Blow-up of the Solution for a kind of Six Order Hyperbolic and Parabolic Evolution Systems

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Abstract: In this paper, we give some results on the blow-up behaviors of the solution to the mixed problem for some higher nonlinear hyperbolic evolution equation in finite time .By introducing the "blow-up factor $K(u, u_t)$ " we get some new results, which generalize the conclusions of [5], [6] and [7].

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1.Introduction and Lemma

A nonlinear pseudo-hyperbolic equation is important in biological mechanics and other fields. In paper [4], Zhang Jian discussed the generalized neural pseudo –hyperbolic equation $\mathbf{u}_{tt} - \Delta \mathbf{u}_{tt} = \mathbf{F}(\mathbf{x}, \mathbf{t}, \mathbf{u}, \nabla u, \nabla \mathbf{u}_{tt})$.

In this paper, we discuss a mixed problem of six-order system. The results obtained generalize some results of [5], [6] and [7]. Let Ω be a bounded domain of \mathbb{R}^n having sufficiently smooth boundary $\partial\Omega$, γ denotes the derivative for outward normal direction, \mathbf{u} (t, x) is a real

function, and
$$\Delta = \sum_{i=1-}^{n} \frac{\partial^2}{\partial x_i^2}, \Delta^2 = \Delta(\Delta), \dots$$

where $\mathbf{x}=(\mathbf{x}_{1,}\ \mathbf{x}_{2,}\ ...,\mathbf{x}_{n})\in\Omega$,(t, $\mathbf{x})\in R^{+}\times\Omega$.Let \mathbf{D} $\mathbf{u}=(\mathbf{u}_{t},\mathbf{D}_{x}\mathbf{u})=(\mathbf{u}_{t},\mathbf{u}_{x},\mathbf{u}_{x},...,\mathbf{u}_{x}),$

(D u)
$$_{x_i} = (u_{tx_i}, u_{x_1x_i}, ... u_{x_nx_i}), (i=1,2,...n);$$
 and $D_x D u = ((D u)_{x_1}, ... (D u))$

u) $_{x_n}$) .The following lemma (see lemma 2.1 in [5]) will be needed for our discussion below.

Lemma 1 .Let $J \in C'(0,\infty)$, J(0)>0, C>0, $\alpha>0$, such that $J'(t) \geq c |J(t)|^{1+\alpha}$, then there exists a $T_0 < \infty$,

such that $\lim_{t\to T_0^-} J(t) = \infty$.

Throughout this paper, we always suppose that $F, \eta, \xi, R, S, u_{0}, u_{1}, a_{0}$ and b_{\perp} are appropriately smooth.

2. "Blow-up" factor $K(u, u_t)$

Let us consider a class of Six -order hyperbolic evolution system:

$$\begin{cases} u_{tt} - \Delta^3 \eta - S(x) \Delta^2 \eta - S(x) \Delta \eta = F(x, u, Du, D_x Du), x \in \Omega, t > 0, \\ u(0, x) = u_0(x), u(0, x) = u_1(x); \\ [a_1(t, x) \frac{\partial \Delta^2 \eta}{\partial \gamma} + b_1(t, x)u]_{\partial \Omega} = 0, [a_2(t, x) \frac{\partial \Delta \eta}{\partial \gamma} + b_2(t, x)u]_{\partial \Omega} = 0, (1) \end{cases}$$

where $\eta = \eta(u, u_t)$ is an appropriately smooth function of 2-variables. By introducing the blow-up of the above equation ,we obtain some results. We assume that the initial value problem of (1) is compatible with the boundary problem of one, and that

(i).
$$\int_{\Omega} K(u_0(x), u_1(x)) dx > 0;$$

(ii).
$$K_{u_t} \left[\frac{\partial \Delta^2 \eta}{\partial \gamma} \right]_{\partial \Omega} = 0, K_{u_t} S(x) \left[\frac{\partial \Delta \eta}{\partial \gamma} \right]_{\partial \Omega} = 0, K_{u_t} S(x) \left[\frac{\partial \eta}{\partial \gamma} \right]_{\partial \Omega} = 0$$
;

(iii).
$$K_{u_t}$$
. $F \ge c |K(u, u_t)|^{1+\alpha} - K_u .u_t + G$,

where
$$G = \sum_{i=1}^{n} [K_{u_i}]_{x_i} [\Delta^2 \eta]_{x_i} + \sum_{i=1}^{n} [K_{u_i} S]_{x_i} [\Delta \eta]_{x_i} + \sum_{i=1}^{n} [K_{u_i} S]_{x_i} [\eta]_{x_i}$$
, $c > 0, \alpha > 0$,

Theorem 1. Assume that there is a real value function K (u, u_t) satisfying (i),(ii),(iii), and that $u \in C^2(0,T;H^6(\Omega))$, is a solution of the mixed problems (1) where $0 < T \le \infty$. Then there exists a T₀ < ∞ such that

$$\lim_{t\to T_0^-}\int_{\Omega}K(u,u_t)dx=\infty.$$

 $K(u, u_t)$ is called the `Blow-up factor 'of (1).

