}\\
\leq & C\left\{\|F-\bar{F}\|_{j+2, \alpha}\|g\|_{1, \alpha}+\|F-\bar{F}\|_{1, \alpha}\|g\|_{j, \alpha}\right\} \tag{4.122}
\end{align*}
$$

using again interpolation inequalities (4.46) to get $\|F-\bar{F}\|_{3, \alpha}\|g\|_{j-2, \alpha} \leq C\|F-\bar{F}\|_{1, \alpha}\|g\|_{j, \alpha}+$ $C\|F-\bar{F}\|_{j, \alpha}\|g\|_{1, \alpha},\|F-\bar{F}\|_{3, \alpha}\|F-\bar{F}\|_{j-2, \alpha} \leq C\|F-\bar{F}\|_{2, \alpha}\|F-\bar{F}\|_{j, \alpha} \leq C\|F-\bar{F}\|_{j, \alpha}$ and to incorporate these into other terms.

6 Estimates on $\|\phi-\bar{\phi}\|_{j+2, \alpha}:\|f(\psi)-\bar{f}(\bar{\psi})\|_{j, \alpha} \quad$ Using again (4.44),

$$
\left.\|f(\psi)-\bar{f}(\bar{\psi})\|_{j, \alpha} \leq C\left\{\|f\|_{j+1, \alpha}\|\psi-\bar{\psi}\|_{0, \alpha}+\|f\|_{2, \alpha}\|\psi-\bar{\psi}\|_{j, \alpha}+\|f-\bar{f}\|_{j, \alpha}\right\} 4.123\right)
$$

Using (4.35), the first two terms in the right-hand side are estimated by

$$
\begin{align*}
& \|f\|_{j+1, \alpha}\|\psi-\bar{\psi}\|_{0, \alpha}+\|f\|_{2, \alpha}\|\psi-\bar{\psi}\|_{j, \alpha}  \tag{4.124}\\
\leq & C\left(\left(\|g\|_{j+1, \alpha}+\|F-\bar{F}\|_{j-1, \alpha}\|g\|_{1, \alpha}\right)\|F-\bar{F}\|_{1, \alpha}+\|g\|_{2, \alpha}\|F-\bar{F}\|_{j-1, \alpha}\right)  \tag{4.125}\\
\leq & C\left(\|g\|_{j+1, \alpha}\|F-\bar{F}\|_{1, \alpha}+\|F-\bar{F}\|_{j-1, \alpha}\|g\|_{2, \alpha}\right) \tag{4.126}
\end{align*}
$$

As for the last term in $\|f(\psi)-\bar{f}(\bar{\psi})\|_{j, \alpha}$, write it as

$$
\begin{equation*}
f-\bar{f}=\mathcal{E} g \circ A-\mathcal{E} g \circ \bar{A} \quad \text { with } \quad A=V\left(\mathcal{E}\left(A_{\psi}^{-1}\right)\right), \quad \bar{A}=V\left(\mathcal{E}\left(A_{\bar{\psi}}^{-1}\right)\right), \tag{4.127}
\end{equation*}
$$

to find again from (4.44)

$$
\begin{equation*}
\|f-\bar{f}\|_{j, \alpha} \leq C\left\{\|\mathcal{E} g\|_{j+1, \alpha}\|A-\bar{A}\|_{0, \alpha}+\|\mathcal{E} g\|_{2, \alpha}\|A-\bar{A}\|_{j, \alpha}\right\} \tag{4.128}
\end{equation*}
$$

In order to estimates $\|A-\bar{A}\|_{j, \alpha}$, pose $\beta=\mathcal{E}\left(A_{\psi}^{-1}\right), \bar{\beta}=\mathcal{E}\left(A_{\bar{\psi}}^{-1}\right)$, and $\beta_{t}=\bar{\beta}+t(\beta-\bar{\beta})$, and use the inversion operator $V$, see Lemma 27 in the Appendix:

$$
\begin{equation*}
A-\bar{A}=V(\beta)-V(\bar{\beta})=\int_{0}^{1} \dot{a}_{t} d t, \quad \dot{a}_{t}=D V\left(\beta_{t}\right) \cdot(\beta-\bar{\beta})=-\frac{(\beta-\bar{\beta}) \circ \beta_{t}^{-1}}{\beta_{t}^{\prime} \circ \beta_{t}^{-1}} . \tag{4.129}
\end{equation*}
$$

We have by (4.40) and (4.41)

$$
\begin{equation*}
\|\beta-\bar{\beta}\|_{m, \alpha} \leq C\left\|A_{\psi}^{-1}-A_{\bar{\psi}}^{-1}\right\|_{m, \alpha} \leq C\|\psi-\bar{\psi}\|_{m+1, \alpha} \leq C\|F-\bar{F}\|_{m, \alpha} \tag{4.130}
\end{equation*}
$$

Using the triangle inequality on $\beta_{t}=\bar{\beta}+t(\beta-\bar{\beta})$ and (4.41),

$$
\begin{equation*}
\left\|\beta_{t}\right\|_{m, \alpha} \leq C\left(1+\|\beta-\bar{\beta}\|_{m, \alpha}\right) \leq C\left\{\|F-\bar{F}\|_{m, \alpha}+1\right\} . \tag{4.131}
\end{equation*}
$$

This along with Lemma 27 implies

$$
\begin{equation*}
\left\|\beta_{t}^{-1}\right\|_{m, \alpha} \leq C\left\{\|F-\bar{F}\|_{m, \alpha}+1\right\} . \tag{4.132}
\end{equation*}
$$

In particular, $\beta_{t}$ and $\beta_{t}^{-1}$ remain in $\|\cdot\|_{2, \alpha}$-neighborhoods for $F, \bar{F} \in \mathcal{V}_{B}(\bar{F})$. Now the fraction in $\dot{a}_{t}$ is linear in the numerator, so that by (5.114) and (5.121),

$$
\begin{align*}
\left\|\dot{a}_{t}\right\|_{j, \alpha} \leq & C\left\|(\beta-\bar{\beta}) \circ \beta_{t}^{-1}\right\|_{j, \alpha}+C\left\|\beta_{t}^{\prime} \circ \beta_{t}^{-1}\right\|_{j, \alpha}\left\|(\beta-\bar{\beta}) \circ \beta_{t}^{-1}\right\|_{0, \alpha}  \tag{4.133}\\
\leq & C\left\{\|\beta-\bar{\beta}\|_{j, \alpha}+\left\|\beta_{t}^{-1}\right\|_{j, \alpha}\|\beta-\bar{\beta}\|_{1, \alpha}\right\}  \tag{4.134}\\
& +C\left(\left\|\beta_{t}\right\|_{j+1, \alpha}+\left\|\beta_{t}^{-1}\right\|_{j, \alpha}\left\|\beta_{t}\right\|_{2, \alpha}\right)\|\beta-\bar{\beta}\|_{1, \alpha}  \tag{4.135}\\
\leq & C\|F-\bar{F}\|_{j, \alpha}+C\left(\|F-\bar{F}\|_{j+1, \alpha}+1\right)\|F-\bar{F}\|_{1, \alpha}  \tag{4.136}\\
\leq & C\|F-\bar{F}\|_{j+1, \alpha} \tag{4.137}
\end{align*}
$$

from (4.130), (4.131), and (4.132). With the above, we have

$$
\begin{equation*}
\|A-\bar{A}\|_{j, \alpha} \leq \int_{0}^{1}\left\|\dot{a}_{t}\right\|_{j, \alpha} d t \leq C\|F-\bar{F}\|_{j+1, \alpha} \tag{4.138}
\end{equation*}
$$

so that by $\|\mathcal{E} g\|_{m, \alpha} \leq C\|g\|_{m, \alpha}$ (see (4.21))

$$
\begin{equation*}
\|f-\bar{f}\|_{j, \alpha} \leq C\left\{\|g\|_{j+1, \alpha}\|F-\bar{F}\|_{1, \alpha}+\|g\|_{2, \alpha}\|F-\bar{F}\|_{j+1, \alpha}\right\} . \tag{4.139}
\end{equation*}
$$

Adding this to (4.126), we find

$$
\begin{equation*}
\|f(\psi)-\bar{f}(\bar{\psi})\|_{j, \alpha} \leq C\left\{\|F-\bar{F}\|_{1, \alpha}\|g\|_{j+1, \alpha}+\|F-\bar{F}\|_{j+1, \alpha}\|g\|_{2, \alpha}\right\} . \tag{4.140}
\end{equation*}
$$

7 Conclusion on III Putting (4.122) and (4.140) together (with $n=j+2$ ) and using interpolation inequalities (4.46),

$$
\begin{equation*}
\left\|I I I_{2}\right\|_{n, \alpha} \leq C\|\phi-\bar{\phi}\|_{n, \alpha} \leq C\left\{\|F-\bar{F}\|_{n, \alpha}\|g\|_{2, \alpha}+\|F-\bar{F}\|_{1, \alpha}\|g\|_{n-1, \alpha}\right\} \tag{4.141}
\end{equation*}
$$

Adding this to the estimate (4.107) on $\left\|I I I_{1}\right\|_{n, \alpha}$ gives

$$
\begin{equation*}
\|I I I\|_{n, \alpha} \leq C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{2, \alpha}+\|F-\bar{F}\|_{1, \alpha}\|g\|_{n-1, \alpha}\right\} . \tag{4.142}
\end{equation*}
$$

8 Conclusion on Lemma 17 Putting (4.99) and (4.142) together, we have

$$
\begin{equation*}
\|K(F) g-K(\bar{F}) g\|_{n, \alpha} \leq C\left\{\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{2, \alpha}+\|F-\bar{F}\|_{2, \alpha}\|g\|_{n-1, \alpha}\right\} . \tag{4.143}
\end{equation*}
$$

Proposition $18(\mathbf{I d}+K(F)$ is invertible for $F$ near $\bar{F})$ Suppose $\bar{\omega}=\bar{F}(\bar{\psi})$ satisfies (ND2). Then, there exists $a\|\cdot\|_{3, \alpha}$-neighborhood

$$
\begin{equation*}
\mathcal{V}_{I}(\bar{F}) \subset \mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F}) \tag{4.144}
\end{equation*}
$$

of $\bar{F}$ such that $h=g+K(F) g$ has a tame family of inverses $g=V M(F) h$ satisfying

$$
\begin{equation*}
\|g\|_{n, \alpha} \leq C\left\{\|h\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|h\|_{2, \alpha}\right\}, \quad n \geq 2 \tag{4.145}
\end{equation*}
$$

Remark The Nash-Moser Inverse Function Theorem only requires a continuous and tame inverse. However, with a little more work, one can show that $g=V M(F) h$ is continuously differentiable, hence smooth tame by Theorem 5.3.1, p. 102, Part I, and Theorem 3.1.1, p. 150, Part II of [7].

