## Quantum Physics

## The cryptohermitian smearedcoordinate representation of wave functions

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The one-dimensional real line of coordinates is replaced, for simplification or approximation purposes, by an N-plet of the so called Gauss-Hermite grid points. These grid points are interpreted as the eigenvalues of a tridiagonal matrix \$ $\$$ mathfrak\{q\}_0\$ which proves rather complicated. Via the "zeroth" Dyson-map \$1Omega_0\$ the "operator of position" \$\mathfrak\{q\}_0\$ is then further simplified into an isospectral matrix \$Q_0\$ which is found optimal for the purpose. As long as the latter matrix appears non-Hermitian it is not an observable in the manifestly "false" Hilbert space $\$\{\backslash \text { cal } H\}^{\wedge}\{(F)\}:=\backslash m a t h b b\{R\}$ ${ }^{\wedge} N \$$. For this reason the optimal operator \$Q_0\$ is assigned the family of its isospectral avatars \$\mathfrak\{h\}_lalpha\$, \$\alpha=(0,)<br>,1,2,..\$. They are, by construction, selfadjoint in the respective \$lalpha-\$dependent image Hilbert spaces $\$\{\backslash c a l ~ H\} \wedge\{(P)\}$ _lalpha\$ obtained from $\$\{\backslash c a l ~ H\} \wedge\{(F)\} \$$ by the respective "new" Dyson maps \$\Omega_lalpha\$. In the ultimate step of simplification, the inner product in the F-superscripted space is redefined in an \{lit ad hoc\}, \$lalpha-\$dependent manner. The resulting "simplest", S-superscripted representations $\$\{\backslash \text { cal } H\}^{\wedge}\{(S)\}$ _lalpha\$ of the eligible physical Hilbert spaces of states (offering different dynamics) then emerge as, by construction, unitary equivalent to the (i.e., indistinguishable from the) respective awkward, P superscripted and \$\alpha-\$subscripted physical Hilbert spaces.

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