Proof: We take $J(t) = \int_{\Omega} K(u, u_t) dx$, then

$$J'(t) = \int_{\Omega} [K_u u_t + K_{ut} u_{tt}] dx$$
$$= \int_{\Omega} [K_u u_t + K_{u_t} (\Delta^3 \eta + S(x) \Delta^2 \eta + S(x) \Delta \eta)] dx$$

By using Green identity, we have

$$\int_{\Omega} [K_{u_t}[\Delta^3 \eta] dx = \int_{\partial \Omega} K_{u_t} \left[\frac{\partial \Delta^2 \eta}{\partial \gamma} \right] ds - \sum_{i=1}^n [K_{u_t}]_{x_i} \left[\Delta^2 \eta \right]_{x_i} dx$$

$$\int_{\Omega} K_{u_{t}} S(x) \cdot [\Delta^{2} \eta] dx = \int_{\partial \Omega} K_{u_{t}} S(x) \left[\frac{\partial \Delta \eta}{\partial \gamma}\right] dx$$

$$\mathrm{s-}\sum_{\scriptscriptstyle i=1}^n \ \int_{\Omega} [K_{u_i}\,\mathrm{S}]_{\,x_i}\,. [\Delta\eta\,]_{\,x_i} \ \mathrm{d}\,\mathrm{x}\;.$$

By the same way, we have

$$\int_{\Omega} K_{u_{t}} \quad \mathbf{S} \quad (\mathbf{x}). \quad [\quad \Delta \eta \quad] \quad \mathbf{d} \quad \mathbf{x} = \int_{\partial \Omega} K_{u_{t}} \quad .\mathbf{S}(\mathbf{x})[\quad \frac{\partial \eta}{\partial \gamma} \quad] \mathbf{d} \quad \mathbf{s} \quad -$$

$$\int_{\Omega} \sum_{i=1}^{n} \left[K_{u_{i}} S \right]_{x_{i}} \left[\eta \right]_{x_{i}} dx.$$

By Holder inequality , we get $J'(t) \ge c \int_{\Omega} \ \left| K(u,u_t) \right|^{1+\alpha} \mathrm{d} \ x \ge c \cdot c_1 \left| J(t) \right|^{1+\alpha}$,

where
$$c_1 = (\int_{\Omega} dx)^{-1} > 0$$
. We choose $T_0 = \frac{1}{c.c.\alpha[J(0)]^{\alpha}} < \infty$, by Lemma 1, we

have
$$\lim_{t \to T_0^-} \int_{\Omega} K(u, u_t) dx = \infty$$
. This ends the proof.

3. "Blow-up factor" $K(x, u_t, v_t)$

Now we use the method of "blow-up factor" to discuss the following mixed problem, some more general results are obtained. We consider

$$\begin{cases} u_{tt} - L^3 \eta - S(x) L^2 \eta - T(x) L \eta - R(x) L \xi = f(u, v, u_t), x \in \Omega, t > 0, \\ v_{tt} - L^3 \xi - S(x) L^2 \xi - T(x) L \xi - R(x) L \eta = g(u, v, u_t) \\ u(0, x) = u_0(x), u_t(0, x) = u_1(x), v(0, x) = v_0(x), v_t(0, x) = v_1(x), \\ \left[\frac{\partial L^2 \eta}{\partial \gamma} \right]_{x \in \partial \Omega} = 0, \left[\frac{\partial L^2 \xi}{\partial \gamma} \right]_{x \in \partial \Omega} = 0, \left[\frac{\partial L \eta}{\partial \gamma} \right]_{x \in \partial \Omega} = 0, \left[\frac{\partial L \xi}{\partial \gamma} \right]_{x \in \partial \Omega} = 0, \\ \left[L^2 \eta \right] = 0, \left[L^2 \xi \right] = 0, \left[L \eta \right] = 0, \left[L \xi \right] = 0, x \in \partial \Omega, \end{cases}$$