## Proof

1 The estimate $\|g\|_{2, \alpha} \leq C\|g+K(\bar{F}) g\|_{2, \alpha}$ holds We first show that (ND2) implies that $\operatorname{Id}_{C_{[0,|\Omega|]}^{\infty}}+K(\bar{F})$ has trivial kernel as a map $C_{[0,|\Omega|]}^{\infty} \rightarrow C_{[0,|\Omega|]}^{\infty}$. Suppose $g \in C_{[0,|\Omega|]}^{\infty}$ is in the kernel of $\operatorname{Id}_{C_{[0,|\Omega|]}^{\infty}}+K(\bar{F}): C_{[0,|\Omega|]}^{\infty} \rightarrow C_{[0,|\Omega|]}^{\infty}$. Set $f=V B(\bar{F}) \cdot g$ so that $D T(\bar{F}) \cdot f=0$ and thus $\nu=\{\bar{\omega}, \alpha\}$ for some $\alpha \in \mathcal{U}$, see Proposition 8 . Then $f=0$ on range $(\bar{\psi})$, precisely by the non-degeneracy condition (ND2) and in turn $g=B(\bar{F}) \cdot V B(\bar{F}) \cdot g=B(\bar{F}) \cdot f=0$, i.e. $\operatorname{Id}_{C_{[0,|\Omega|]}^{\infty}}+K(\bar{F}): C_{[0,|\Omega|]}^{\infty} \rightarrow C_{[0,|\Omega|]}^{\infty}$ has trivial kernel.

Next, we show that $\operatorname{Id}_{C_{[0,|\Omega|]}^{2, \alpha}}+K(\bar{F}): C_{[0,|\Omega|]}^{2, \alpha} \rightarrow C_{[0,|\Omega|]}^{2, \alpha}$, satisfies

$$
\begin{equation*}
\|g\|_{2, \alpha} \leq C\|g+K(\bar{F}) g\|_{2, \alpha} . \tag{4.146}
\end{equation*}
$$

First observe from the tame estimates on $K(F) g$ derived in the proof of Lemma 16 that $K(\bar{F})$ maps $C_{[0,|\Omega|]}^{n, \alpha}$ into $C_{[0,|\Omega|]}^{n+1, \alpha}$ for each $n \geq 2$. Let $g \in C_{[0,|\Omega|]}^{2, \alpha}$ be in the kernel of $\operatorname{Id}_{C_{[0,|\Omega|]}^{2, \alpha}}+K(\bar{F}): C_{[0,|\Omega|]}^{2, \alpha} \rightarrow C_{[0,|\Omega|]}^{2, \alpha}$. Then $g=-K(\bar{F}) g \in C_{[0,|\Omega|]}^{3, \alpha}$ and, repeating, $g \in C_{[0,|\Omega|]}^{\infty}$. Hence $g=0$ which shows that $\operatorname{Id}_{C_{[0,|\Omega|]}^{2, \alpha}}+K(\bar{F}): C_{[0,|\Omega|]}^{2, \alpha} \rightarrow C_{[0,|\Omega|]}^{2, \alpha}$ has trivial kernel. By the Fredholm alternative (for Banach spaces), $\operatorname{Id}_{C_{[0,|\Omega|]}^{2, \alpha}}+K(\bar{F}): C_{[0,|\Omega|]}^{2, \alpha} \rightarrow C_{[0,|\Omega|]}^{2, \alpha}$ is an isomorphism satisfying (4.146).

2 Invertibility for $F$ near $\bar{F} \quad$ From Lemma 17

$$
\begin{align*}
\|g\|_{2, \alpha} & \leq C\|g+K(\bar{F}) g\|_{2, \alpha}  \tag{4.147}\\
& \leq C\left(\|g+K(F) g\|_{2, \alpha}+\|K(F) g-K(\bar{F}) g\|_{2, \alpha}\right)  \tag{4.148}\\
& \leq C\left(\|g+K(F) g\|_{2, \alpha}+\|F-\bar{F}\|_{3, \alpha}\|g\|_{2, \alpha}\right) . \tag{4.149}
\end{align*}
$$

Define now $\mathcal{V}_{I}(\bar{F}) \subset \mathcal{V}_{B}(\bar{F}) \subset \mathcal{V}_{S}(\bar{F})$ to be a sufficiently small $\|\cdot\|_{3, \alpha}$-neighborhood of $\bar{F}$ so that the last term can be incorporated into the left-hand side:

$$
\begin{equation*}
\|g\|_{2, \alpha} \leq C\|g+K(F) g\|_{2, \alpha} \tag{4.150}
\end{equation*}
$$

Now from (4.66) we have for $n \geq 2$

$$
\begin{align*}
\|g\|_{n, \alpha} & \leq C\left(\|g+K(F) g\|_{n, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|g\|_{1, \alpha}\right)  \tag{4.151}\\
& \leq C\left(\|g+K(F) g\|_{n, \alpha}+\|F-\bar{F}\|_{n+1, \alpha}\|g+K(\bar{F}) g\|_{2, \alpha}\right) \tag{4.152}
\end{align*}
$$

which implies the estimates (4.145). That is, $h=M(F) g=g+K(F) g$ is a family of invertible linear maps of Fréchet spaces for $F \in \mathcal{V}_{I}(\bar{F})$ with tame inverse denoted $g=V M(F) h$.

3 Continuity Let $F, \tilde{F} \in \mathcal{V}_{I}(\bar{F}), g, \tilde{g} \in C_{[0,|\Omega|]}^{\infty}$, and set $h=g+K(F) g, \tilde{h}=\tilde{g}+K(\tilde{F}) \tilde{g}$. Then,

$$
\begin{equation*}
(g-\tilde{g})+K(\tilde{F})(g-\tilde{g})=h-\tilde{h}-(K(F) g-K(\tilde{F}) g) \tag{4.153}
\end{equation*}
$$

and from (4.152) we deduce that

$$
\begin{align*}
\|g-\tilde{g}\|_{n, \alpha} \leq & C\left\{\|h-\tilde{h}\|_{n, \alpha}+\|K(F) g-K(\tilde{F}) g\|_{n, \alpha}\right.  \tag{4.154}\\
& \left.+\|F-\tilde{F}\|_{n+1, \alpha}\left(\|h-\tilde{h}\|_{2, \alpha}+\|K(F) g-K(\bar{F}) g\|_{2, \alpha}\right)\right\} \tag{4.155}
\end{align*}
$$

Now from the tame estimates (4.145) it is clear that $\|g\|_{n, \alpha}$ remains bounded as $\|h-\tilde{h}\|_{n, \alpha}$ and $\|F-\tilde{F}\|_{n+1, \alpha}$ tend to zero. In turn the estimates (4.74) (also valid for $\left.K(F) g-K(\tilde{F}) g\right)$ show that $\|K(F) g-K(\tilde{F}) g\|_{n, \alpha}$ tends to zero as well, and thus clearly $\|g-\tilde{g}\|_{n, \alpha}$ tends to zero.

### 4.4 Injective part of Theorem 1

The injective part of Theorem 1 requires a modification of the injective part of the NashMoser theorem since the map $T$ cannot be injective: defining $T$ on $C_{I}^{\infty}$ where $I \supset \operatorname{range}(\bar{\psi})$ (see (4.3)), changing $F$ outside the range of the corresponding solution $\Delta \psi=F(\psi)$ clearly does not affect this solution.

Proposition 19 There exists a $C^{2}$-neighborhood of $\bar{F}$ such that, if $F, \tilde{F}$ are such that $T(F)=T(\tilde{F})$, then the corresponding solutions $\psi=S(F)$ and $\tilde{\psi}=S(\tilde{F})$ are the same.

Proof Let $F_{n} \neq \tilde{F}_{n} \rightarrow_{n} \bar{F}$ in $C^{2}$ such that $T\left(F_{n}\right)=T\left(\tilde{F}_{n}\right)$, and let

$$
\begin{equation*}
\tilde{F}_{n}-F_{n}=\epsilon_{n} G_{n}, \quad \psi_{n}=S\left(F_{n}\right), \quad \tilde{\psi}_{n}=S\left(\tilde{F}_{n}\right), \quad \tilde{\psi}_{n}-\psi_{n}=\epsilon_{n} v_{n} \tag{4.157}
\end{equation*}
$$

where $\epsilon_{n}$ is to be chosen appropriately. Assume without loss of generality that range $\left(\psi_{n}\right) \subset$ range $\left(\tilde{\psi}_{n}\right)$. Supposing that we can normalize according to

$$
\begin{equation*}
\left\|G_{n}\right\|_{C_{\text {range }\left(\bar{\psi}_{n}\right)}^{0}}=1, \tag{4.158}
\end{equation*}
$$

we will arrive at a contradiction, thus proving our claim.
Remark One might expect to normalize in the $C^{2}$-norm since the $F$ 's converge in that norm. However, (strong) compactness fails in infinite dimensions. On the other hand, one can use the gain of regularity provided by $\Delta \psi=F(\psi)$.

Changing $F_{n}$ and $\tilde{F}_{n}$ outside of range $\left(\tilde{\psi}_{n}\right)$ does not affect $\psi_{n}$ nor $\tilde{\psi}_{n}$, and in turn the assumption $T\left(\tilde{F}_{n}\right)=T\left(F_{n}\right)$ is preserved. Therefore, we may adjust $F_{n}$ and $\tilde{F}_{n}$ in such a way that, without loss of generality, for the first derivatives we have the bounds

$$
\begin{equation*}
\left\|G_{n}\right\|_{C_{[\bar{c}, 0]}^{1}} \leq 2\left\|G_{n}\right\|_{C_{\text {range }}^{1}\left(\tilde{\psi}_{n}\right)} \tag{4.159}
\end{equation*}
$$

Since $T\left(\tilde{F}_{n}\right)=T\left(F_{n}\right)$, we have $\epsilon_{n} G_{n}=F_{n} \circ A_{\psi_{n}}^{-1} \circ A_{\tilde{\psi}_{n}}-F_{n}$ hence

$$
\begin{equation*}
\epsilon_{n} G_{n}(\lambda)=\left(\int_{0}^{1} F_{n}^{\prime}\left(\lambda+t\left(A_{\psi_{n}}^{-1}\left(A_{\tilde{\psi}_{n}}(\lambda)\right)-\lambda\right)\right) d t\right)\left(A_{\psi_{n}}^{-1}\left(A_{\tilde{\psi}_{n}}(\lambda)\right)-\lambda\right), \quad \lambda \in \operatorname{range}\left(\tilde{\psi}_{n}\right) \tag{4.160}
\end{equation*}
$$

Letting $d_{n}$ denote the integral factor, we write this as

$$
\begin{equation*}
\epsilon_{n} G_{n}(\lambda)=d_{n}(\lambda)\left(A_{\psi_{n}}^{-1}-A_{\tilde{\psi}_{n}}^{-1}\right)\left(A_{\tilde{\psi}_{n}}(\lambda)\right), \quad \lambda \in \operatorname{range}\left(\tilde{\psi}_{n}\right) . \tag{4.161}
\end{equation*}
$$

Now $F_{n}$ and $\tilde{F}_{n}$ are bounded in $C^{2}, F \in C^{2} \mapsto \psi \in C^{2}$ is continuous by (2.58), and $\psi \in C^{2} \mapsto A_{\psi}^{-1} \in C^{1}$ is continuous by Lemma 12. Thus, $d_{n}$ and $A_{\tilde{\psi}_{n}}$ are bounded in $C^{1}$ and in turn

$$
\begin{equation*}
\epsilon_{n}\left\|G_{n}\right\|_{C_{\text {range }\left(\tilde{\psi}_{n}\right)}^{1}} \leq C\left\|A_{\psi_{n}}^{-1}-A_{\tilde{\psi}_{n}}^{-1}\right\|_{1} . \tag{4.162}
\end{equation*}
$$

Write

$$
\begin{equation*}
A_{\tilde{\psi}_{n}}^{-1}-A_{\psi_{n}}^{-1}=\epsilon_{n}\left(\int_{0}^{1} D Q\left(\psi_{n}+t\left(\tilde{\psi}_{n}-\psi_{n}\right)\right) d t\right) v_{n} \tag{4.163}
\end{equation*}
$$

where $\psi_{n}$ and $\tilde{\psi}_{n}$ are bounded in $C^{2}$ since $F_{n}$ and $\tilde{F}_{n}$ converge in $C^{2}$. The proof of Lemma 14 easily shows that (4.38) holds in the $\|\cdot\|_{n}$-grading, so that

$$
\begin{equation*}
\epsilon_{n}\left\|G_{n}\right\|_{C_{\text {range }\left(\tilde{\psi}_{n}\right)}^{1}} \leq C\left\|A_{\tilde{\psi}_{n}}^{-1}-A_{\psi_{n}}^{-1}\right\|_{1} \leq \epsilon_{n} C\left\|v_{n}\right\|_{1} . \tag{4.164}
\end{equation*}
$$

Finally, rewrite $\Delta \tilde{\psi}_{n}-\Delta \psi_{n}=\tilde{F}_{n}\left(\tilde{\psi}_{n}\right)-F_{n}\left(\psi_{n}\right)$ as

$$
\begin{equation*}
\epsilon_{n} \Delta v_{n}=\epsilon_{n} G_{n}\left(\tilde{\psi}_{n}\right)+\left(F_{n}\left(\tilde{\psi}_{n}\right)-F_{n}\left(\psi_{n}\right)\right) \tag{4.165}
\end{equation*}
$$

This can be written as $\Delta v_{n}+c_{n} v_{n}=G_{n}\left(\tilde{\psi}_{n}\right)$

$$
\begin{equation*}
-c_{n}=\left(\int_{0}^{1} F_{n}^{\prime}\left(\psi_{n}+t\left(\tilde{\psi}_{n}-\psi_{n}\right)\right) d t\right) v_{n} \tag{4.166}
\end{equation*}
$$

$\bar{F}$ satisfies (ND1) hence for large $n, c_{n}$ is sufficiently close to $\bar{F}^{\prime}(\bar{\psi})$ and in turn $\Delta+c_{n}$ is invertible with

$$
\begin{equation*}
\left\|v_{n}\right\|_{1} \leq C\left\|G_{n}\right\|_{0} \tag{4.167}
\end{equation*}
$$

Putting (4.159), (4.164), and (4.167) together we obtain

$$
\begin{equation*}
\left\|G_{n}\right\|_{1} \leq 2\left\|G_{n}\right\|_{C_{\text {range }\left(\tilde{\psi}_{n}\right)}^{1}} \leq C\left\|v_{n}\right\|_{1} \leq C\left\|G_{n}\right\|_{0} \leq C \tag{4.168}
\end{equation*}
$$

Passing to a subsequence, we may assume that

$$
\begin{equation*}
G_{n} \rightarrow_{n} G \quad \text { in } \quad C^{0} \tag{4.169}
\end{equation*}
$$

Lemma 12 says that $(\omega, \nu) \mapsto D Q(\omega) \nu$ is continuous as an operator $C^{2} \times C^{0} \rightarrow C^{0}$, (2.103) shows that $F \mapsto \omega=F(\psi)$ is continuous $C^{2} \rightarrow C^{2}$, while $(F, f) \mapsto \nu=\Delta \phi=$ $F^{\prime}(\psi) \phi+f(\psi)$ is continuous $C^{2} \times C^{0} \rightarrow C^{0}$ by (2.104). In conclusion, we have the continuous operator

$$
\left\{\begin{array}{ccccc}
C^{2} & \times & C^{0} & & C^{0}  \tag{4.170}\\
F & & f & & D T(F) f
\end{array}\right\}
$$

Taking limits in

$$
\begin{equation*}
0=\frac{1}{\epsilon_{n}}\left(T\left(F_{n}+\epsilon_{n} G_{n}\right)-T\left(F_{n}\right)\right)=\int_{0}^{1} D T\left(F_{n}+\epsilon_{n}\left(\tilde{F}_{n}-F_{n}\right)\right) G_{n} d t \tag{4.171}
\end{equation*}
$$

one finds $D T(\bar{F}) G=0$. This means that $G$ vanishes on range $(\bar{\psi})$, contradicting the normalization $\|G\|_{C_{\text {range }(\bar{\psi})}^{0}}=1$ guaranteed by (4.158).

## 5 Appendix: the Nash-Moser Inverse Function Theorem

Inverse function theorems express the fact that nonlinear problems are as solvable as their linearizations: a nonlinear map $T$ is (locally) surjective where its first derivative $D T$ is surjective, and $T$ is (locally) injective where $D T$ is injective. In case $T:(B \subset X) \rightarrow Y$ is a sufficiently smooth (e.g. twice continuously differentiable) map of Banach spcaes, Newton's
scheme constructs successive approximations which converge very rapidly. This "accelerated convergence" is visible through an estimate of the form $x_{n+1} \leq x_{n}^{2}$. For comparison, a proof by the Picard approximation method would involve an iteration of the form $x_{n+1} \leq \lambda x_{n}$ with some fixed $0<\lambda<1$.

Loss of derivatives occurs when, for instance, a surjective first derivative $D T(F)$ uses a number of derivatives, which are not recovered by its right-inverse $L(F)$. (This is the case of our map $T(F)=A_{\omega}^{-1}$.) The Newton algorithm can no longer be implemented as such. On the other hand, its accelerated convergence indicates that it should tolerate some adjustments made in order to overcome loss of derivatives.

In [15] Nash introduced smoothing operators for his solution to the isometric imbedding problem of Riemannian manifolds into Euclidean spaces. From [15] Moser extracted a simple algorithm solving an inverse function problem even when loss of derivatives occurs [14]. We will refer to this algorithm as the Moser scheme. It is a modified Newton scheme where the smoothing operators of Nash introduce an error term having no effect on the convergence of the algorithm provided the maps satisfy certain ("tame") estimates. The solution so obtained is "rough" a priori (e.g. if one works with spaces of functions, the solution may have fewer derivatives than the formulation of the problem actually allows). In a second step, one verifies that this rough solution is in fact smooth. This again uses the "tameness" of certain operators, as well as interpolation inequalities available in "tame Fréchet-spaces" in a crucial way.

Various extensions and improvements have been developed subsequently. In particular, Hamilton introduced in [7] the tame Fréchet category (see Sections II. 1 and II.2), essentially that introduced by Sergeraert in [19], in which the modified Newton algorithm is applicable and therefore an Inverse Function Theorem holds. That is, an inverse function exists and lives in the tame Fréchet category.

We emphasize that the Moser scheme is used to construct a rough solution to $T(x)=y$ when the map $T$ has surjective first derivative ( $D T$ is not required to be injective.) That this solution is smooth is a consequence of the interpolation inequalities available on "tame Fréchet-spaces", and surely the estimates on the successive approximations play a part in the proof. This is the surjective part of the Inverse Function Theorem. In case $T$ has injective first derivative $D T$, then the interpolation inequalities (1.39) again show that $T$ is injective as well. The Moser scheme plays no rôle in this injective part of the Inverse Function Theorem. Our map $T(F)=A_{\omega}^{-1}$ cannot be injective (see the discussion in the Introduction). Nevertheless, the injective part of the Inverse Function Theorem for tame Fréchet spaces gives the idea for the proof of the injective part of Theorem 1 .