(2)

where $L = \sum_{i,j=1}^{n} \frac{\partial}{\partial x_i} \left\{ a_{ij}(x) \frac{\partial}{\partial x_j} \right\}$, $|\Omega| = \int_{\Omega} dx$. Throughout this paper we

suppose also that f, g S, R are smooth, and $\eta = \eta(u, u_t)$, $\xi = \xi(v, v_t)$ are real functions of two variables ,we get the following theorems;

Theorem 2. Assume that there is a real value function K (x, u_t, v_t) of 3 - variables satisfying:

$$\int_{\Omega} K(x, u_1(x), v_1(x)) d \times >0 , \text{ and } K_{u_t} f + K_{v_t} g \ge C \left| K(x, u_t, v_t) \right|^{1+\alpha} + G_L,$$

where C>0,
$$\alpha$$
 >0, $G_L = \sum_{i=1}^n [(K_{u_i})_{x_i} (L^2 \eta)_{x_i} + (K_{v_t})_{x_i} (L^2 \xi)_{x_i} + (K_{u_t} S)_{x_i}$
 $(L \eta)_{x_i} + (K_{v_t} S)_{x_i} (L \xi)_{x_i} + (K_{u_t} T)_{x_i} (\eta)_{x_i} + (K_{v_t} T)_{x_i}$
 $(\xi)_{x_i} + (K_{u_t} R)_{x_i} (\xi)_{x_i} + (K_{v_t} R)_{x_i} (\eta)_{x_i}$. Then there exists a $T_0 < \infty$ such that

$$\lim_{t\to T_0^-}\int_{\Omega}K(x,u_t,v_t)\,\mathrm{d}\,\mathrm{x}=\infty\,,$$

where (u(t, x), v(t x)) is the classical solution of the systems (2).

Proof. Let
$$J(t) = \int_{\Omega} K(x, u_t, v_t) dx$$
, then

$$J'(t) = \int_{\Omega} [K_{u_t} \cdot u_{tt} + K_{v_t} \cdot v_{tt}] dx$$

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$$\int_{\Omega} [K_{u_{i}}(L^{3}\eta + S(x)L^{2}\eta + T(x)L\eta + R(x)L\xi) + K_{v_{i}}(L^{3}\xi + S(x)L^{2}\xi + T(x)L\xi + R(x)L\eta)]dx$$

Similar to the proof of Theorem 1, we use Green identity to obtain $J'(t) \ge C \int_{\Omega} \left| K(x, u_t, v_t) \right|^{1+\alpha} dx \quad \text{, that} \quad \text{is} \qquad J'(t) \ge C \left| J(t) \right|^{1+\alpha} \quad \text{.} \quad \text{By Lemma} \quad 1,$ and J(0) > 0 ,we get $\lim_{t \to T_0^-} \int_{\Omega} K(x, u_t, v_t) dx = \infty \text{ .This completes the}$ proof of the theorem. \square

4. "Blow-up factor" K(u, v)

We consider again the coupled systems for higher order hyperbolic and parabolic equations:

$$\begin{cases} u_{tt} - \Delta^3 \eta - S(x) \Delta^2 \eta - T(x) \Delta \eta = f(u, v, u_t), x \in \Omega, t > 0, \\ v_t - \Delta^2 \xi - R(x) \Delta \xi = g(u, v, u_t) \end{cases}$$
(3)

where the initial condition: $u\big|_{t=0} = u_0(x), u_t\big|_{t=0} = u_1(x); \quad v\big|_{t=0} = v_0(x)$, (3)';

And the boundary value condition:

$$\left[a_1(t,x)\frac{\partial\Delta^2\eta}{\partial\gamma}+b_1(t,x)u\right]_{\partial\Omega}=0, \qquad \left[a_2(t,x)\frac{\partial\Delta\eta}{\partial\gamma}+b_2(t,x)u\right]_{\partial\Omega}=0,$$

$$[a_3(t,x)\frac{\partial \eta}{\partial \gamma} + b_3(t,x)]_{\partial\Omega} = 0; \qquad [a_4(t,x)\frac{\partial \Delta \xi}{\partial \gamma} + b_4(t,x)u]_{\partial\Omega} = 0.$$

$$[a_5(t,x)\frac{\partial \xi}{\partial \gamma} + b_5(t,x)u]_{\partial\Omega} = 0.(3)$$

there functions f, g ,S,T,R,...are suite smooth ,and satisfying following conditions