### 5.1 Tame estimates on $T(F)=A_{\omega}^{-1}$

In this Section we derive the precise tame estimates on $T(F)$ so as to set the parameters for the proof of the (existence part of the) Nash-Moser Inverse Function Theorem, see Theorem [21,

Recall from Proposition 18 that $\mathcal{V}_{I}(\bar{F})$ is a $\|\cdot\|_{3, \alpha}$-neighborhood of $\bar{F}$.
Proposition 20 For $F \in V_{I}(\bar{F})$, any $f_{1}, f_{2} \in C_{I}^{\infty}$, and $n \geq 0$,

$$
\begin{align*}
\|T(F)\|_{n, \alpha} \leq & C\left(\|F\|_{n, \alpha}+1\right)  \tag{5.1}\\
\|D T(F) f\|_{n, \alpha} \leq & C\left(\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}\right)  \tag{5.2}\\
\left\|D^{2} T(F)\left(f_{1}, f_{2}\right)\right\|_{n, \alpha} \leq & C\left(\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{2, \alpha}+\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}\right.  \tag{5.3}\\
& \left.+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right), \tag{5.4}
\end{align*}
$$

and for $n \geq 2$ and $h \in C_{[0,|\Omega|]}^{\infty}$,

$$
\begin{equation*}
\|L(F) h\|_{n, \alpha} \leq C\left(\|h\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|h\|_{2, \alpha}\right) \tag{5.5}
\end{equation*}
$$

In the proof of the Nash-Moser theorem below, we will shift the indices in the norms and use the notation

$$
\begin{equation*}
|\cdot|_{n}:=\|\cdot\|_{n+2, \alpha}, \quad n \geq 0 \tag{5.6}
\end{equation*}
$$

Note that $\mathcal{V}_{I}(\bar{F})$ is then a $|\cdot|_{1}$-neighborhood.
Proof With $T(F)=Q(\omega)$ where $\omega=\Delta \psi$ and $\psi=S(F)$, from Propositions 11 and 6, we have for $n \geq 0$ and $F$ in the $\|\cdot\|_{2, \alpha}$-neighborhood $\mathcal{V}_{S}(\bar{F})$, see Proposition 6,

$$
\begin{equation*}
\|T(F)\|_{n, \alpha} \leq C\left(\|\omega\|_{n, \alpha}+1\right) \leq C\left(\|\psi\|_{n+2, \alpha}+1\right) \leq C\left(\|F\|_{n, \alpha}+1\right) \tag{5.7}
\end{equation*}
$$

Write the first derivative as $D T(F) f=D Q(\omega) \nu$ where $\nu=\Delta \phi=F^{\prime}(\psi) \phi+f(\psi)$. Again from Propositions 11 and 6, we have for $n \geq 0, F \in \mathcal{V}_{S}(\bar{F})$, and any $f \in C_{I}^{\infty}$,

$$
\begin{align*}
& \|D T(F) f\|_{n, \alpha}  \tag{5.8}\\
\leq & C\left(\|\nu\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\|\nu\|_{1, \alpha}\right)  \tag{5.9}\\
\leq & C\left(\|\phi\|_{n+2, \alpha}+\left(\|F\|_{n+1, \alpha}+1\right)\|\phi\|_{3, \alpha}\right)  \tag{5.10}\\
\leq & C\left(\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}+\left(\|F\|_{n+1, \alpha}+1\right)\left(\|f\|_{1, \alpha}+\|F\|_{2, \alpha}\|f\|_{1, \alpha}\right)\right.  \tag{5.11}\\
\leq & C\left(\|f\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|f\|_{1, \alpha}\right) \tag{5.12}
\end{align*}
$$

Write the second derivative as

$$
\begin{equation*}
D^{2} T(F)\left(f_{1}, f_{2}\right)=D^{2} Q(\omega)\left(\nu_{1}, \nu_{2}\right)+D Q(\omega) \nu_{12} \tag{5.13}
\end{equation*}
$$

where

$$
\begin{gather*}
\nu_{1}=\Delta \phi_{1}=F^{\prime}(\psi) \phi_{1}+f_{1}(\psi),  \tag{5.14}\\
\nu_{2}=\Delta \phi_{2}=F^{\prime}(\psi) \phi_{2}+f_{2}(\psi),  \tag{5.15}\\
\nu_{12}=\Delta \phi_{12}=F^{\prime}(\psi) \phi_{12}+F^{\prime}(\psi) \phi_{1} \phi_{2}+f_{1}^{\prime}(\psi) \phi_{2}+f_{2}^{\prime}(\psi) \phi_{1}, \tag{5.16}
\end{gather*}
$$

see Proposition 6. Thus, for $F$ in the $\|\cdot\|_{3, \alpha}$-neighborhood $\mathcal{V}_{I}(\bar{F})$, see Proposition 18, and any $f_{1}, f_{2} \in C_{I}^{\infty}$, we have for $n \geq 0$

$$
\begin{align*}
& \left\|D^{2} T(F)\left(f_{1}, f_{2}\right)\right\|_{n, \alpha}  \tag{5.17}\\
\leq & C\left\{\left\|\nu_{1}\right\|_{n+1, \alpha}\left\|\nu_{2}\right\|_{1, \alpha}+\left\|\nu_{1}\right\|_{1, \alpha}\left\|\nu_{2}\right\|_{n+1, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\nu_{1}\right\|_{2, \alpha}\left\|\nu_{2}\right\|_{2, \alpha}\right.  \tag{5.18}\\
& \left.+\left\|\nu_{12}\right\|_{n, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\nu_{12}\right\|_{1, \alpha}\right\}  \tag{5.19}\\
\leq & C\left\{\left\|\phi_{1}\right\|_{n+3, \alpha}\left\|\phi_{2}\right\|_{3, \alpha}+\left\|\phi_{1}\right\|_{3, \alpha}\left\|\phi_{2}\right\|_{n+3, \alpha}+\|\omega\|_{n+2, \alpha}\left\|\phi_{1}\right\|_{4, \alpha}\left\|\phi_{2}\right\|_{4, \alpha}\right.  \tag{5.20}\\
& \left.+\left\|\phi_{12}\right\|_{n+2, \alpha}+\|\omega\|_{n+1, \alpha}\left\|\phi_{12}\right\|_{3, \alpha}\right\}  \tag{5.21}\\
\leq & C\left\{\left(\left\|f_{1}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{1, \alpha}\right)\left(\left\|f_{2}\right\|_{1, \alpha}+\|F\|_{2, \alpha}\left\|f_{2}\right\|_{1, \alpha}\right)\right.  \tag{5.22}\\
& +\left(\left\|f_{1}\right\|_{1, \alpha}+\|F\|_{2, \alpha}\left\|f_{1}\right\|_{1, \alpha}\right)\left(\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{2}\right\|_{1, \alpha}\right)  \tag{5.23}\\
& +\left(\|F\|_{n+2, \alpha}+1\right)\left(\left\|f_{1}\right\|_{2, \alpha}+\|F\|_{3, \alpha}\left\|f_{1}\right\|_{1, \alpha}\right)\left(\left\|f_{2}\right\|_{2, \alpha}+\|F\|_{3, \alpha}\left\|f_{2}\right\|_{1, \alpha}\right)  \tag{5.24}\\
& +\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{1, \alpha}+\left\|f_{1}\right\|_{1, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}  \tag{5.25}\\
& \left.+\left(\|F\|_{n+1, \alpha}+1\right)\left(\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}+\|F\|_{3, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right)\right\}  \tag{5.26}\\
\leq & C\left\{\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{2, \alpha}+\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right\}  \tag{5.27}\\
\leq & C\left\{\left\|f_{1}\right\|_{n+1, \alpha}\left\|f_{2}\right\|_{2, \alpha}+\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{n+1, \alpha}+\|F\|_{n+2, \alpha}\left\|f_{1}\right\|_{2, \alpha}\left\|f_{2}\right\|_{2, \alpha}\right\} . \tag{5.28}
\end{align*}
$$

Finally we compute the tame estimates for the right-inverse $L(F) h$ to $D T(F) f$. Recall that it is given by $f=L(F) h=V B(F) \cdot V M(F) h$. With the tame estimates on $f=V B(F) g$
from Lemma 15 and those on $g=V M(F) h$ from Proposition 18, we deduce that, for $n \geq 2$,

$$
\begin{align*}
\|f\|_{n, \alpha} & \leq C\left\{\|g\|_{n, \alpha}+\|F\|_{n-2, \alpha}\|g\|_{1, \alpha}\right\}  \tag{5.29}\\
& \leq C\left\{\|h\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|h\|_{2, \alpha}+\|F\|_{n-2, \alpha}\left(\|h\|_{1, \alpha}+\|F\|_{2, \alpha}\|h\|_{2, \alpha}\right\}\right.  \tag{5.30}\\
& \leq C\left\{\|h\|_{n, \alpha}+\|F\|_{n+1, \alpha}\|h\|_{2, \alpha}\right\} \tag{5.31}
\end{align*}
$$

### 5.2 Surjective part of the Nash-Moser Inverse Function Theorem

Our presentation of the Nash-Moser Inverse Function Theorem is a blend of [7], [14], and [19]. We will limit ourselves to constructing a right-inverse, since in our application to Theorem 1 we do not need further properties of this right-inverse (smoothness and tameness). We refer to Section III. 1 of [7] for further details.

Consider $\mathcal{X}, \mathcal{Y}$ tame Fréchet spaces, with smoothing operators $S(t), t>0$, satisfying the estimates (1.38) described in the Introduction, and set

$$
\begin{equation*}
\mathcal{B}:=\left\{u \in \mathcal{X}| | u-\left.\bar{u}\right|_{1}<\eta\right\} . \tag{5.32}
\end{equation*}
$$

Let $T:(\mathcal{B} \subset \mathcal{X}) \rightarrow \mathcal{Y}$ such that for any $u \in \mathcal{B}, v_{1}, v_{2} \in \mathcal{X}, h \in \mathcal{Y}$, and $n \geq 0$,

$$
\begin{align*}
|T(u)|_{n} & \leq C\left(|u-\bar{u}|_{n}+1\right)  \tag{5.33}\\
|D T(u) v|_{n} & \leq\left(|v|_{n}+|u-\bar{u}|_{n+1}|v|_{0}\right)  \tag{5.34}\\
\left|D^{2} T(u)\left(v_{1}, v_{2}\right)\right|_{n} & \leq C\left(\left|v_{1}\right|_{n+1}\left|v_{2}\right|_{0}+\left.\left.\left|v_{1}\right|_{0}\right|_{2}\right|_{n+1}+|u-\bar{u}|_{n+2}\left|v_{1}\right|_{0}\left|v_{2}\right|_{0}\right)  \tag{5.35}\\
|L(u) h|_{n} & \leq C\left(|h|_{n}+|u-\bar{u}|_{n+1}|h|_{0}\right) . \tag{5.36}
\end{align*}
$$

(It is clear that our map $T(F)=A_{\omega}^{-1}$ satisfies these conditions where $|\cdot|_{n}=\|\cdot\|_{n+2, \alpha}$.) Suppose given a solution $T(\bar{u})=\bar{g}$.

## Theorem 21 (Existence part of the Nash-Moser Inverse Function Theorem)

There is a neighborhood $\mathcal{G} \subset \mathcal{Y}$ of $\bar{g}$ in which $T(u)=g$ has a solution $u \in \mathcal{B}$ for any $g \in \mathcal{G}$.
Remark The neighborhood $\mathcal{G}$ is defined in (5.80) in terms of a parameter $j$ given in (5.96). In particular, $\mathcal{G}$ is a $|\cdot|_{8}=\|\cdot\|_{10, \alpha}$-neighborhood. Further, since $\omega \mapsto A_{\omega}^{-1}$ is continuous $C^{11} \rightarrow C^{11} \hookrightarrow C^{10, \alpha}$, see Proposition 11, the $C^{\infty}$-neighborhood of $\bar{\omega}$ in which each co-adjoint orbit has a steady-state can then be taken as a $\|\cdot\|_{11}$-neighborhood.

## Proof

1 The modified Newton scheme For $g \in \mathcal{Y}$, we pose

$$
\begin{equation*}
P(g, u)=T(u)-g \tag{5.37}
\end{equation*}
$$

and the problem is to solve $P(g, u)=0$. We will think of $g$ as a parameter. Then $P$ satisfies the following estimates for any $u \in \mathcal{B}, v_{1}, v_{2} \in \mathcal{X}$, and $n \geq 0$,

$$
\begin{align*}
|P(g, u)|_{n} & \leq C\left(|u-\bar{u}|_{n}+|g-\bar{g}|_{n}+1\right)  \tag{5.38}\\
\left|D_{u} P(g, u) v\right|_{n} & \leq\left(|v|_{n}+|u-\bar{u}|_{n+1}|v|_{0}\right)  \tag{5.39}\\
\left|D_{u u}^{2} P(g, u)\left(v_{1}, v_{2}\right)\right|_{n} & \leq C\left(\left|v_{1}\right|_{n+1}\left|v_{2}\right|_{0}+\left|v_{1}\right|_{0}\left|v_{2}\right|_{n+1}+\left.\left.|u-\bar{u}|_{n+2}\left|v_{1}\right|_{0}\right|_{2}\right|_{0}\right) \tag{5.40}
\end{align*}
$$

and $v=L(u) h$ is again a right-inverse to $h=D_{u} P(g, u) v$ :

$$
\begin{equation*}
D_{u} P(g, u) L(u) h=h \tag{5.41}
\end{equation*}
$$

The solution is constructed by the Moser scheme, which is a modified Newton scheme:

$$
\begin{equation*}
u_{n+1}-u_{n}:=-S\left(t_{n}\right) L\left(u_{n}\right) P\left(u_{n}\right), \quad n \geq 0, \quad u_{0}:=\bar{u}, \quad t_{n}:=A^{\kappa^{n}} \tag{5.42}
\end{equation*}
$$

for some $A>1$ and $0<\kappa<2$ to be determined. Fix $j \geq 1$ which will be specified later. Let $M, M_{j}>1$ be constants such that for all $u$, $w$ such that $u, u+w \in \mathcal{B}$, any $v \in \mathcal{X}$, any $h \in \mathcal{Y}$, and $t>0$, we have:

$$
\begin{align*}
|S(t) v|_{1} & \leq M t|v|_{0}  \tag{5.43}\\
|S(t) v|_{j} & \leq M_{j} t|v|_{j-1}  \tag{5.44}\\
|v-S(t) v|_{0} & \leq M_{j} t^{1-j}|v|_{j-1}  \tag{5.45}\\
\left|D_{u} P(u) v\right|_{0} & \leq M|v|_{0}  \tag{5.46}\\
\left|P(u+w)-P(u)-D_{u} P(u) w\right|_{0} & \leq M|w|_{1}^{2}  \tag{5.47}\\
|L(u) h|_{0} & \leq M|h|_{0} \tag{5.48}
\end{align*}
$$

(That these hold is immediate from the tame estimates on $T$, the estimates (1.38) on the smoothing operators $S(t)$, and Taylor's expansion with remainder (1.37) given in the Introduction). Next, with the first requirement that $\mathcal{G}$ be contained in a neighborhood of the form

$$
\begin{equation*}
\mathcal{G} \subset\left\{|g-\bar{g}|_{j-1}<C\right\} \subset \mathcal{Y} \tag{5.49}
\end{equation*}
$$

and increasing $M_{j}$ if necessary, we have for any $u \in \mathcal{B}$ and $g \in \mathcal{G}$

$$
\begin{equation*}
|L(u) P(g, u)|_{j-1} \leq M_{j}\left(1+|u-\bar{u}|_{j}\right) \tag{5.50}
\end{equation*}
$$

which holds since $|u|_{0}$ and $|g|_{0}$ remain bounded and

$$
\begin{align*}
|L(u) P(g, u)|_{j-1} & \leq C\left(|P(g, u)|_{j-1}+|u-\bar{u}|_{j}|P(g, u)|_{0}\right)  \tag{5.51}\\
& \leq C\left(|u-\bar{u}|_{j-1}+|g-\bar{g}|_{j-1}+1+|u-\bar{u}|_{j}\right) . \tag{5.52}
\end{align*}
$$

Write now

$$
\begin{align*}
\left|u_{n+1}-u_{n}\right|_{1} & =\left|S\left(t_{n}\right) L\left(u_{n}\right) P\left(g, u_{n}\right)\right|_{1}  \tag{5.53}\\
& \leq M t_{n}\left|L\left(u_{n}\right) P\left(g, u_{n}\right)\right|_{0}  \tag{5.54}\\
& \leq M^{2} t_{n}\left|P\left(g, u_{n}\right)\right|_{0} \tag{5.55}
\end{align*}
$$

and by (5.42), (5.41)

$$
\begin{align*}
& \left|P\left(g, u_{n}\right)\right|_{0}  \tag{5.56}\\
\leq & \left|P\left(g, u_{n}\right)-P\left(g, u_{n-1}\right)-D_{u} P\left(g, u_{n-1}\right)\left(u_{n}-u_{n-1}\right)\right|_{0}  \tag{5.57}\\
& +\left|P\left(g, u_{n-1}\right)+D_{u} P\left(g, u_{n-1}\right)\left(u_{n}-u_{n-1}\right)\right|_{0}  \tag{5.58}\\
\leq & M\left|u_{n}-u_{n-1}\right|_{1}^{2}+\left|D_{u} P\left(g, u_{n-1}\right)\left(1-S\left(t_{n-1}\right)\right) L\left(u_{n-1}\right) P\left(g, u_{n-1}\right)\right|_{0}  \tag{5.59}\\
\leq & M\left|u_{n}-u_{n-1}\right|_{1}^{2}+M\left|\left(1-S\left(t_{n-1}\right)\right) L\left(u_{n-1}\right) P\left(g, u_{n-1}\right)\right|_{0}  \tag{5.60}\\
\leq & M\left|u_{n}-u_{n-1}\right|_{1}^{2}+M M_{j} t_{n-1}^{1-j}\left|L\left(u_{n-1}\right) P\left(g, u_{n-1}\right)\right|_{j-1}  \tag{5.61}\\
\leq & M\left|u_{n}-u_{n-1}^{2}\right|_{1}^{2}+M M_{j}^{2} t_{n-1}^{1-j}\left(1+\left|u_{n-1}-\bar{u}\right|_{j}\right) \tag{5.62}
\end{align*}
$$

so that

$$
\begin{equation*}
\left|u_{n+1}-u_{n}\right|_{1} \leq t_{n} M^{3}\left|u_{n}-u_{n-1}\right|_{1}^{2}+M^{3} M_{j}^{2} t_{n} t_{n-1}^{1-j}\left(1+\left|u_{n-1}-\bar{u}\right|_{j}\right) . \tag{5.63}
\end{equation*}
$$

For some $\mu>0$ to be determined later, let

$$
\begin{equation*}
\delta_{n}:=t_{n}^{\mu} M^{3}\left|u_{n}-u_{n-1}\right|_{1} . \tag{5.64}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\delta_{n+1} \leq A^{\kappa^{n}(1+\mu(\kappa-2))} \delta_{n}^{2}+e_{n}, \quad e_{n}:=M^{6} M_{j}^{2} A^{\mu \kappa^{n+1}+\kappa^{n}+(1-j) \kappa^{n-1}}\left(1+\left|u_{n-1}-\bar{u}\right|_{j}\right) . \tag{5.65}
\end{equation*}
$$

The parameters $A, \kappa$, etc. will be determined in order to view (5.65) as a perturbation of $x_{n+1} \leq x_{n}^{2}$. With

$$
\begin{equation*}
1+\mu(\kappa-2) \leq 0 \tag{5.66}
\end{equation*}
$$

we have $\delta_{n+1} \leq \delta_{n}^{2}+e_{n}$. By inspection, the graphs of $y=x$ and $y=x^{2}+\frac{1}{8}$ intersect at some $x \in\left[\frac{2}{3}, 1\right]$. Thus, if one can impose

$$
\begin{equation*}
\delta_{1}<\frac{2}{3}, \quad e_{n}<\frac{1}{8} \tag{5.67}
\end{equation*}
$$

the inequality $\delta_{n+1} \leq \delta_{n}^{2}+e_{n}$ guarantees that $\delta_{n}$ is bounded for all $n$ (say by 1 ) and therefore

$$
\begin{equation*}
\left|u_{n}-u_{n-1}\right|_{1} \leq \frac{A^{-\mu \kappa^{n}}}{M^{3}} \tag{5.68}
\end{equation*}
$$

In order to control $e_{n}$ we need to estimate the growth of $1+\left|u_{n}-\bar{u}\right|_{j}$ :

$$
\begin{align*}
1+\left|u_{n+1}-\bar{u}\right|_{j} & \leq 1+\left|u_{n}-\bar{u}\right|_{j}+\left|u_{n+1}-u_{n}\right|_{j}  \tag{5.69}\\
& =1+\left|u_{n}-\bar{u}\right|_{j}+\left|S\left(t_{n}\right) L\left(u_{n}\right) P\left(g, u_{n}\right)\right|_{j}  \tag{5.70}\\
& \leq 1+\left|u_{n}-\bar{u}\right|_{j}+M_{j} t_{n}\left|L\left(u_{n}\right) P\left(g, u_{n}\right)\right|_{j-1}  \tag{5.71}\\
& \leq 1+\left|u_{n}-\bar{u}\right|_{j}+M_{j}^{2} t_{n}\left(1+\left|u_{n}-\bar{u}\right|_{j}\right)  \tag{5.72}\\
& \leq 2 M_{j}^{2} A^{\kappa^{n}}\left(1+\left|u_{n}-\bar{u}\right|_{j}\right) . \tag{5.73}
\end{align*}
$$

Let $\beta \geq 0$ to be determined later and write

$$
\begin{align*}
A^{-\beta \kappa^{n+1}}\left(1+\left|u_{n+1}-\bar{u}\right|_{j}\right) & \leq 2 M_{j}^{2} A^{(-\beta(\kappa-1)+1) \kappa^{n}}\left(A^{-\beta \kappa^{n}}\left(1+\left|u_{n}-\bar{u}\right|_{j}\right)\right)  \tag{5.74}\\
& \leq 2 M_{j}^{2} A^{-\beta(\kappa-1)+1}\left(A^{-\beta \kappa^{n}}\left(1+\left|u_{n}-\bar{u}\right|_{j}\right)\right) \tag{5.75}
\end{align*}
$$

provided $-\beta(\kappa-1)+1<0$. Since we also want the multiplicative factor to be $\leq 1$, we impose the more stringent condition that

$$
\begin{equation*}
(-\beta(\kappa-1)+1) \ln A+\ln \left(2 M_{j}^{2}\right) \leq 0 . \tag{5.76}
\end{equation*}
$$