(i)
$$\int_{\Omega} K(u_1(x), v_0(x)) dx > 0$$
,

(ii)
$$K_{u_{i}} \left[\frac{\partial \Delta^{2} \eta}{\partial \gamma} \right]_{\partial \Omega} = 0,$$

$$K_{u_{t}}\left[\frac{\partial\Delta\eta}{\partial\gamma}\right]_{\partial\Omega}=0,\ K_{u_{t}}\left[\frac{\partial\eta}{\partial\gamma}\right]_{\partial\Omega}=0,\ K_{v}\left[\frac{\partial\Delta\xi}{\partial\gamma}\right]_{\partial\Omega}=0,\ K_{v}\left[\frac{\partial\xi}{\partial\gamma}\right]_{\partial\Omega}=0,$$

(iii)
$$K_u.f + K_v.g \ge c |K(u_t,v)|^{1+\alpha} + G, (\alpha > 0),$$

where

$$G = \sum_{i=1}^{n} \{ [K_{u_i}]_{x_i} [\Delta^2 \eta]_{x_i} + [K_{u_i} S]_{x_i} [\Delta \eta]_{x_i} + [K_{u_i} T]_{x_i} [\eta]_{x_i} + [K_{v}]_{x_i} [\Delta \xi]_{x_i} + [K_{v} R]_{x_i} [\xi]_{x_i} \}.$$

Theorem 3. Assume that there is a real value function K (u_t, v) satisfying (i),(iii),(iii), and that $u \in C^2(0,T;H^6(\Omega)), v \in C^2(0,T;H^4(\Omega))$, are a solution of the mixed problems (1) where $0 < T \le \infty$. Then there exists a T $_0$

<∞ such that

$$\lim_{t\to T_0^-}\int_{\Omega}K(u_t,v)dx=\infty\,,$$

where $K(u_t, v)$ is called the "Blow-up factor" of (3).

Proof. Let J (t) = $\int_{\Omega} K(u_t, v) dx$, then

$$J' (t) = \int_{\Omega} (K_{u_t} u_{tt} + K_{v} v_{t}) dx = \int_{\Omega} [K_{u_t} (\Delta^3 \eta + S(x) \Delta^2 \eta + T(x) \Delta \eta)] dx + \int_{\Omega} [K_{v} (\Delta^2 \xi + R(x) \Delta \xi)] dx + \int_{\Omega} [K_{u_t} f + K_{v} g] dx$$

By Green identity we get:

$$\int_{\Omega} K_{u_t}[\Delta^3 \eta] dx = \int_{\partial \Omega} K_{u_t}[\frac{\partial \Delta^2 \eta}{\partial \gamma}] ds - \sum_{i=1}^n \int_{\Omega} [K_{u_t}]_{x_i} [\Delta^2 \eta]_{x_i} dx =$$

$$-\sum_{i=1}^{n}\int_{\Omega} [K_{u_i}]_{x_i}.[\Delta^2 \eta]_{x_i} dx,$$

$$\int_{\Omega} K_{u_t} S(\mathbf{x}). \left[\Delta^2 \eta \right] d \mathbf{x} = \int_{\partial \Omega} K_{u_t} S(\mathbf{x}) \left[\frac{\partial \Delta \eta}{\partial \gamma} \right] d \mathbf{s} - \sum_{i=1}^{n} \int_{\Omega} \left[K_{u_t} S \right]_{x_i} . \left[\Delta \eta \right]_{x_i} d$$

X .

$$= - \sum_{i=1}^{n} \int_{\Omega} [K_{u_i}. S] \qquad x_i \qquad .[\Delta \eta] \qquad x_i \qquad .$$

$$\mathbf{x} \quad \text{,and} \int_{\Omega} K_{u_t}[\Delta \eta] dx = \int_{\partial \Omega} K_{u_t}[\frac{\partial \eta}{\partial \gamma}] ds - \sum_{i=1}^n \int_{\Omega} [K_{u_t}]_{x_i}[\eta]_{x_i} dx.$$

$$\int_{\Omega} K_{v} \cdot [\Delta^{2} \xi] dx = \int_{\partial \Omega} K_{v} \left[\frac{\partial \Delta \xi}{\partial \gamma} \right] ds - \sum_{i=1}^{n} \int_{\Omega} \left[K_{v} \right]_{x_{i}} \left[\Delta \xi \right]_{x_{i}} dx ,$$

$$\int_{\Omega} K_{\nu} R(x) [\Delta \xi] dx = \int_{\partial \Omega} K_{\nu} R[\frac{\partial \xi}{\partial \gamma}] ds - \int_{\Omega} \{ \sum_{i=1}^{n} [K_{\nu} R]_{x_{i}} [\xi]_{x_{i}} \} dx$$

Thus,
$$\int_{\Omega} (K_{u_t}.f + K_{v}.g) dx \ge c \int_{\Omega} |K(u_t,v)|^{1+\alpha} dx,$$

that is $J'(t) \ge c.c_1. |J(t)|^{1+\alpha}$. Where $c_1 = (\int_{\Omega} dx)^{-1} > 0$. By Lemma 1, we get $\lim_{t \to T_0^-} \int_{\Omega} K(u_t, v) dx = \infty$. This ends this proof $.\Box$

Remark .We consider again this problems with no homogeneous condition.