In turn, as long as the terms $u_{n} \in \mathcal{B}$ exist, we have (recall $u_{0}=\bar{u}, A>1$, and $\beta \geq 0$ )

$$
\begin{equation*}
1+\left|u_{n}-\bar{u}\right|_{j} \leq A^{\beta \kappa^{n}} \tag{5.77}
\end{equation*}
$$

hence $e_{n} \leq M^{6} M_{j}^{2} A^{\left(\mu \kappa^{2}+\kappa+1-j+\beta\right) \kappa^{n-1}}$. In order to satisfy (5.67), we will therefore impose that

$$
\begin{equation*}
\left(\mu \kappa^{2}+\kappa+1-j+\beta\right) \ln A+\ln \left(M^{6} M_{j}^{2}\right)<\ln \frac{1}{8} . \tag{5.78}
\end{equation*}
$$

Finally, we show that the $u_{n} \in \mathcal{B}$ exists for all $n \geq 0$ provided $|P(g, \bar{u})|_{0}<\epsilon$ is sufficiently small, and that the sequence is Cauchy in the $|\cdot|_{1}$-norm. Estimate (5.68) holds provided $\delta_{1}=A^{\mu \kappa} M^{3}\left|u_{1}-u_{0}\right|_{1}<2 / 3$. But $\left|u_{1}-u_{0}\right|_{1} \leq M t_{0}^{2}|L(\bar{u}) P(\bar{u})|_{0} \leq M^{2} A^{2}|P(\bar{u})|_{0}$ so we take

$$
\begin{equation*}
\epsilon<\frac{2}{3} \frac{1}{M^{5} A^{\mu \kappa+2}} . \tag{5.79}
\end{equation*}
$$

This in turn determines the neighborhood $\mathcal{G}$ : since $P(g, \bar{u})=T(\bar{u})-g=\bar{g}-g$,

$$
\begin{equation*}
\mathcal{G}=\left\{g \in \mathcal{Y}| | g-\left.\bar{g}\right|_{0}<\epsilon,|g|_{j-1}<C\right\} . \tag{5.80}
\end{equation*}
$$

Now we verify that $u_{n} \in \mathcal{B}$ is defined for all $n$ :
$\left|u_{n+1}\right|_{1} \leq \sum_{m=0}^{n}\left|u_{m+1}-u_{m}\right|_{1} \leq \frac{1}{M^{3}} \sum_{m=0}^{\infty} A^{-\mu \kappa^{m+1}} \leq \frac{1}{M^{3}} \sum_{m=0}^{\infty} A^{-\mu \kappa(1+m \ln \kappa)}=\frac{1}{M^{3}} \frac{A^{-\mu \kappa}}{1-A^{-\mu \kappa} \ln \kappa}$
(we have used $\kappa^{i} \geq \kappa^{l}+(i-l) \ln \kappa$ ). We thus impose

$$
\begin{equation*}
\frac{1}{M^{3}} \frac{A^{-\mu \kappa}}{1-A^{-\mu \kappa \ln \kappa}}<\eta . \tag{5.82}
\end{equation*}
$$

The sequence is Cauchy since

$$
\begin{align*}
\left|u_{m}-u_{n}\right|_{1} & \leq \sum_{l=n}^{m-1}\left|u_{l+1}-u_{l}\right|_{1}  \tag{5.83}\\
& \leq \frac{1}{M^{3}} \sum_{l=n}^{\infty} A^{-\mu \kappa^{l+1}}  \tag{5.84}\\
& \leq \frac{1}{M^{3}} \sum_{l=n}^{\infty} A^{-\mu \kappa\left(\kappa^{n}+\ln \kappa(l-n)\right)}  \tag{5.85}\\
& =\frac{1}{M^{3}} \frac{A^{-\mu \kappa^{n+1}}}{1-A^{-\mu \kappa \ln \kappa}} \underset{n, m \rightarrow \infty}{\rightarrow} \tag{5.86}
\end{align*}
$$

We denote $u_{\infty}$ its limit in the $|\cdot|_{1}$-norm. Observe from (5.62) that

$$
\begin{equation*}
\left|P\left(g, u_{n}\right)\right|_{0} \leq M\left|u_{n}-u_{n-1}\right|_{1}^{2}+M M_{j}^{2} A^{(1-j+\beta) \kappa^{n-1}} \rightarrow 0 \tag{5.87}
\end{equation*}
$$

since $1-j+\beta<0$ by (5.78). Thus, once $u_{\infty}$ is proven to be the limit in each $|\cdot|_{k}$-norm, the above shows that it is in fact a solution to $P(g, u)=0$ in $\mathcal{B}$.

2 Setting the parameters The constants $M, M_{j}$ are imposed by the problem. The conditions on the parameters $\kappa, \mu, \beta, j$, and $A$ are (5.66), (5.76), (5.78), (5.82):

$$
\begin{gather*}
1+\mu(\kappa-2) \leq 0, \quad 0<\kappa<2  \tag{5.88}\\
(-\beta(\kappa-1)+1) \ln A+\ln \left(2 M_{j}^{2}\right) \leq 0  \tag{5.89}\\
\left(\mu \kappa^{2}+\kappa+1-j+\beta\right) \ln A+\ln \left(M^{6} M_{j}^{2}\right)<\ln \frac{1}{8}  \tag{5.90}\\
\frac{1}{M^{3}} \frac{A^{-\mu \kappa}}{1-A^{-\mu \kappa \ln \kappa}}<\eta . \tag{5.91}
\end{gather*}
$$

These conditions are satisfied if the parameters $\kappa, \mu, \beta, j$, and $A$ are chosen in this order so as to satisfy the following:

$$
\begin{gather*}
1<\kappa<2  \tag{5.92}\\
\mu \geq \frac{1}{2-\kappa}  \tag{5.93}\\
-\beta(\kappa-1)+1<0  \tag{5.94}\\
\mu \kappa^{2}+\kappa+1-j+\beta<0 \tag{5.95}
\end{gather*}
$$

and $A$ sufficiently large so that the three inequalities where it is involved are satisfied. It is not difficult to see that, in order to minimize $j$, it should be chosen the smallest integer
strictly larger than $\frac{\kappa^{2}}{2-\kappa}+\kappa+1+\frac{1}{\kappa-1}$ over $\kappa \in(1,2)$. Using a computer one finds that

$$
\begin{equation*}
j=9 \quad \text { is attained with any } \quad 1.2<\kappa<1.5 \text {. } \tag{5.96}
\end{equation*}
$$

3 The rough solution is smooth We will use that for each $m \geq 1$ there exists a constant $C_{m}$ such that

$$
\begin{equation*}
1+\left|u_{n}-\bar{u}\right|_{m} \leq C_{m} A^{\beta \kappa^{n}} \tag{5.97}
\end{equation*}
$$

which is proven below. The important point is that the inequality holds with the same $\beta$ regardless of $m$.

Fix then $i \geq 1$, and let $m \geq 1$ which will be determined later. Denote $M_{m}$ (or make $M_{j}$ larger if $j-1=m$ ) a constant such that the following estimates hold for any $v \in \mathcal{Y}, u \in \mathcal{B}$, and $t>0$ :

$$
\begin{align*}
|S(t) v|_{m} & \leq M_{m}|v|_{m}  \tag{5.98}\\
|L(u) P(g, u)|_{m} & \leq M_{m}\left(1+|u-\bar{u}|_{m+1}\right) . \tag{5.99}
\end{align*}
$$

By interpolation inequalities, ( $C$ denotes constants depending on $i, m$, but not on $n$ )

$$
\begin{align*}
\left|u_{n+1}-u_{n}\right|_{i} & \leq\left|u_{n+1}-u_{n}\right|_{1}^{\frac{m-i}{m-1}}\left|u_{n+1}-u_{n}\right|^{\frac{i-1}{m-1}}  \tag{5.100}\\
& \leq C A^{-\mu \kappa^{n+1} \frac{m-i}{m-1}}\left|S\left(t_{n}\right) L\left(u_{n}\right) P\left(g, u_{n}\right)\right|_{m^{\frac{i-1}{m-1}}}^{\frac{i-1}{m}}  \tag{5.101}\\
& \leq\left. C A^{-\mu \kappa^{n+1} \frac{m-i}{m-1}} L\left(u_{n}\right) P\left(g, u_{n}\right)\right|_{m-1} ^{m-1}  \tag{5.102}\\
& \leq C A^{-\mu \kappa^{n+1} \frac{m-i}{m-1}}\left(1+\left|u_{n}\right|_{m+1}\right)^{\frac{i-1}{m-1}}  \tag{5.103}\\
& \leq C A^{-\mu \kappa^{n+1} \frac{m-i}{m-1}} A^{\beta \kappa^{n} \frac{i-1}{m-1}}  \tag{5.104}\\
& \leq C A^{(-\mu \kappa(m-i)+\beta(i-1)) \frac{\kappa^{n}}{m-1}} . \tag{5.105}
\end{align*}
$$

Now choosing $m$ sufficiently large that

$$
\begin{equation*}
-\mu \kappa(m-i)+\beta(i-1)<0 \tag{5.106}
\end{equation*}
$$

makes the exponent negative and the increment $\left|u_{n+1}-u_{n}\right|_{i}$ decays with a double exponential rate. It is then easy to see that $u_{n}$ is Cauchy in the $|\cdot|_{i}$-norm.

4 Proof of estimates (5.97) Fix $m \geq 1$. As for (5.73) we find

$$
\begin{equation*}
1+\left|u_{n+1}-\bar{u}\right|_{m} \leq 2 M_{m}^{2} A^{\kappa^{n}}\left(1+\left|u_{n}-\bar{u}\right|_{m}\right) \tag{5.107}
\end{equation*}
$$

so that

$$
\begin{equation*}
A^{-\beta \kappa^{n+1}}\left(1+\left|u_{n+1}-\bar{u}\right|_{m}\right) \leq 2 M_{m}^{2} A^{(-\beta(\kappa-1)+1) \kappa^{n}}\left(A^{-\beta \kappa^{n}}\left(1+\left|u_{n}-\bar{u}\right|_{m}\right)\right) . \tag{5.108}
\end{equation*}
$$

Given $m, \beta, \kappa$, let $n^{*}(m)$ such that for $n \geq n^{*}(m), 2 M_{m}^{2} A^{(-\beta(\kappa-1)+1) \kappa^{n}} \leq 1$ (recall that $-\beta(\kappa-1)+1<0)$. Then choose $C_{m}$ so that $A^{-\beta \kappa^{n}}\left(1+\left|u_{n}-\bar{u}\right|_{m}\right) \leq C_{m}$ for $n<n^{*}(m)$.

### 5.3 The injective part of the Nash-Moser Inverse Function Theorem

Even though the injective part of Nash-Moser Inverse Function Theorem cannot be used as such for the injective part of Theorem [1 it is instructive to show its proof as our result is an adaptation of it. We follow [7].

The assumptions here are different from those of Theorem 21. We are not concerned with existence of solutions, only uniqueness. Thus, $D_{u} P(g, u)$ is assumed injective, with a left-inverse again denoted $L(u)$. Note that the Moser scheme plays no rôle here.

Theorem 22 (Nash-Moser IFT - injective part) Consider $g \in \mathcal{Y}$ in a $|\cdot|_{0}$-neighborhood of $\bar{g}$. Suppose that $D_{u} P(g, u) v$ has a left-inverse $L(u) h$, which is a tame of degree 0 in $g$ and $h$, and 1 in $u$. Then, there exists a $|\cdot|_{1}$-neighborhood $\mathcal{B}^{\prime}:=\left\{u \in \mathcal{X}| | u-\left.\bar{u}\right|_{1}<\eta^{\prime}\right\}$ of $\bar{u}$ such that, if $P\left(g, u_{1}\right)=P\left(g, u_{2}\right)$ where $u_{1}, u_{2} \in \mathcal{B}^{\prime}$, then $u_{1}=u_{2}$.

Proof Use Taylor's formula,
$P\left(g, u_{2}\right)=P\left(g, u_{1}\right)+D_{u} P\left(g, u_{1}\right)\left(u_{2}-u_{1}\right)+\int_{0}^{1}(1-t) D_{u u}^{2} P\left(g, u_{1}+t\left(u_{2}-u_{1}\right)\right)\left(u_{2}-u_{1}, u_{2}-u_{1}\right) d t$
so that

$$
\begin{equation*}
u_{2}-u_{1}=-L\left(u_{1}\right) \int_{0}^{1}(1-t) D_{u u}^{2} P\left(g, u_{1}+t\left(u_{2}-u_{1}\right)\right)\left(u_{2}-u_{1}, u_{2}-u_{1}\right) d t . \tag{5.110}
\end{equation*}
$$

Tame estimates on $D_{u u}^{2} P(u)\left(v_{1}, v_{1}\right)$ give

$$
\begin{equation*}
\left|u_{2}-u_{1}\right|_{0} \leq c\left|u_{2}-u_{1}\right|_{0}\left|u_{2}-u_{1}\right|_{1} \tag{5.111}
\end{equation*}
$$

where the constant is independent of $u_{1}, u_{2} \in \mathcal{B}^{\prime}$ and $g$ in the restricted neighborhood. Making $\eta^{\prime}$ sufficiently small, we can make $c\left|u_{2}-u_{1}\right|_{1}<1$ for any $u_{1}, u_{2} \in \mathcal{B}^{\prime}$. This forces $\left|u_{2}-u_{1}\right|=0$.

### 5.4 Examples of smooth tame maps

We list in this Section some smooth tame maps which are used throughout the present work. In this Section, by smoothness we do mean that derivatives of all orders exist. $K$ denotes a compact subset of Euclidean space with smooth boundary. $\mathcal{V}$ denotes an open subset of some Fréchet space.

## 1 Linear differential operators with constant coefficients

Lemma 23 A linear differential operator of order $r$ with constant coefficients $L: C_{K}^{\infty} \rightarrow$ $C_{K}^{\infty}$ is a smooth tame map of Fréchet spaces: $L: C_{K}^{n+r} \rightarrow C_{K}^{n}$ is continuous for each $n \geq 0$. Lu has degree $r$ and base 0: for all $u \in C_{K}^{\infty}$.

$$
\begin{equation*}
\|L u\|_{n, \alpha} \leq C \cdot\|u\|_{n+r, \alpha}, \quad n \geq 0 . \tag{5.112}
\end{equation*}
$$

## 2 Product of functions

Lemma 24 The bilinear map

$$
B:\left\{\begin{array}{ccc}
C_{K}^{\infty} & \times & C_{K}^{\infty}  \tag{5.113}\\
F & & G
\end{array} \longrightarrow \begin{array}{c}
C_{K}^{\infty} \\
F G
\end{array}\right\}
$$

is a smooth tame map of Fréchet spaces: for each $n \geq 0, B: C_{K}^{n} \times C_{K}^{n} \rightarrow C_{K}^{n}$ is continuous as well as $B: C_{K}^{n, \alpha} \times C_{K}^{n, \alpha} \rightarrow C_{K}^{n, \alpha} . B(F, G)$ has degree 0 in $F$ and $G$, and base 0 :

$$
\begin{equation*}
\|B(F, G)\|_{n, \alpha} \leq C \cdot\left(\|G\|_{0, \alpha}\|F\|_{n, \alpha}+\|F\|_{0, \alpha}\|G\|_{n, \alpha}\right), \quad n \geq 0 \tag{5.114}
\end{equation*}
$$

for all $F$ and $G$. The first derivative is given by

$$
\begin{equation*}
B(F, G) \cdot(f, g)=f G+F g . \tag{5.115}
\end{equation*}
$$

There are obvious generalizations of the above for the product of an arbitrary number of functions $\left(F_{1}, \ldots, F_{l}\right) \mapsto F_{1} \cdots F_{l}$.

3 The Nemitskii operator If $p(x, z)=p: K \times \mathbb{R} \rightarrow \mathbb{R}$ is a smooth function, define $P(F)(x):=p(x, F(x)), x \in K, F \in \mathcal{V} \subset C_{K}^{\infty}$.

Lemma $25 P: \mathcal{V} \rightarrow C_{K}^{\infty}$ is a smooth tame map of Fréchet spaces: for each $n \geq 0$, $P: C_{K}^{n} \rightarrow C_{K}^{n}$ is continuous as well as $P: C_{K}^{n, \alpha} \rightarrow C_{K}^{n, \alpha} . P(F)$ has degree 0 in $F$ and base 1:

$$
\begin{equation*}
\|P(F)\|_{n, \alpha} \leq C \cdot\left(\|F\|_{n, \alpha}+1\right), \quad n \geq 1 \tag{5.116}
\end{equation*}
$$

for $F$ in a neighborhood where $\|F\|_{1, \alpha}$ is bounded. The first derivative $D P(F) \cdot f \in C_{K}^{\infty}$ is given by

$$
\begin{equation*}
(D P(F) \cdot f)(x)=D_{z} p(x, F(x)) f(x), \quad x \in K \tag{5.117}
\end{equation*}
$$

4 Composition of functions Let $K \subset \mathbb{R}^{d}, K^{\prime} \subset \mathbb{R}^{d^{\prime}}, K^{\prime \prime} \subset \mathbb{R}^{d^{\prime \prime}}$ be compact subsets of Euclidean spaces, and let $G_{0} \in C_{\left(K, \mathbb{R}^{d^{\prime}}\right)}^{\infty}$ such that $G_{0}(K) \subset V^{\prime} \subset K^{\prime}$ for some open set $V^{\prime}$. If $G$ is in a suitable $\|\cdot\|_{0}$-neighborhood $\mathcal{V}$ of $G_{0}$, then $G(K) \subset K^{\prime}$. Thus, we may define the composition operator

$$
C:\left\{\begin{array}{ccc}
C_{\left(K^{\prime}, \mathbb{R}^{d^{\prime \prime}}\right)}^{\infty} & \times & \mathcal{V}  \tag{5.118}\\
F & & G
\end{array} \longrightarrow \begin{array}{c}
C_{\left(K, \mathbb{R}^{d^{\prime \prime}}\right)}^{\infty} \\
F \circ G
\end{array}\right\} .
$$

Lemma $26 C$ is a smooth tame map of Fréchet spaces. In the $C^{n}$-grading we have that

$$
\begin{equation*}
C: C_{\left(K^{\prime}, \mathbb{R}^{d^{\prime \prime}}\right)}^{n} \times C_{K, \mathbb{R}^{d^{\prime}}}^{n} \rightarrow C_{\left(K, \mathbb{R}^{d^{\prime \prime}}\right)}^{n}, \quad n \geq 0 \tag{5.119}
\end{equation*}
$$

is continuous, while in the $C^{n, \alpha}$-grading we only have that

$$
\begin{equation*}
C: C_{\left(K^{\prime}, \mathbb{R}^{d^{\prime \prime}}\right)}^{n+1, \alpha} \times C_{K, \mathbb{R}^{d^{\prime}}}^{n, \alpha} \rightarrow C_{\left(K, \mathbb{R}^{d^{\prime \prime}}\right)}^{n, \alpha}, \quad n \geq 0 \tag{5.120}
\end{equation*}
$$

is continuous. For $G$ in $a\|\cdot\|_{1, \alpha}$-neighborhood, and all $F$ (without restriction),

$$
\begin{equation*}
\|F \circ G\|_{n, \alpha} \leq C \cdot\left(\|F\|_{n, \alpha}+\|G\|_{n, \alpha}\|F\|_{1, \alpha}\right), \quad n \geq 1 \tag{5.121}
\end{equation*}
$$

The derivative is given by

$$
\begin{equation*}
D C(F, G) \cdot(f, g)=F^{\prime}(G) g+f(G) \tag{5.122}
\end{equation*}
$$

Proof See [3]. Note that composition $C^{n, \alpha} \times C^{n, \alpha} \rightarrow C^{n, \alpha}$ is well-defined, even though it is not continuous.

5 The inversion operator Let $I$ be a compact interval denote $\mathcal{D}_{I}^{\infty}$ the group of increasing $C^{\infty}$-diffeomorphisms of $I$. Denote $V(F)=F^{-1}$ the inverse of $F \in \mathcal{D}_{I}^{\infty}$.