The system of coupled equations for hyperbolic and parabolic are appearing in thermal elastic mechanics. We consider the more generalized stronger type systems for higher order evolution systems:

$$\begin{cases}
 u_{tt} - \beta \Delta^2 u - \Delta u = f(u, v, u_t), & x \in \Omega, t > 0, \\
 v_t - \beta \Delta^2 v - \Delta v = g(u, v, u_t),
\end{cases}$$
(4)

the initial condition : $u|_{t=0} = u_0(x), u_t|_{t=0} = u_1(x), v|_{t=0} = v_0(x),$ (4)'the boundary

condition :
$$\left[\frac{\partial \Delta u}{\partial \gamma}\right]_{\partial \Omega} = k_1(t, x), \left[\frac{\partial v}{\partial \gamma}\right]_{\partial \Omega} = k_2(t, x)$$

$$\left[\frac{\partial u}{\partial \gamma}\right]_{\partial \Omega} = k_3(t,x), \left[\frac{\partial v}{\partial \gamma}\right]_{\partial \Omega} = k_4(t,x), (4).$$

where functions f,g,...etc are suite smooth ,they satisfy conditions:

(i)
$$\int_{\partial \Omega} [ak_1(t,x) + bk_2(t,x)]ds \ge 0$$
, $\int_{\partial \Omega} (ak_3(x,t) + bk_4(x,t))ds \ge 0$,

(ii)
$$af + bg \ge c_1 |au_t + bv|^{1+\alpha}, c_1 > 0, \alpha > 0,$$

(iii)
$$\int_{\Omega} [au_1(x) + bv_0(x)]dx > 0.$$

which generalized some results along the direction in [9], and we get the following theorem.

Theorem 4. suppose that conditions (i), (ii), (iii) holds .If the system (4) exits this solution $u \in C^2(0,T;H^4)$, $v \in C^1(0,T;H^4)$.Then there exists a T₀ $<\infty$ such that

$$\lim_{t\to T_0^-}\int_{\Omega}(au_t+bv)dx=\infty,$$

that is $\lim_{t \to T_0^-} \sup |u_t(x,t)| = \infty$, or $\lim_{t \to T_0^-} \sup |v(x,t)| = \infty$, where a > 0, b>0. In other words, the solution (u, v) of (4) blow-up in the finite time.

Proof. Let J (t)= $\int_{\Omega} (au_t + bv) dx$, then by using condition (i) and Green identity, we obtain:

$$J'$$

$$(t) = \int_{\Omega} (au_{tt} + bv_{t}) dx = \int_{\Omega} \beta (a\Delta^{2}u + b\Delta^{2}v) dx + \int_{\Omega} (a\Delta u + b - \Delta v) dx + \int_{\Omega} (af + bg) dx.$$

$$=\int_{\partial\Omega}\beta[ak_1(x,t)+bk(x,t)]ds+\delta\int_{\partial\Omega}[ak_3(x,t)+bk_4(x,t)]ds+\int_{\Omega}(af+bg)dx,$$
 that is $J'(t)\geq c_1\int_{\Omega}\left|au_t+bv\right|^{1+\alpha}dx$. Thus, by Holder inequality, we get easy that $J'(t)\geq c\left|J(t)\right|^{1+\alpha}$, where $C=c_1(\int_{\Omega}dx)^{-1}$.>0. Similar to the proof of theorem 1, and $J(0)>0$. By Lemma 1, there exists a $T_0<\infty$, such that

$$\lim_{t\to T_0^-} J(t) = \lim_{t\to T_0^-} \int_{\Omega} (au_t + bv) dx = \infty.$$

That is $\lim_{t \to T_0} Sup |u_t(x,t)| = \infty$, or $\lim_{t \to T_0^-} Sup |v(x,t)| = \infty$. Thus the solution (u ,v) of system (4) blow-up in the finite time. This completes this proof. \square Corollary. Let $\beta = 0$, we get the systems of complete equation with 2- order that is also the similar results.

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