Lemma 27 The inversion operator

$$
V:\left\{\begin{array}{c}
\mathcal{D}_{I}^{\infty}  \tag{5.123}\\
F
\end{array} \longrightarrow \begin{array}{c}
\mathcal{D}_{I}^{\infty} \\
F^{-1}
\end{array}\right\}
$$

is a smooth tame map. It is continuous $\mathcal{D}_{I}^{n} \rightarrow \mathcal{D}_{I}^{n}$ for each $n \geq 1 . V(F)$ has degree 0 in $F$ and base 1 in the $\|\cdot\|_{n, \alpha}$-grading:

$$
\begin{equation*}
\|V(F)\|_{n, \alpha} \leq C \cdot\left(\|F\|_{n, \alpha}+1\right), \quad n \geq 1 \tag{5.124}
\end{equation*}
$$

for $F$ in a neighborhood where $\|F\|_{1, \alpha}$ is bounded. The first derivative is given by

$$
\begin{equation*}
D V(F) \cdot f=-\frac{f\left(F^{-1}\right)}{F^{\prime}\left(F^{-1}\right)} \tag{5.125}
\end{equation*}
$$

Proof That $V: \mathcal{D}_{I}^{n} \rightarrow \mathcal{D}_{I}^{n}$ is continuous for each $n \geq 1$ is standard, see Example 4.4.6, p. 92, Part I, and Theorem 2.3.5, p. 148, Part II of [7]. We need tame estimates in the $C^{n, \alpha}$, while those of [7] are given in the $C^{n}$-grading.

Let $y=g(x)$ be a diffeomorphism in $\mathcal{D}_{I}^{\infty}$ with inverse $x=f(y)$. Then, $f^{\prime}(y)=\frac{1}{g^{\prime}(x)}$ so that

$$
\begin{align*}
\left|f^{\prime}\left(y_{1}\right)-f^{\prime}\left(y_{2}\right)\right| & \leq\left|\frac{1}{g^{\prime}\left(x_{1}\right)}-\frac{1}{g^{\prime}\left(x_{2}\right)}\right|  \tag{5.126}\\
& \leq \frac{\left|g^{\prime}\left(x_{2}\right)-g^{\prime}\left(x_{1}\right)\right|}{\left|g^{\prime}\left(x_{1}\right) g^{\prime}\left(x_{2}\right)\right|}  \tag{5.127}\\
& \leq C\left[g^{\prime}\right]_{\alpha}\left|x_{2}-x_{1}\right|^{\alpha} \tag{5.128}
\end{align*}
$$

for $g$ in a $\|\cdot\|_{1}$-neighborhood. Since $\left|x_{2}-x_{1}\right| \leq\left\|f^{\prime}\right\|_{0}\left|y_{2}-y_{1}\right|$, we are done with the tame estimates for $n=1$.

Suppose the tame estimates verified for $1 \leq m<n$. From (3.48) and tame estimates on product of functions, we have for $g$ in a $\|\cdot\|_{1, \alpha}$-neighborhood

$$
\begin{align*}
\left\|f^{(n)}\right\|_{0, \alpha} & \leq C\left\|f^{\prime}\right\|_{0, \alpha}^{n} \sum_{k=1}^{n-1} \sum_{j_{1}+\cdots+j_{k}=n}\left\|f^{(k)}\right\|_{0, \alpha}\left\|g^{\left(j_{1}\right)} \circ f\right\|_{0, \alpha} \cdots\left\|g^{\left(j_{k}\right)} \circ f\right\|_{0, \alpha}  \tag{5.129}\\
& \leq C\left\|f^{\prime}\right\|_{0, \alpha}^{n} \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\ldots, j_{k} \geq 1 \\
j_{1}, \ldots, j_{k} \geq 1}}\left\|f^{(k)}\right\|_{0, \alpha}\left\|g^{\left(j_{1}\right)}\right\|_{0, \alpha}\|f\|_{0, \alpha} \cdots\left\|g^{\left(j_{k}\right)}\right\|_{0, \alpha}\|f\|_{0, \bar{k} \cdot} . \\
& \leq C \sum_{k=1}^{n-1} \sum_{\substack{j_{1}+\cdots+j_{k}=n \\
j_{1}, \ldots, j_{k} \geq 1}}\left(1+\left\|g^{(k)}\right\|_{0, \alpha}\right)\|g\|_{j_{1}, \alpha} \cdots\|g\|_{j_{k}, \alpha} . \tag{5.131}
\end{align*}
$$

We interpolate each factor between their $\|\cdot\|_{1, \alpha^{-}}$and $\|\cdot\|_{n, \alpha}$-norms:

$$
\begin{align*}
& \|g\|_{k, \alpha}\|g\|_{j_{1}, \alpha} \cdots\|g\|_{j_{k}, \alpha}  \tag{5.132}\\
\leq & C \cdot\|g\|_{1, \alpha}^{\frac{n-k}{n-1}}\|g\|_{n, \alpha}^{\frac{k-1}{n-1}} \cdot\|g\|_{1, \alpha}^{\frac{n-j_{1}}{n-1}}\|g\|_{n, \alpha}^{\frac{j_{1}-1}{n-1}} \cdots\|g\|_{1, \alpha}^{\frac{n-j_{k}}{n-1}}\|g\|_{n, \alpha}^{\frac{j_{k}-1}{n-1}}  \tag{5.133}\\
\leq & C \cdot\|g\|_{n, \alpha} \tag{5.134}
\end{align*}
$$

since $\|g\|_{1, \alpha}$ remains bounded and $j_{1}+\cdots+j_{k}=n$. The other terms in the sum are handled similarly.

## References

[1] V. I. Arnold, Sur la géométrie différentielle des groupes de Lie de dimension infinie et ses applications à l'hydrodynamique des fluides parfaits, Ann. Inst. Fourier 16 (1966), 316-361.
[2] V. I. Arnold and B. Khesin, Topological Methods in Hydrodynamics, Applied Mathematical Sciences, no. 125, Springer-Verlag, 1998.
[3] R. de la Llave and R. Obaya, Regularity of the composition operator in spaces of Hölder functions, Discrete Contin. Dynam. Systems 5 (1999), no. 1, 157-184.
[4] D. G. Ebin and J. E. Marsden, Groups of diffeomorphisms and the motion of an incompressible fluid, Ann. of Math. (2) 92 (1970), 102-163.
[5] D. G. Ebin, G. Misiołek, and S. C. Preston, Singularities of the exponential map on the volume-preserving diffeomorphism group, Geom. Funct. Anal. 16 (2006), no. 4, 850-868.
[6] H. Federer, Geometric measure theory, Die Grundlehren der mathematischen Wissenschaften, Band 153, Springer-Verlag, 1969.
[7] R. S. Hamilton, The inverse function theorem of Nash and Moser, Bull. Amer. Math. Soc. (N.S.) 7 (1982), no. 1, 65-222.
[8] A. A. Kirillov, Lectures on the orbit method, Graduate Studies in Mathematics, no. 64, American Mathematical Society, 2004.
[9] N. V. Krylov, Lectures on elliptic and parabolic equations in Hölder spaces, Graduate Studies in Mathematics, no. 12, American Mathematical Society, 1996.
[10] C. Marchioro and M. Pulvirenti, Mathematical theory of incompressible nonviscous fluids, Applied Mathematical Sciences, no. 96, Springer-Verlag, 1994.
[11] J. E. Marsden and A. Weinstein, Coadjoint orbits, vortices, and Clebsch variables for incompressible fluids, Order in chaos (Los Alamos, N.M., 1982). Phys. D 7 (1983), no. 1-3, 305-323.
[12] J. Miller, P. Weichman, and M. C. Cross, Statistical mechanics, Euler's equation, and Jupiter's Red Spot, Phys. Rev. A 45 (1992), 2328-2359.
[13] G. Misiołek and S. C. Preston, Fredholm properties of Riemannian exponential maps on diffeomorphism groups, Invent. Math. 179 (2010), no. 1, 191-227.
[14] J. Moser, A new technique for the construction of solutions of nonlinear differential equations, Proc. Nat. Acad. Sci. U.S.A. (1961), no. 47, 1824-1831.
[15] J. Nash, The imbedding problem for Riemannian manifolds, Ann. of Math. (2) 63 (1956), 20-63.
[16] S. C. Preston, The WKB method for conjugate points in the volumorphism group, Indiana Univ. Math. J. 57 (2008), no. 7, 3303-3327.
[17] T. S. Ratiu, R. Tudoran, L. Sbano, E. Sousa Dias, and G. Terra, A crash course in geometric mechanics, notes of the courses given by Ratiu, London Math. Soc. Lecture Note Ser., no. 306, Cambridge Univ. Press, Cambridge, 2005.
[18] R. Robert, A maximum-entropy principle for two-dimensional perfect fluid dynamics, J. Statist. Phys. 65 (1991), no. 3-4, 531-553.
[19] F. Sergeraert, Un théorème de fonctions implicites sur certains espaces de Fréchet et quelques applications, Ann. Sci. École Norm. Sup. (4) 5 (1972), 599-660.
[20] A. I. Shnirelman, Lattice theory and flows of ideal incompressible fluid, Russian J. Math. Phys. 1 (1993), no. 1, 105-114.
[21] B. Turkington, Statistical equilibrium measures and coherent states in two-dimensional turbulence, Comm. Pure Appl. Math. 52 (1999), no. 7, 781-809.


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[^1]:    ${ }^{[1]}$ We refer the reader to monograph [2] for examples.

[^2]:    ${ }^{[2]}$ It is important to point out that the space of vorticities is formally a Poisson manifold, not a symplectic manifold. The orbits $\mathcal{O}_{\omega}$ can be considered as symplectic leaves of this Poisson manifold. See e. g. 11 for details. See also [8, 17. The quantities $I_{f}=\int_{\Omega} f(\omega)$ are Casimir functions. They will also be conserved if $\mathcal{E}$ is replaced by any other Hamiltonian, and they do not generate any symmetries. Such situation typically arises in the process of symplectic reduction, and our situation is an example of this: Euler's equations appear as a result of the reduction of the geodesic flow in the co-tangent bundle of $\mathcal{D}_{\text {vol }}$ by the group $\mathcal{D}_{\text {vol }}$. The space of vorticities can be identified with the dual of the Lie algebra of $\mathcal{D}_{\text {vol }}$, and the orbits $\mathcal{O}_{\omega}$ are the orbits of the co-adjoint representation of $\mathcal{D}_{\text {vol }}$.

[^3]:    ${ }^{[3]}$ The loss of derivative is in the $s$-direction only. To regain this derivative, write the inverse mapping as $z^{-1}=\left(\frac{\omega-\min \omega}{\max \omega-\min \omega}, \theta\right)$, apply a suitable smoothing operator to the second factor $\theta$, and note that this perserves the property (3.2).

